## Citation for published version:

Benoit Vicedo, and Charles Young, 'Cyclotomic Gaudin models with irregular singularities', Journal of Geometry and Physics, Vol. 121: 247-278, November 2017.

DOI:
https://doi.org/10.1016/j.geomphys.2017.07.013

## Document Version:

This is the Accepted Manuscript version.
The version in the University of Hertfordshire Research Archive may differ from the final published version.

## Copyright and Reuse:

© 2017 Elsevier B.V.
This manuscript version is made available under the terms of the Creative Commons Attribution-NonCommercialNoDerivatives License ( http://creativecommons.org/licenses/by-nc-nd/4.0/ ), which permits non-commercial re-use, distribution, and reproduction in any medium, provided the original work is properly cited, and is not altered, transformed, or built upon in any way.

## Enquiries

If you believe this document infringes copyright, please contact the Research \& Scholarly Communications Team at rsc@herts.ac.uk

# CYCLOTOMIC GAUDIN MODELS WITH IRREGULAR SINGULARITIES 

BENOÎT VICEDO AND CHARLES YOUNG


#### Abstract

Generalizing the construction of the cyclotomic Gaudin algebra from [VY16a], we define the universal cyclotomic Gaudin algebra. It is a cyclotomic generalization of the Gaudin models with irregular singularities defined in [FFT10].

We go on to solve, by Bethe ansatz, the special case in which the Lax matrix has simple poles at the origin and arbitrarily many finite points, and a double pole at infinity.


## Contents

1. Introduction ..... 1
2. The cyclotomic Gaudin model ..... 3
3. Statement of main results ..... 11
4. Proofs ..... 14
Appendix A. Proof of Proposition 2.5 ..... 33
Appendix B. $\quad Y_{W}$-map ..... 38
References ..... 42

Keywords: Quantum integrable model; cyclotomic Gaudin model; irregular singularities; coinvariants; Bethe ansatz

## 1. Introduction

Pick a primitive $T$ th root of unity $\omega \in \mathbb{C}^{\times}$, for some non-negative integer $T$ and let $\Gamma:=\langle\omega\rangle \subset \mathbb{C}^{\times}$ denote a copy of the cyclic group $\mathbb{Z} / T \mathbb{Z}$. Let $\mathfrak{g}$ be a finite-dimensional semisimple Lie algebra and $\sigma: \mathfrak{g} \rightarrow \mathfrak{g}$ an automorphism whose order divides $T$.

Associated to these data is a cyclotomic Gaudin algebra [VY16a]. It is a large commutative subalgebra of $U\left(\mathfrak{g}^{\oplus N}\right)^{\mathfrak{g}^{\sigma}}$, depending on a choice of non-zero marked points $\boldsymbol{z}=\left\{z_{1}, \ldots, z_{N}\right\}$ in the complex plane whose $\Gamma$-orbits are pairwise disjoint. It is generated by a hierarchy of Hamiltonians, among which are quadratic Hamiltonians $\mathcal{H}_{1}, \ldots, \mathcal{H}_{N}$ that have appeared previously in [Skr06, Skr13] - see also [CY07] - and, in the context of cyclotomic KZ equations, in [Bro10]. It defines a quantum integrable model generalizing the quantum Gaudin model [Gau76], to which it reduces in the special case $T=1$.

The cyclotomic Gaudin algebra was constructed in [VY16a] using the technology of coinvariants/conformal blocks of $\widehat{\mathfrak{g}}$-modules of critical level, following [FFR94]. The relevant coinvariants in this case are $\Gamma$-equivariant; see [VY16b].

Now, in fact, this approach using coinvariants naturally gives commutative subalgebras not just of $U\left(\mathfrak{g}^{\oplus N}\right)$ but of the larger algebra $U\left(\bigoplus_{i=1}^{N} \mathfrak{g}\left[\left[t-z_{i}\right]\right]\right)$, where $\mathfrak{g}[[t-z]] \cong \mathfrak{g}[[t]]$ is the half loop algebra. Moreover, in the cyclotomic setting it is natural to include also 0 and $\infty$ as marked points. These are the fixed points of the action of $\Gamma$, and to them one attaches twisted half loop algebras, respectively $\mathfrak{g}[[t]]^{\Gamma}$ and $\left(t^{-1} \mathfrak{g}^{\text {op }}\left[\left[t^{-1}\right]\right]\right)^{\Gamma}$ (see $\S 2.2$ ). The first main result of the present sequel to [VY16a] is thus to construct, in $\S 2$, a large commutative subalgebra

$$
\mathscr{Z}_{\infty, z, 0}(\mathfrak{g}, \sigma)^{\Gamma} \subset U\left(\left(t^{-1} \mathfrak{g}^{\mathrm{op}}\left[\left[t^{-1}\right]\right]\right)^{\Gamma} \oplus \bigoplus_{i=1}^{N} \mathfrak{g}\left[\left[t-z_{i}\right]\right] \oplus(\mathfrak{g}[[t]])^{\Gamma}\right)^{\mathfrak{g}^{\sigma}}
$$

It is the cyclotomic generalization of the universal Gaudin algebra defined in [FFT10].
Quotients of half loop algebras of the form $\mathfrak{g}[[t]] / t^{n} \mathfrak{g}[[t]]$ are called (generalized) Takiff algebras. Taking such quotients of $\mathscr{Z}_{\infty, \boldsymbol{z}, 0}(\mathfrak{g}, \sigma)^{\Gamma}$ one obtains commutative algebras $\mathscr{Z}_{\infty, z, 0}^{n_{\infty}, \boldsymbol{n}, n_{0}}(\mathfrak{g}, \sigma)^{\Gamma}$. ${ }^{1}$ In particular, one recovers the cyclotomic Gaudin algebra of [VY16a] as the special case

$$
\mathscr{Z}_{\infty, z, 0}^{1,(1), 0}(\mathfrak{g}, \sigma)^{\Gamma} \subset U\left(0 \oplus \mathfrak{g}^{\oplus N} \oplus 0\right)^{\mathfrak{g}^{\sigma}} .
$$

More generally, following [FFT10] it is natural to call the integrable models defined by representing $\mathscr{Z}_{\infty, z, 0}^{n_{\infty}, \boldsymbol{n}, n_{0}}(\mathfrak{g}, \sigma)^{\Gamma}$ on tensor products of modules over Takiff algebras, cyclotomic Gaudin models with irregular singularities. Among the simplest possibilities is to introduce one irregular singularity, as mild as possible, at $\infty$; that is, to consider

$$
\mathscr{Z}_{\infty, z, 0}^{2,(1), 1}(\mathfrak{g}, \sigma)^{\Gamma} \subset U\left(\left(\Pi_{-1} \mathfrak{g}\right) \oplus \mathfrak{g}^{\oplus N} \oplus \mathfrak{g}^{\sigma}\right)^{\mathfrak{g}^{\sigma}}
$$

Here $\Pi_{-1} \mathfrak{g} \subset \mathfrak{g}$ denotes the $\omega^{-1}$-eigenspace of $\sigma$. It is to be regarded here as a commutative Lie algebra: it arises as the quotient $\Pi_{-1} \mathfrak{g} \cong_{\mathbb{C}}\left(t^{-1} \mathfrak{g}\left[\left[t^{-1}\right]\right]\right)^{\Gamma} /\left(t^{-2} \mathfrak{g}\left[\left[t^{-1}\right]\right]\right)^{\Gamma}$. Suppose we now pick a one-dimensional representation of this commutative Lie algebra i.e. a linear map $\chi: \Pi_{-1} \mathfrak{g} \rightarrow \mathbb{C}$. Applying this map we obtain a commutative subalgebra

$$
\mathcal{A}_{\infty, z, 0}^{1,(1), 1}(\mathfrak{g}, \sigma, \chi)^{\Gamma} \subset\left(U(\mathfrak{g})^{\otimes N} \otimes U\left(\mathfrak{g}^{\sigma}\right)\right)^{\mathfrak{g}_{\chi}^{\sigma}}
$$

where $\mathfrak{g}_{\chi}^{\sigma}$ denotes the centraliser of $\chi$ under the coadjoint action of $\mathfrak{g}^{\sigma}$ on $\left(\Pi_{-1} \mathfrak{g}\right)^{*}$. In the special case of $N=0$ (i.e. only one marked point, at the origin) this is a cyclotomic generalization of the quantum shift-of-argument subalgebra of [Ryb06]; see also [FFRb10]. The latter is a quantisation of the shift-of-argument subalgebra $\mathcal{A}_{\chi}\left[\right.$ MF78], the Poisson commutative subalgebra of $S(\mathfrak{g}) \simeq P\left(\mathfrak{g}^{*}\right)$ generated by all derivatives of every element of $S(\mathfrak{g})^{\mathfrak{g}}$ in the direction of some fixed $\chi \in \mathfrak{g}^{*}$. The quantum shift-of-argument subalgebra has important connections to $\mathfrak{g}$-crystals, certain limits of $U_{q}(\mathfrak{g})$-modules, and cactus group actions on these, see e.g. [Ryb16]. It would be interesting to investigate these various connections in the cyclotomic setting. (Note also that quantum KdV theory is closely related to an affine shift-of-argument subalgebra [FF07].)

In the remainder of the paper we go on to diagonalize the Hamiltonians generating $\mathcal{A}_{\infty, \boldsymbol{z}, 0}^{1,(1), 1}(\mathfrak{g}, \sigma, \chi)^{\Gamma}$ on tensor products of Verma modules, by means of a Bethe ansatz. We assume that $\chi$, and the highest weights $\lambda_{1}, \ldots, \lambda_{N}, \lambda_{0}$ of these Verma modules, all belong to the dual of a single Cartan subalgebra, and that this Cartan subalgebra is stable under $\sigma$. Under these assumptions one can apply

[^0]the approach to the Bethe ansatz for Gaudin models from [FFR94, Fre05], which uses coinvariants of a particular class of $\widehat{\mathfrak{g}}$-modules at critical level called Wakimoto modules. See $\S 3$, Theorem 3.2, for the precise statement of the result.

Finally, in the special case $\chi=0$ we prove that the Bethe vectors are singular. See Theorem 3.4.
Let us conclude this introduction with some remarks and open questions.
As discussed in [FFT10], the origin of the term irregular singularities comes from the description of the spectrum of Gaudin algebras in terms of opers. The notion of opers with regular singularities was recently extended to the cyclotomic setting in [LV06], and it was conjectured that the spectrum of the cyclotomic Gaudin algebra $\mathscr{Z}_{\infty, z, 0}^{1,(1), 0}(\mathfrak{g}, \sigma)^{\Gamma}$ admits a description in terms of such cyclotomic opers, or $\Gamma$-equivariant opers. It would be interesting to extend the definition of cyclotomic opers to include the case of irregular singularities and relate these to the spectrum of $\mathscr{Z}_{\infty, \boldsymbol{z}, 0}^{n_{\infty}, \boldsymbol{n}, n_{0}}(\mathfrak{g}, \sigma)^{\Gamma}$ defined in the present paper.

The quadratic Hamiltonians of the algebra $\mathcal{A}_{\infty, z, 0}^{1,(1), 1}(\mathfrak{g}, \sigma, \chi)^{\Gamma}$ are the cyclotomic analogs of the Gaudin models considered in [FMTV00], which exhibit a certain bispectrality property. It would be interesting to investigate bispectrality in the cyclotomic setting, in the spirit of that paper.

Cyclotomic analogs of the KP hierarchy were defined recently in [CS]. This construction involves a generalization of (the completion of) Calogero-Moser phase space, which can be seen as a quiver variety whose underlying quiver has a single loop, to quiver varieties for cyclic quivers. CalogeroMoser space is known to be related to Gaudin algebras (Bethe algebras) [MTV14], so it is natural to hope for a similar relation in the cyclotomic setting.

## 2. The cyclotomic Gaudin model

2.1. Rational functions and formal series. We work over $\mathbb{C}$. For any formal variable $t$, we have the ring of polynomials $\mathbb{C}[t]$, the ring of formal power series $\mathbb{C}[[t]]$, and the field of formal Laurent series $\mathbb{C}((t))$. Given a finite collection of points $\boldsymbol{x}=\left\{x_{1}, \ldots, x_{p}\right\} \subset \mathbb{C}$ let $\mathbb{C}_{\infty, \boldsymbol{x}}(t)$ denote the localization of $\mathbb{C}[t]$ by the multiplicative subset generated by $t-x_{1}, \ldots, t-x_{p}$. Elements of $\mathbb{C}_{\infty, \boldsymbol{x}}(t)$ are rational functions in $t$ with poles at most at the points $x_{1}, \ldots, x_{p}$ and at infinity.

For any $z \in \mathbb{C}$ we have the map $\iota_{t-z}: \mathbb{C}_{\infty, x}(t) \rightarrow \mathbb{C}((t-z))$ which returns the Laurent expansion $\iota_{t-z} f(t)$ of a rational function $f(t)$ about $t=z$. We have also $\iota_{t^{-1}}: \mathbb{C}_{\infty, \boldsymbol{x}}(t) \rightarrow \mathbb{C}\left(\left(t^{-1}\right)\right)$ which returns the Laurent expansion $\iota_{t^{-1}} f(t)$ of $f(t)$ in powers of $t^{-1}$. The maps $\iota_{t-z}$ and $\iota_{t^{-1}}$ are both injective homomorphisms of $\mathbb{C}$-algebras.

Let $\operatorname{res}_{t}: \mathbb{C}((t)) \rightarrow \mathbb{C}$ be the map which returns the coefficient of $t^{-1}$. For any $f(t) \in \mathbb{C}_{\infty, \boldsymbol{x}}(t)$ we have

$$
\begin{equation*}
-\operatorname{res}_{t^{-1}} t^{2} \iota_{t^{-1}} f(t)+\sum_{i=1}^{p} \operatorname{res}_{t-x_{i}} \iota_{t-x_{i}} f(t)=0 . \tag{2.1}
\end{equation*}
$$

(This is equivalent to the statement that the sum of the residues of a meromorphic one-form $f(t) d t$ on $\mathbb{C} P^{1}$ vanishes.)
2.2. Opposite Lie algebras and left vs. right modules. Given a complex Lie algebra $\mathfrak{a}$ with Lie product $[\cdot, \cdot]: \mathfrak{a} \otimes \mathfrak{a} \rightarrow \mathfrak{a}$ we write $\mathfrak{a}^{\text {op }}$ for the opposite Lie algebra, namely the vector space $\mathfrak{a}$ endowed with the Lie product $[X, Y]^{\mathrm{op}}:=[Y, X]$. The Lie algebras $\mathfrak{a}$ and $\mathfrak{a}^{\mathrm{op}}$ are isomorphic (by e.g. $X \mapsto-X$ ) but it will be useful to regard them as two distinct Lie algebra structures on the
same underlying vector space. Modules over $\mathfrak{a}$ are naturally identified with left modules over the envelope $U(\mathfrak{a})$; modules over $\mathfrak{a}^{\text {op }}$ are naturally identified with right modules over the envelope $U(\mathfrak{a})$.
2.3. Marked points and the group $\Gamma$. Let $\omega$ be a root of unity of order $T \in \mathbb{Z}_{\geq 1}$. The cyclic group $\Gamma:=\langle\omega\rangle \cong \mathbb{Z} / T \mathbb{Z}$ acts on the Riemann sphere $\mathbb{C} P^{1}=\mathbb{C} \cup\{\infty\}$ by multiplication. The fixed points of this action are 0 and $\infty$. Pick $N \in \mathbb{Z}_{\geq 0}$ points $z_{1}, \ldots, z_{N} \in \mathbb{C} P^{1} \backslash\{0, \infty\}$ whose $\Gamma$-orbits are disjoint: $\Gamma z_{i} \cap \Gamma z_{j}=\emptyset$ whenever $i \neq j$. We write $\boldsymbol{z}=\left\{z_{1}, \ldots, z_{N}\right\}$.

Let $\mathfrak{g}$ be a finite-dimensional simple Lie algebra over $\mathbb{C}$ and $\sigma: \mathfrak{g} \rightarrow \mathfrak{g}$ an automorphism of $\mathfrak{g}$ whose order divides $T$. Let $\langle\cdot, \cdot\rangle$ denote the Killing form on $\mathfrak{g}$ normalised such that long roots have square length 2 . This is $\sigma$-invariant since for any pair of elements $X, Y \in \mathfrak{g}$ we have $\operatorname{tr}(\operatorname{ad}(\sigma X) \circ \operatorname{ad}(\sigma Y))=$ $\operatorname{tr}\left(\sigma \circ \operatorname{ad} X \circ \operatorname{ad} Y \circ \sigma^{-1}\right)=\operatorname{tr}(\operatorname{ad} X \circ \operatorname{ad} Y)$ using the fact that $\operatorname{ad}(\sigma X)=\sigma \circ \operatorname{ad} X \circ \sigma^{-1}$. Let $\Pi_{k}$, $k \in \mathbb{Z} / T \mathbb{Z}$, be the projectors

$$
\begin{equation*}
\Pi_{k}:=\frac{1}{T} \sum_{m=0}^{T-1} \omega^{-m k} \sigma^{m}: \mathfrak{g} \rightarrow \mathfrak{g} \tag{2.2}
\end{equation*}
$$

onto the eigenspaces of $\sigma$. They obey $\sum_{k \in \mathbb{Z} / T \mathbb{Z}} \Pi_{k}=$ id. We write $\mathfrak{g}^{\sigma}$ for the subalgebra of invariants,

$$
\mathfrak{g}^{\sigma}:=\Pi_{0} \mathfrak{g}
$$

Denote by $\mathfrak{g}_{\infty, \boldsymbol{z}, 0}^{\Gamma, k}, k \in \mathbb{Z} / T \mathbb{Z}$, the Lie algebra of those $\mathfrak{g}$-valued rational functions $f(t)$ of a formal variable $t$ that have no poles outside the set of points $\{0, \infty\} \cup \Gamma \boldsymbol{z}$ and that obey the equivariance condition $\omega^{*} f=\omega^{k} \sigma f$, i.e.

$$
f(\omega t)=\omega^{k} \sigma f(t)
$$

Let also $\left(\mathfrak{g} \otimes \mathbb{C}\left(\left(t^{ \pm 1}\right)\right)\right)^{\Gamma, k}:=\left\{f\left(t^{ \pm 1}\right) \in \mathfrak{g} \otimes \mathbb{C}\left(\left(t^{ \pm 1}\right)\right): f\left(\omega^{ \pm 1} t^{ \pm 1}\right)=\omega^{k} \sigma f\left(t^{ \pm 1}\right)\right\}$. For brevity we write $\mathfrak{g}_{\infty, \boldsymbol{z}, 0}^{\Gamma}:=\mathfrak{g}_{\infty, \boldsymbol{z}, 0}^{\Gamma, 0}$ and $\left(\mathfrak{g} \otimes \mathbb{C}\left(\left(t^{ \pm 1}\right)\right)\right)^{\Gamma}:=\left(\mathfrak{g} \otimes \mathbb{C}\left(\left(t^{ \pm 1}\right)\right)\right)^{\Gamma, 0}$, etc.

There is an injective homomorphism of Lie algebras

$$
\mathfrak{g} \otimes \mathbb{C}_{\infty, z, 0}(t) \longrightarrow \mathfrak{g}^{\mathrm{op}} \otimes \mathbb{C}\left(\left(t^{-1}\right)\right) \oplus \bigoplus_{i=1}^{N} \mathfrak{g} \otimes \mathbb{C}\left(\left(t-z_{i}\right)\right) \oplus \mathfrak{g} \otimes \mathbb{C}((t))
$$

defined by

$$
f(t) \longmapsto\left(-\iota_{t^{-1}} f(t) ; \iota_{t-z_{1}} f(t), \ldots, \iota_{t-z_{N}} f(t) ; \iota_{t}(t)\right)
$$

(note the op and minus sign in our conventions).
Lemma 2.1 ( $\Gamma$-equivariant Strong residue theorem). A tuple of formal series

$$
\left(f_{\infty} ; f_{z_{1}}, \ldots, f_{z_{N}} ; f_{0}\right) \in\left(\mathfrak{g}^{\mathbf{o p}} \otimes \mathbb{C}\left(\left(t^{-1}\right)\right)\right)^{\Gamma, k} \oplus \bigoplus_{i=1}^{N} \mathfrak{g} \otimes \mathbb{C}\left(\left(t-z_{i}\right)\right) \oplus(\mathfrak{g} \otimes \mathbb{C}((t)))^{\Gamma, k}
$$

belongs to $\left(-\iota_{t-1} ; \iota_{t-z_{1}}, \ldots, \iota_{t-z_{N}} ; \iota_{t}\right)\left(\mathfrak{g}_{\infty, \boldsymbol{z}, 0}^{\Gamma, k}\right)$, i.e. they are the Laurent expansions of some rational function in $\mathfrak{g}_{\infty, \boldsymbol{z}, 0}^{\Gamma, k}$, if and only if

$$
\begin{equation*}
-\frac{1}{T} \operatorname{res}_{t^{-1}} t^{2}\left\langle f_{\infty},-\iota_{t^{-1}} g(t)\right\rangle+\sum_{i=1}^{N} \operatorname{res}_{t-z_{i}}\left\langle f_{z_{i}}, \iota_{t-z_{i}} g(t)\right\rangle+\frac{1}{T} \operatorname{res}_{t}\left\langle f_{0}, \iota_{t}(g)\right\rangle=0 \tag{2.3}
\end{equation*}
$$

for all $g(t) \in \mathfrak{g}_{\infty, \boldsymbol{z}, 0}^{\Gamma,-k-1}$.

Proof. The proof is as in [VY16a, Lemma A.1], but including the poles at $\infty$. Compare (2.1).

Let $\widehat{\mathfrak{g}}_{z_{i}}$ denote the extension of $\mathfrak{g} \otimes \mathbb{C}\left(\left(t-z_{i}\right)\right)$ by a one-dimensional centre $\mathbb{C} K_{z_{i}}$, defined by the cocycle

$$
\begin{equation*}
\Omega_{z_{i}}\left(f_{z_{i}}, g_{z_{i}}\right):=\operatorname{res}_{t-z_{i}}\left\langle f_{z_{i}}, \partial_{t-z_{i}} g_{z_{i}}\right\rangle K_{z_{i}}, \quad f_{z_{i}}, g_{z_{i}} \in \mathfrak{g} \otimes \mathbb{C}\left(\left(t-z_{i}\right)\right) . \tag{2.4}
\end{equation*}
$$

Thus, each $\widehat{\mathfrak{g}}_{z_{i}}, i=1, \ldots, N$, is a copy of the (untwisted) affine Lie algebra $\widehat{\mathfrak{g}}$.
Let $\widehat{\mathfrak{g}}_{0}^{\Gamma}$ denote the extension of $(\mathfrak{g} \otimes \mathbb{C}((t)))^{\Gamma}$ by a one-dimensional centre $\mathbb{C} K_{0}$, defined by the cocycle

$$
\begin{equation*}
\Omega_{0}\left(f_{0}, g_{0}\right):=\operatorname{res}_{t}\left\langle f_{0}, \partial_{t} g_{0}\right\rangle K_{0}, \quad f_{0}, g_{0} \in(\mathfrak{g} \otimes \mathbb{C}((t)))^{\Gamma} . \tag{2.5}
\end{equation*}
$$

Let $\widehat{\mathfrak{g}}_{\infty}^{\text {op, }, \Gamma}$ denote the extension of $\left(\mathfrak{g}^{\text {op }} \otimes \mathbb{C}\left(\left(t^{-1}\right)\right)\right)^{\Gamma}$ by a one-dimensional centre $\mathbb{C} K_{\infty}$, defined by the cocycle

$$
\begin{align*}
\Omega_{\infty}\left(f_{\infty}, g_{\infty}\right) & :=\operatorname{res}_{t^{-1}}\left\langle f_{\infty}, \partial_{t^{-1}} g_{\infty}\right\rangle K_{\infty} \\
& =-\operatorname{res}_{t^{-1}} t^{2}\left\langle f_{\infty}, \partial_{t} g_{\infty}\right\rangle K_{\infty}, \quad f_{\infty}, g_{\infty} \in\left(\mathfrak{g}^{\mathrm{op}} \otimes \mathbb{C}\left(\left(t^{-1}\right)\right)\right)^{\Gamma} . \tag{2.6}
\end{align*}
$$

Given any $X \in \mathfrak{g}, n \in \mathbb{Z}$ we introduce the notations

$$
X[n]_{z_{i}}:=X \otimes\left(t-z_{i}\right)^{n} \in \widehat{\mathfrak{g}}_{z_{i}}, \quad X[n]_{0}:=X \otimes t^{n} \in \widehat{\mathfrak{g}}_{0}, \quad X[n]_{\infty}:=X \otimes t^{n} \in \widehat{\mathfrak{g}}_{\infty} .
$$

Note in particular our conventions for the $n^{\text {th }}$-modes at $\infty$.
The algebras $\widehat{\mathfrak{g}}_{0}^{\Gamma}$ and $\hat{\mathfrak{g}}_{\infty}^{\Gamma}$ are both copies of an algebra $\hat{\mathfrak{g}}^{\Gamma}$ which is either a twisted affine Lie algebra (if $\sigma$ is an outer automorphism) or else isomorphic to $\widehat{\mathfrak{g}}$ (if $\sigma$ is an inner automorphism).

Let $\widehat{\mathfrak{g}}_{\infty, N, 0}$ denote the extension of $\left(\mathfrak{g}^{\text {op }} \otimes \mathbb{C}\left(\left(t^{-1}\right)\right)\right)^{\Gamma} \oplus \bigoplus_{i=1}^{N} \mathfrak{g} \otimes \mathbb{C}\left(\left(t-z_{i}\right)\right) \oplus(\mathfrak{g} \otimes \mathbb{C}((t)))^{\Gamma}$, by a one-dimensional centre $\mathbb{C} K$, defined by the cocycle

$$
\begin{equation*}
\Omega(f, g):=\left(-\frac{1}{T} \operatorname{res}_{t^{-1}} t^{2}\left\langle f_{\infty}, \partial_{t} g_{\infty}\right\rangle+\sum_{i=1}^{N} \operatorname{res}_{t-z_{i}}\left\langle f_{z_{i}}, \partial_{t} g_{z_{i}}\right\rangle+\frac{1}{T} \operatorname{res}_{t}\left\langle f_{0}, \partial_{t} g_{0}\right\rangle\right) K, \tag{2.7}
\end{equation*}
$$

where $f=\left(f_{\infty} ; f_{z_{1}}, \ldots, f_{z_{N}} ; f_{0}\right)$ and $g=\left(g_{\infty} ; g_{z_{1}}, \ldots, g_{z_{N}} ; g_{0}\right)$ are in $\left(\mathfrak{g}^{\text {op }} \otimes \mathbb{C}\left(\left(t^{-1}\right)\right)\right)^{\Gamma} \oplus \bigoplus_{i=1}^{N} \mathfrak{g} \otimes$ $\mathbb{C}\left(\left(t-z_{i}\right)\right) \oplus(\mathfrak{g} \otimes \mathbb{C}((t)))^{\Gamma}$. In other words, $\widehat{\mathfrak{g}}_{\infty, N, 0}$ is the quotient of the direct sum $\widehat{\mathfrak{g}}_{\infty}^{\circ p, \Gamma} \oplus \bigoplus_{i=1}^{N} \widehat{\mathfrak{g}}_{z_{i}} \oplus \widehat{\mathfrak{g}}_{0}^{\Gamma}$ by the ideal spanned by $K_{z_{i}}-T K_{0}, i=1, \ldots, N$, and $K_{\infty}-K_{0}$, leaving one central generator, say $K_{z_{1}}$, which we call $K$.

We have an embedding of Lie algebras

$$
\left(-\iota_{t^{-1}} ; \iota_{t-z_{1}}, \ldots, \iota_{t-z_{N}} ; \iota_{t}\right): \mathfrak{g}_{\infty, z, 0}^{\Gamma} \longleftrightarrow\left(\mathfrak{g}^{\text {op }} \otimes \mathbb{C}\left(\left(t^{-1}\right)\right)\right)^{\Gamma} \oplus \bigoplus_{i=1}^{N} \mathfrak{g} \otimes \mathbb{C}\left(\left(t-z_{i}\right)\right) \oplus(\mathfrak{g} \otimes \mathbb{C}((t)))^{\Gamma} .
$$

By Lemma 2.1 the restriction of the cocycle $\Omega$ to the image of $\mathfrak{g}_{\infty, z, 0}^{\Gamma}$ under this embedding vanishes. Therefore the embedding lifts to an embedding

$$
\begin{equation*}
\mathfrak{g}_{\infty, z, 0}^{\Gamma} \longleftrightarrow \widehat{\mathfrak{g}}_{\infty, N, 0} . \tag{2.8}
\end{equation*}
$$

2.4. Induced $\widehat{\mathfrak{g}}_{\infty, N, 0}$-modules. Let $\mathcal{M}_{z_{i}}$ be a module over $\mathfrak{g} \otimes \mathbb{C}\left[\left[t-z_{i}\right]\right]$, for each $i=1, \ldots, N$. We then make it into a module over $\mathfrak{g} \otimes \mathbb{C}\left[\left[t-z_{i}\right]\right] \oplus \mathbb{C} K_{z_{i}}$ by declaring that $K_{z_{i}}$ acts by multiplication by $k \in \mathbb{C}$. Then we have the induced left $U\left(\widehat{\mathfrak{g}}_{z_{i}}\right)$-module,

$$
\begin{equation*}
\mathbb{M}_{z_{i}}^{k}:=U\left(\widehat{\mathfrak{g}}_{z_{i}}\right) \otimes_{\left.U\left(\mathfrak{g} \otimes \mathbb{C}\left[t-z_{i}\right]\right] \oplus \mathbb{C} K_{z_{i}}\right)} \mathcal{M}_{z_{i}} . \tag{2.9}
\end{equation*}
$$

Let $\mathcal{M}_{0}$ be a module over $(\mathfrak{g} \otimes \mathbb{C}[[t]])^{\Gamma}$. We make it into a module over $(\mathfrak{g} \otimes \mathbb{C}[[t]])^{\Gamma} \oplus \mathbb{C} K_{0}$ by declaring that $K_{0}$ acts by multiplication by $k / T \in \mathbb{C}$. Then we have the induced left $U\left(\hat{\mathfrak{g}}_{0}^{\Gamma}\right)$-module,

$$
\begin{equation*}
\mathbb{M}_{0}^{k / T}:=U\left(\hat{\mathfrak{g}}_{0}^{\Gamma}\right) \otimes_{\left.U((\mathfrak{g} \otimes \mathbb{C}[t t]])^{\Gamma} \oplus \mathbb{C} K_{0}\right)} \mathcal{M}_{0} \tag{2.10}
\end{equation*}
$$

Let $\mathcal{M}_{\infty}$ be a module over $\left(\mathfrak{g}^{\text {op }} \otimes t^{-1} \mathbb{C}\left[\left[t^{-1}\right]\right]\right)^{\Gamma}$. We make it into a module over $\left(\mathfrak{g}^{\text {op }} \otimes t^{-1} \mathbb{C}\left[\left[t^{-1}\right]\right]\right)^{\Gamma} \oplus$ $\mathbb{C} K_{\infty}$ by declaring that $K_{\infty}$ acts by multiplication by $k / T \in \mathbb{C}$. We have the induced left $U\left(\widehat{\mathfrak{g}}_{\infty}^{\circ p, \Gamma}\right)$ module,

$$
\begin{equation*}
\mathbb{M}_{\infty}^{k / T}:=U\left(\widehat{\mathfrak{g}}_{\infty}^{\mathrm{op}, \Gamma}\right) \otimes_{\left.U\left(\left(\mathfrak{g}^{\mathrm{op}} \otimes t^{-1} \mathbb{C}\left[\left[t^{-1}\right]\right]\right)\right)^{\Gamma} \oplus \mathbb{C} K_{\infty}\right)} \mathcal{M}_{\infty} \tag{2.11}
\end{equation*}
$$

Let us write

$$
\begin{equation*}
\mathcal{M}:=\mathcal{M}_{\infty} \otimes \bigotimes_{i=1}^{N} \mathcal{M}_{z_{i}} \otimes \mathcal{M}_{0} \tag{2.12}
\end{equation*}
$$

and

$$
\begin{equation*}
\widehat{\mathfrak{g}}_{\infty, N, 0}^{++}:=\left(\mathfrak{g}^{\mathrm{op}} \otimes t^{-1} \mathbb{C}\left[\left[t^{-1}\right]\right]\right)^{\Gamma} \oplus \bigoplus_{i=1}^{N} \mathfrak{g} \otimes \mathbb{C}\left[\left[t-z_{i}\right]\right] \oplus(\mathfrak{g} \otimes \mathbb{C}[[t]])^{\Gamma} \oplus \mathbb{C} K . \tag{2.13}
\end{equation*}
$$

Then the tensor product

$$
\begin{equation*}
\mathbb{M}:=\mathbb{M}_{\infty}^{k / T} \otimes \bigotimes_{i=1}^{N} \mathbb{M}_{z_{i}}^{k} \otimes \mathbb{M}_{0}^{k / T}=U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}\right) \otimes_{U\left(\mathfrak{g}_{\infty, N, 0}^{+}\right)} \mathcal{M} \tag{2.14}
\end{equation*}
$$

is a module over $\widehat{\mathfrak{g}}_{\infty, N, 0}$ on which $K$ acts as $k$. Pulling back by the embedding (2.8), we have that $\mathbb{M}$ becomes a module over $\mathfrak{g}_{\infty, \boldsymbol{z}, 0}^{\Gamma}$ and we can form the space of coinvariants,

$$
\mathbb{M} / \mathfrak{g}_{\infty, z, 0}^{\Gamma}:=\mathbb{M} /\left(\mathfrak{g}_{\infty, z, 0}^{\Gamma} \cdot \mathbb{M}\right)
$$

Proposition 2.2. The Lie algebras $\widehat{\mathfrak{g}}_{\infty, N, 0}^{+}$and $\mathfrak{g}_{\infty, z, 0}^{\Gamma}$ embed as a pair of complementary Lie subalgebras in $\widehat{\mathfrak{g}}_{\infty, N, 0}$, i.e.

$$
\widehat{\mathfrak{g}}_{\infty, N, 0}=\widehat{\mathfrak{g}}_{\infty, N, 0}^{++}+\mathfrak{g}_{\infty, \boldsymbol{z}, 0}^{\Gamma}
$$

as vector spaces. Therefore there is a canonical isomorphism of vector spaces

$$
\begin{equation*}
\mathbb{M} / \mathfrak{g}_{\infty, \boldsymbol{z}, 0}^{\Gamma} \cong_{\mathbb{C}} \mathcal{M} \tag{2.15}
\end{equation*}
$$

Proof. As in [VY16a], Lemma 2.1 and Corollary 2.4.
2.5. Vacuum verma module $\mathbb{V}_{0}^{k}$. Now let $u \in \mathbb{C}^{\times}$be an additional non-zero marked point, whose orbit $\Gamma u$ is disjoint from $\Gamma \boldsymbol{z}$. Then we have the algebras $\mathfrak{g}_{\infty, \boldsymbol{z}, u, 0}^{\Gamma}, \hat{\mathfrak{g}}_{\infty, N+1,0}$, etc. defined as above but with the point $u$ included. To the point $u$ we assign a copy of the vacuum Verma module $\mathbb{V}_{0}^{k}$ over the local copy $\widehat{\mathfrak{g}}_{u}$ of the affine Lie algebra $\widehat{\mathfrak{g}}$. Recall that by definition $\mathbb{V}_{0}^{k}$ is the induced module

$$
\begin{equation*}
\mathbb{V}_{0}^{k}:=U\left(\widehat{\mathfrak{g}}_{u}\right) \otimes_{\left.U(\mathfrak{g} \otimes \mathbb{C}[t-u]] \oplus \mathbb{C} K_{u}\right)} \mathbb{C} v_{0} \tag{2.16}
\end{equation*}
$$

Here $\mathbb{C} v_{0}$ denotes the one-dimensional module over $\mathfrak{g} \otimes \mathbb{C}[[t-u]] \oplus \mathbb{C} K_{u}$ on which $\mathfrak{g} \otimes \mathbb{C}[[t-u]]$ acts trivially and $K_{u}$ acts by multiplication by $k \in \mathbb{C}$.

A vector $X \in \mathbb{V}_{0}^{k}$ is singular if $A . X=0$ for all $A \in \mathfrak{g} \otimes \mathbb{C}[t]$. The singular vectors form a linear subspace of $\mathbb{V}_{0}^{k}$ denoted $\mathfrak{z}(\widehat{\mathfrak{g}})$.

Proposition 2.3. There is a canonical isomorphism of vector spaces

$$
\left(\mathbb{M} \otimes \mathbb{V}_{0}^{k}\right) / \mathfrak{g}_{\infty, z, u, 0}^{\Gamma} \cong_{\mathbb{C}} \mathcal{M} \otimes \mathbb{C} v_{0} \cong_{\mathbb{C}} \mathcal{M}
$$

It follows that, given any $X \in \mathbb{V}_{0}^{k}$, there is a linear map $X(u): \mathcal{M} \rightarrow \mathcal{M}$ defined by

$$
\begin{equation*}
\mathcal{M} \longrightarrow \mathbb{M} \xrightarrow{\otimes X} \mathbb{M} \otimes \mathbb{V}_{0}^{k} \longrightarrow\left(\mathbb{M} \otimes \mathbb{V}_{0}^{k}\right) / \mathfrak{g}_{\infty, \boldsymbol{z}, u, 0}^{\Gamma} \xrightarrow{\sim} \mathcal{M} \tag{2.17}
\end{equation*}
$$

where $\mathcal{M} \hookrightarrow \mathbb{M}$ is the natural embedding. The map $X(u)$ depends rationally on $u$, with poles at most at the points $0, \omega^{k} z_{i}(1 \leq i \leq N$ and $k \in \mathbb{Z} / T \mathbb{Z})$ and $\infty$.
2.6. Generalized Takiff algebras. For any $n \in \mathbb{Z}_{\geq 1}$ there is an ideal $\mathfrak{g} \otimes t^{n} \mathbb{C}[t] \subset \mathfrak{g} \otimes \mathbb{C}[t]$. Define the Lie algebra $\mathcal{T}_{n} \mathfrak{g}$ to be the quotient

$$
\begin{aligned}
\mathcal{T}_{n} \mathfrak{g} & :=(\mathfrak{g} \otimes \mathbb{C}[t]) /\left(\mathfrak{g} \otimes t^{n} \mathbb{C}[t]\right) \\
& \cong_{\mathbb{C}} \mathfrak{g} \oplus t \mathfrak{g} \cdots \oplus t^{n-1} \mathfrak{g} .
\end{aligned}
$$

Thus $\mathcal{T}_{1} \mathfrak{g}=\mathfrak{g}$. The Lie algebra $\mathcal{T}_{n} \mathfrak{g}$ is known as a (generalized) Takiff algebra.
The Lie algebras $\mathcal{T}_{n} \mathfrak{g}$ together with the canonical projections $\mathcal{T}_{n} \mathfrak{g} \rightarrow \mathcal{T}_{m} \mathfrak{g}, n>m$, form an inverse system, and $\mathfrak{g} \otimes \mathbb{C}[[t]]$ is the inverse limit $\lim ^{\sim} \mathcal{T}_{n} \mathfrak{g}$.

Define also the twisted Takiff algebra $\mathcal{T}_{n} \mathfrak{g}^{\Gamma}$ :

$$
\begin{aligned}
\mathcal{T}_{n} \mathfrak{g}^{\Gamma} & :=(\mathfrak{g} \otimes \mathbb{C}[t])^{\Gamma} /\left(\mathfrak{g} \otimes t^{n} \mathbb{C}[t]\right)^{\Gamma} \\
& \cong_{\mathbb{C}} \mathfrak{g}^{\sigma} \oplus t \Pi_{1} \mathfrak{g} \cdots \oplus t^{n-1} \Pi_{n-1} \mathfrak{g} .
\end{aligned}
$$

In particular $\mathcal{T}_{1} \mathfrak{g}^{\Gamma}=\mathfrak{g}^{\sigma}$.
We use the notation $X_{\underline{p}}$ for the class of the element $t^{p} X=X \otimes t^{p}$ in $\mathcal{T}_{n} \mathfrak{g}$.
For any $z \in \mathbb{C}$ we have the naive isomorphism $\mathcal{T}_{n} \mathfrak{g} \cong(\mathfrak{g} \otimes \mathbb{C}[[t-z]]) /\left(\mathfrak{g} \otimes(t-z)^{n} \mathbb{C}[[t-z]]\right)$ which sends $X_{\underline{p}}$ to the class of $X \otimes(t-z)^{p}$. By means of this isomorphism, modules over $\mathcal{T}_{n} \mathfrak{g}$ pull back to modules over $\mathfrak{g} \otimes \mathbb{C}[[t-z]]$.
2.7. Universal Cyclotomic Gaudin Algebra. Given any $n_{0}, n_{z_{1}}, \ldots, n_{z_{N}}, n_{\infty} \in \mathbb{Z}_{\geq 1}$ we write $\boldsymbol{n}=\left\{n_{z_{1}}, \ldots, n_{z_{N}}\right\}$. Let $\mathscr{I}_{n_{\infty}, \boldsymbol{n}, n_{0}} \subset U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}^{+}\right)$denote the two-sided ideal in $U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}^{+}\right)$generated by $\left(\mathfrak{g}^{\text {op }} \otimes t^{-n_{\infty}} \mathbb{C}\left[\left[t^{-1}\right]\right]\right)^{\Gamma}, \mathfrak{g} \otimes\left(t-z_{i}\right)^{n_{z_{i}}} \mathbb{C}\left[\left[t-z_{i}\right]\right], i=1, \ldots, N$, and $\left(\mathfrak{g} \otimes t^{n_{0}} \mathbb{C}[[t]]\right)^{\Gamma}$. Define

$$
\begin{equation*}
U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}^{+}\right)_{n_{\infty}, \boldsymbol{n}, n_{0}}:=U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}^{++}\right) / \mathscr{I}_{n_{\infty}, \boldsymbol{n}, n_{0}} . \tag{2.18}
\end{equation*}
$$

These form an inverse system whose inverse limit is $U\left(\hat{\mathfrak{g}}_{\infty, N, 0}^{+}\right)$.
Let us now take the module $\mathcal{M}$ in (2.12) to be a copy of $\left.U\left(\hat{\mathfrak{g}}_{\infty}^{+}+, N, 0\right)\right)_{n_{\infty}, \boldsymbol{n}, n_{0}}$, regarded as a left module over itself. For any $X \in \mathbb{V}_{0}^{k}$ we have a map $\left.X(u): U\left(\hat{\mathfrak{g}}_{\infty}^{+}+N, 0\right)\right)_{n_{\infty}, \boldsymbol{n}, n_{0}} \rightarrow U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}^{+}\right)_{n_{\infty}, \boldsymbol{n}, n_{0}}$ as in (2.17). By construction this can be written in terms of the left action of $U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}^{+}\right)_{n_{\infty}, \boldsymbol{n}, n_{0}}$, which commutes with the right action of $\left.U\left(\hat{\mathfrak{g}}_{\infty}^{+}, N, 0\right)\right)_{n_{\infty}, \boldsymbol{n}, n_{0}}$. So $X(u)$ commutes with the right action of $U\left(\hat{\mathfrak{g}}_{\infty, N, 0}^{++}\right)_{n_{\infty}, n, n_{0}}$. Hence for all $\left.a \in U\left(\widehat{\mathfrak{g}}_{\infty}^{+}+N, 0\right)\right)_{n_{\infty}, \boldsymbol{n}, n_{0}}, X(u) \cdot a=X(u) \cdot(1 a)=(X(u) .1) a$. That is, $X(u)$ acts by left-multiplication by the element $\left.X(u) .1 \in U\left(\hat{\mathfrak{g}}_{\infty}^{+}, N, 0\right)\right)_{n_{\infty}, \boldsymbol{n}, n_{0}}$. Since the latter depends on the choice of $n_{\infty}, \boldsymbol{n}, n_{0}$, we will denote it by $X(u)_{n_{\infty}, \boldsymbol{n}, n_{0}}$. When the choice of $n_{\infty}, \boldsymbol{n}, n_{0}$ is clear from the context we will write $X(u)_{n_{\infty}, \boldsymbol{n}, n_{0}}$ simply as $X(u)$. By construction, whenever
$n_{0}^{\prime}>n_{0}, n_{z_{i}}^{\prime}>n_{z_{i}}$ and $n_{\infty}^{\prime}>n_{\infty}$ then

$$
\begin{equation*}
X(u)_{n_{\infty}, \boldsymbol{n}, n_{0}}=X(u)_{n_{0}^{\prime}, \boldsymbol{n}^{\prime}, n_{\infty}^{\prime}}+\mathscr{I}_{n_{\infty}, \boldsymbol{n}, n_{0}} \tag{2.19}
\end{equation*}
$$

In other words, the elements $X(u)_{n_{\infty}, \boldsymbol{n}, n_{0}} \in U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}^{++}\right)_{n_{\infty}, \boldsymbol{n}, n_{0}}$ are compatible with the above inverse system and hence define an element of the inverse limit $U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}^{+}\right)$. By a slight abuse of notation we will also call this element simply $X(u) \in U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}^{++}\right)$.

We have the natural inclusion $\mathfrak{g}^{\sigma} \hookrightarrow \mathfrak{g}$ and hence the "diagonal" embedding

$$
\begin{align*}
& \mathfrak{g}^{\sigma} \longleftrightarrow \mathfrak{g}^{\sigma, \mathrm{op}} \oplus \bigoplus_{i=1}^{N} \mathfrak{g} \oplus \mathfrak{g}^{\sigma} \longleftrightarrow \widehat{\mathfrak{g}}_{\infty, N, 0} \\
& X \longmapsto(-X ; X, \ldots, X ; X) \longmapsto\left(-X[0]_{\infty} ; X[0]_{z_{1}}, \ldots, X[0]_{z_{N}} ; X[0]_{0}\right) . \tag{2.20}
\end{align*}
$$

Identifying $\mathfrak{g}^{\sigma}$ with its image under this embedding, this gives an action of $\mathfrak{g}^{\sigma}$ on $U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}\right)$ by leftand right-multiplication. In particular, we can define the adjoint action of $\mathfrak{g}^{\sigma}$ on $U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}\right)$. Note that the adjoint action stabilises $U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}^{+}\right)$but the actions by left- and right-multiplication do not, because the zero-modes at $\infty$ are not present in $U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}^{+}\right)$. Let us write

$$
\begin{equation*}
U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}^{++}\right)^{\mathfrak{g}^{\sigma}}:=\left\{x \in U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}^{++}\right):[a, x]=0 \text { for all } a \in \mathfrak{g}^{\sigma}\right\} \tag{2.21}
\end{equation*}
$$

for the invariant subspace of the adjoint action of $\mathfrak{g}^{\sigma}$ on $U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}^{+}\right)$. Define $U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}^{++}\right)_{n_{\infty}, \boldsymbol{n}, n_{0}}^{\mathfrak{g}^{\sigma}}$ likewise.
Now suppose $X$ is a singular vector, $X \in \mathfrak{z}(\widehat{\mathfrak{g}}) \subset \mathbb{V}_{0}^{k}$. Then $X$ is in particular $\mathfrak{g}^{\sigma}$-invariant, $a . X=0$ for all $a \in \mathfrak{g}^{\sigma}$. Hence we have

$$
0=[1 \otimes a . X]=-[a .1 \otimes X]=-\left[X(u) a .1 \otimes v_{0}\right]=-[X(u) a .1]
$$

where we "swapped using the constant rational function $a$ ", i.e. used $[a .(1 \otimes X)]=0$, in the second equality and used the definition of $X(u)$ in the third. On the other hand, in the space of coinvariants $\mathbb{M} / \mathfrak{g}_{\infty, \boldsymbol{z}, 0}^{\Gamma}$ we have

$$
0=[a \cdot(X(u) \cdot 1)]=[a X(u) \cdot 1] .
$$

Taking the difference of the two equalities above, we get

$$
0=[[a, X(u)] \cdot 1]=[a, X(u)]
$$

where in the last equality we can use the identification $\mathbb{M} / \mathfrak{g}_{\infty, \boldsymbol{z}, 0}^{\Gamma} \cong \mathcal{M}$. (The point is that neither $a X(u)$ nor $X(u) a$ need belong to $\mathcal{M}=U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}^{+}\right)_{n_{\infty}, \boldsymbol{n}, n_{0}}$, but the commutator $[a, X(u)]$ does, as we noted above.) This shows that if $X$ is singular then $X(u) \in U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}^{++}\right)_{n_{\infty}, \boldsymbol{n}, n_{0}}^{\mathfrak{g}^{\sigma}}$.

For each $X \in \mathbb{V}_{0}^{k}$, the element $X(u)$ depends rationally on $u$ with poles at most at $0, \omega^{k} z_{i}, \infty$, $i=1, \ldots, N$. Define the algebra $\mathscr{Z}_{\infty, \boldsymbol{z}, 0}^{n_{\infty}, n_{0}}(\mathfrak{g}, \sigma)^{\Gamma}$ to be the span, in $U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}^{+}\right)_{n_{\infty}, \boldsymbol{n}, n_{0}}^{\mathfrak{g}^{\sigma}}$, of all the coefficients of singular terms of Laurent expansions of the elements $Z(u)$ as $Z$ varies in the space of singular vectors $\mathfrak{z}(\widehat{\mathfrak{g}}) \subset \mathbb{V}_{0}^{-h^{\vee}}$.

By virtue of (2.19), the algebras $\mathscr{Z}_{\infty, \boldsymbol{z}, 0}^{n_{\infty}, \boldsymbol{n}, n_{0}}(\mathfrak{g}, \sigma)^{\Gamma}$ form an inverse system. Define the universal cyclotomic Gaudin algebra $\mathscr{Z}_{\infty, \boldsymbol{z}, 0}(\mathfrak{g}, \sigma)^{\Gamma}$ to be the inverse limit,

$$
\mathscr{Z}_{\infty, \boldsymbol{z}, 0}(\mathfrak{g}, \sigma)^{\Gamma}:=\lim _{\longleftarrow} \mathscr{Z}_{\infty, \boldsymbol{z}, 0}^{n_{\infty}, \boldsymbol{n}, n_{0}}(\mathfrak{g}, \sigma)^{\Gamma}
$$

By the argument in [VY16a], following [FFR94], we have

Theorem 2.4. Each $\mathscr{Z}_{\infty, \boldsymbol{z}, 0}^{n_{\infty}, \boldsymbol{n}, n_{0}}(\mathfrak{g}, \sigma)^{\Gamma}$ is a commutative subalgebra of $U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}^{+}\right)_{n_{\infty}, \boldsymbol{n}, n_{0}}^{\mathfrak{g}_{0}^{\sigma}}$. Hence $\mathscr{Z}_{\infty, z, 0}(\mathfrak{g}, \sigma)^{\Gamma}$ is a commutative subalgebra of $U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}^{+}\right)^{\mathfrak{g}^{\sigma}}$.
2.8. Quadratic cyclotomic Hamiltonians. Let $I_{a} \in \mathfrak{g}$ and $I^{a} \in \mathfrak{g}, a=1, \ldots$, dim $\mathfrak{g}$, be dual bases of $\mathfrak{g}$ with respect to $\langle\cdot, \cdot\rangle$, i.e. $\left\langle I_{a}, I^{b}\right\rangle=\delta_{a}^{b}$. Let $\mathcal{C}:=\frac{1}{2} I^{a} I_{a} \in Z(U(\mathfrak{g}))$, the quadratic Casimir of $\mathfrak{g}$. Here and below we employ summation convention on the index $a=1, \ldots, \operatorname{dim} \mathfrak{g}$. Define an element $F \in \mathfrak{g}^{\sigma}$ and number $K \in \mathbb{C}$ by

$$
\begin{equation*}
F:=\frac{1}{2} \sum_{p=1}^{T-1} \frac{\omega^{p}\left[\sigma^{p} I^{a}, I_{a}\right]}{\omega^{p}-1}, \quad K:=\frac{1}{2} \sum_{p=1}^{T-1} \frac{\omega^{p}\left\langle\sigma^{p} I^{a}, I_{a}\right\rangle k}{\left(\omega^{p}-1\right)^{2}} . \tag{2.22}
\end{equation*}
$$

The quadratic Segal-Sugawara vector $S$ is by definition

$$
\begin{equation*}
S:=\frac{1}{2} I^{a}[-1] I_{a}[-1] v_{0} \in \mathbb{V}_{0}^{k} \tag{2.23}
\end{equation*}
$$

At the critical level $k=-h^{\vee}$, the vector $S$ is singular.

Proposition 2.5. The corresponding element $S(u) \in U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}^{+}\right)^{\mathfrak{g}^{\sigma}}$ is given by

$$
\begin{equation*}
S(u)-\frac{1}{u^{2}} K=\sum_{i=1}^{N} \sum_{k=0}^{T-1} \sum_{p=0}^{\infty} \frac{\omega^{-k p+k} \mathcal{H}_{i, p}}{\left(u-\omega^{-k} z_{i}\right)^{p+1}}+\sum_{\substack{p=0 \\ p \equiv 1 \bmod T}}^{\infty} \frac{\mathcal{H}_{0, p}}{u^{p+1}}+\sum_{\substack{p=0 \\ p \equiv-2 \bmod T}}^{\infty} u^{p} \mathcal{H}_{\infty, p} \tag{2.24}
\end{equation*}
$$

where

$$
\begin{align*}
& \mathcal{H}_{i, p}= \sum_{\substack{j=1 \\
j \neq i}}^{N} \sum_{l=0}^{T-1} \sum_{n, m=0}^{\infty} \frac{(-1)^{n}\left[\begin{array}{c}
n+m \\
m
\end{array}\right]}{\left(z_{i}-\omega^{-l} z_{j}\right)^{n+m+1}} I_{a}[n+p]_{z_{i}} \omega^{-l m}\left(\sigma^{l} I^{a}\right)[m]_{z_{j}} \\
&+\sum_{l=1}^{T-1} \sum_{r, m=0}^{\infty} \frac{\omega^{-l m}(-1)^{r}\left[{ }_{m}^{r+m}{ }_{m}\right]}{\left(\left(1-\omega^{-l}\right) z_{i}\right)^{r+m+1}} \frac{1}{2}\left\{I_{a}[r+p]_{z_{i}},\left(\sigma^{l} I^{a}\right)[m]_{z_{i}}\right\}+\sum_{n=0}^{p-1} \frac{1}{2} I_{a}[n]_{z_{i}} I^{a}[p-n-1]_{z_{i}} \\
& \quad+T \sum_{n, m=0}^{\infty} \frac{(-1)^{n}\left[\begin{array}{c}
n+m \\
m
\end{array}\right]}{z_{i}^{n+m+1}} I_{a}[n+p]_{z_{i}}\left(\Pi_{m} I^{a}\right)[m]_{0}+\sum_{n=0}^{\infty} \frac{(-1)^{n}}{z_{i}^{n+1}} F[n+p]_{z_{i}} \\
& \quad+T \sum_{n, m=0}^{\infty} z_{i}^{n}\left[\begin{array}{c}
n+m \\
m
\end{array}\right]\left(\Pi_{-n-m-1} I_{a}\right)[-n-m-1]_{\infty} I^{a}[p+m]_{z_{i}}, \tag{2.25}
\end{align*}
$$

for $i=1, \ldots, N$, and

$$
\begin{align*}
& \mathcal{H}_{0, p}=T^{2} \sum_{m=0}^{\infty} \sum_{n=0}^{\infty} \sum_{i=1}^{N} \frac{(-1)^{n+1}\left[\begin{array}{l}
n+m \\
n
\end{array}\right]}{z_{i}^{n+m+1}}\left(\Pi_{-m-1} I_{a}\right)[n]_{z_{i}}\left(\Pi_{m+p} I^{a}\right)[m+p]_{0} \\
&+T^{2} \sum_{n=p}^{\infty}\left(\Pi_{n} I_{a}\right)[n]_{0}\left(\Pi_{-n+p-1} I^{a}\right)[-n+p-1]_{\infty} \\
&+\frac{T^{2}}{2} \sum_{n=0}^{p-1}\left(\Pi_{n} I_{a}\right)[n]_{0}\left(\Pi_{p-n-1} I^{a}\right)[p-n-1]_{0}+T\left(\Pi_{p-1} F\right)[p-1]_{0}, \tag{2.26}
\end{align*}
$$

and where finally

$$
\begin{align*}
\mathcal{H}_{\infty, p}= & T^{2} \sum_{i=1}^{N} \sum_{n=0}^{\infty} \sum_{m=p+n+1}^{\infty} z_{i}^{m-n-1-p}\left[\begin{array}{c}
m-1-p \\
n
\end{array}\right]\left(\Pi_{-m-1} I_{a}\right)[-m-1]_{\infty}\left(\Pi_{m-p-1} I^{a}\right)[n]_{z_{i}} \\
& +T^{2} \sum_{n=0}^{\infty}\left(\Pi_{n} I_{a}\right)[n]_{0}\left(\Pi_{-p-n-2} I^{a}\right)[-p-n-2]_{\infty} \\
+ & \frac{T^{2}}{2} \sum_{n=0}^{p}\left(\Pi_{-n-1} I_{a}\right)[-n-1]_{\infty}\left(\Pi_{-p+n-1} I^{a}\right)[-p+n-1]_{\infty}+T\left(\Pi_{-p-2} F\right)[-p-2]_{\infty} \tag{2.27}
\end{align*}
$$

Proof. The proof is given in Appendix A.

Remark 2.6. The expression (2.26) for $\mathcal{H}_{0, p}$ vanishes unless $p \equiv 1 \bmod T$, in accordance with the restriction in the second sum on the right hand side of (2.24). Indeed, the first three terms involve expressions of the form $\left(\Pi_{r} I_{a}\right)[n]\left(\Pi_{s} I^{a}\right)[m]$ for some $r, s \in \mathbb{Z}_{T}$ and $m, n \in \mathbb{Z}_{\geq 0}$. However, using the $\sigma$-invariance of the bilinear pairing on $\mathfrak{g}$ it follows that this is equal to $\omega^{r+s}\left(\Pi_{r} I_{a}\right)[n]\left(\Pi_{s} I^{a}\right)[m]$ and therefore vanishes unless $r+s \equiv 0 \bmod T$. Likewise, for the last term, $\left(\Pi_{p-1} F\right)[p-1]$ vanishes unless $p=1$ since $F \in \mathfrak{g}^{\sigma}$. The expression (2.27) for $\mathcal{H}_{\infty, p}$ vanishes unless $p \equiv-2 \bmod T$ for similar reasons, in accordance with the restriction in the third sum on the right hand side of (2.24).
2.9. Regular singularities and shift-of-argument. In the special case when $n_{z_{i}}=1$ for $i=$ $1, \ldots, N$ and $n_{0}=1$, we obtain commutative subalgebras

$$
\mathscr{Z}_{\infty, \boldsymbol{z}, 0}^{n_{\infty},(1), 1}(\mathfrak{g}, \sigma)^{\Gamma} \subset\left(U\left(\left(\mathfrak{g}^{\mathrm{op}} \otimes t^{-1} \mathbb{C}\left[t^{-1}\right]\right)^{\Gamma} /\left(\mathfrak{g}^{\mathrm{op}} \otimes t^{-n_{\infty}} \mathbb{C}\left[t^{-1}\right]\right)^{\Gamma}\right) \otimes U(\mathfrak{g})^{\otimes N} \otimes U\left(\mathfrak{g}^{\sigma}\right)\right)^{\mathfrak{g}^{\sigma}}
$$

If furthermore we set $n_{\infty}=1$ then we obtain the commutative subalgebra

$$
\mathscr{Z}_{\infty, \boldsymbol{z}, 0}^{1,(1), 1}(\mathfrak{g}, \sigma)^{\Gamma} \subset\left(U(\mathfrak{g})^{\otimes N} \otimes U\left(\mathfrak{g}^{\sigma}\right)\right)^{\mathfrak{g}^{\sigma}}
$$

Now consider setting $n_{\infty}=2$. Pick any linear map

$$
\chi: \Pi_{-1} \mathfrak{g} \rightarrow \mathbb{C}
$$

The Lie algebra $\left(\mathfrak{g}^{\text {op }} \otimes t^{-1} \mathbb{C}\left[t^{-1}\right]\right)^{\Gamma} /\left(\mathfrak{g}^{\text {op }} \otimes t^{-2} \mathbb{C}\left[t^{-1}\right]\right)^{\Gamma}$ is commutative, and is canonically isomorphic to $\Pi_{-1} \mathfrak{g}$ as a vector space. We may therefore regard $\chi$ as an algebra homomorphism $\chi: U\left(\left(\mathfrak{g}^{\text {op }} \otimes t^{-1} \mathbb{C}\left[t^{-1}\right]\right)^{\Gamma} /\left(\mathfrak{g}^{\text {op }} \otimes t^{-2} \mathbb{C}\left[t^{-1}\right]\right)^{\Gamma}\right) \rightarrow \mathbb{C}$. Let then

$$
\mathcal{A}_{\infty, \boldsymbol{z}, 0}^{1,(1), 1}(\mathfrak{g}, \sigma, \chi)^{\Gamma}:=\left(\chi \otimes \mathrm{id}^{\otimes N} \otimes \mathrm{id}\right)\left(\mathscr{Z}_{\infty, \boldsymbol{z}, 0}^{2,(1), 1}(\mathfrak{g}, \sigma)^{\Gamma}\right)
$$

This defines a commutative subalgebra

$$
\mathcal{A}_{\infty, \boldsymbol{z}, 0}^{1,(1), 1}(\mathfrak{g}, \sigma, \chi)^{\Gamma} \subset\left(U(\mathfrak{g})^{\otimes N} \otimes U\left(\mathfrak{g}^{\sigma}\right)\right)^{\mathfrak{g}_{\chi}^{\sigma}}
$$

where $\mathfrak{g}_{\chi}^{\sigma}=\left\{X \in \mathfrak{g}^{\sigma}: \chi([X, Y])=0\right.$ for all $\left.Y \in \Pi_{-1} \mathfrak{g}\right\}$ denotes the centralizer of $\chi$ under the coadjoint action. Note that in the special case $\chi=0$ we recover $\mathscr{Z}_{\infty, \boldsymbol{z}, 0}^{1,(1), 1}(\mathfrak{g}, \sigma)^{\Gamma}$.

For any $X \in \mathfrak{g}$ and $i=1, \ldots, N$ we let $X^{(i)}$ denote the element of $U(\mathfrak{g})^{\otimes N} \otimes U\left(\mathfrak{g}^{\sigma}\right)$ with $X$ in the $i^{\text {th }}$ tensor factor and a 1 everywhere else. Similarly, for $X \in \mathfrak{g}^{\sigma}$ we let $X^{(0)}$ be the element
$1^{\otimes N} \otimes X$ of $U(\mathfrak{g})^{\otimes N} \otimes U\left(\mathfrak{g}^{\sigma}\right)$. The only non-zero Hamiltonians of Proposition 2.5 above are then

$$
\begin{gathered}
\mathcal{H}_{i, 0}=\sum_{\substack{j=1 \\
j \neq i}}^{N} \sum_{l=0}^{T-1} \frac{I_{a}^{(i)}\left(\sigma^{l} I^{a}\right)^{(j)}}{\left(z_{i}-\omega^{-l} z_{j}\right)}+\sum_{l=1}^{T-1} \frac{\left(\sigma^{l} I_{a}\right)^{(i)} I^{a(i)}}{\left(1-\omega^{-l}\right) z_{i}}+T \frac{1}{z_{i}} I_{a}^{(i)}\left(\Pi_{0} I^{a}\right)^{(0)}+T I^{a(i)} \chi\left(\Pi_{-1} I_{a}\right), \\
\mathcal{H}_{i, 1}=\frac{1}{2} I_{a}^{(i)} I^{a(i)}, \\
\mathcal{H}_{0,0}=T^{2} \sum_{i=1}^{N} \frac{(-1)}{z_{i}}\left(\Pi_{-1} I_{a}\right)^{(i)}\left(\Pi_{0} I^{a}\right)^{(0)}+T^{2}\left(\Pi_{0} I_{a}\right)^{(0)} \chi\left(\Pi_{-1} I^{a}\right), \\
\mathcal{H}_{0,1}=\frac{T^{2}}{2}\left(\Pi_{0} I_{a}\right)^{(0)}\left(\Pi_{0} I^{a}\right)^{(0)}+T F^{(0)},
\end{gathered}
$$

and

$$
\mathcal{H}_{\infty, 0}=\frac{T^{2}}{2} \chi\left(\Pi_{-1} I_{a}\right) \chi\left(\Pi_{-1} I^{a}\right)
$$

Note that $\mathcal{H}_{0,0}=0$ unless $T=1$ and $\mathcal{H}_{\infty, 0}=0$ unless $T=1$ or 2 , cf. Remark 2.6.
Remark 2.7. The cyclotomic Gaudin algebra introduced in [VY16a] is the commutative subalgebra $\mathscr{Z}_{\infty, \boldsymbol{z}, 0}^{1,(1), 0}(\mathfrak{g}, \sigma)^{\Gamma} \subset\left(U(\mathfrak{g})^{\otimes N}\right)^{\mathfrak{g}^{\sigma}}$. It is the image of $\mathscr{Z}_{\infty, \boldsymbol{z}, 0}^{1,(1), 1}(\mathfrak{g}, \sigma)^{\Gamma}$ under id ${ }^{\otimes N} \otimes \epsilon$, where $\epsilon: U\left(\mathfrak{g}^{\sigma}\right) \rightarrow \mathbb{C}$ is the counit.

Remark 2.8. The algebra $\mathcal{A}_{\infty, \boldsymbol{z}, 0}^{1,(1), 1}(\mathfrak{g}, \sigma, \chi)^{\Gamma}$ is a cyclotomic generalisation of the quantum shift-ofargument subalgebra; see [Ryb06, FFT10, FFRb10].

Remark 2.9. Sometimes setting $\chi \neq 0$ is called adding twisted boundary conditions. The name comes from the Heisenberg XXX spin chain of which the usual Gaudin model is a limit.

## 3. Statement of main Results

3.1. Cartan data and Verma modules. We fix a Cartan decomposition $\mathfrak{g}=\mathfrak{n}-\oplus \mathfrak{h} \oplus \mathfrak{n}$ of $\mathfrak{g}$. Let $\Delta^{+} \subset \mathfrak{h}^{*}$ be the set of positive roots of $\mathfrak{g}$ and $\left\{\alpha_{i}\right\}_{i \in I} \subset \Delta^{+}$the set of simple roots, where $i$ runs over the set $I$ of nodes of the Dynkin diagram of $\mathfrak{g}$. Let $E_{\alpha}$ (resp. $F_{\alpha}$ ) be a root vector of weight $\alpha$ (resp. $-\alpha$ ) for each root $\alpha \in \Delta^{+}$, and $H_{\alpha} \equiv \alpha^{\vee}:=\left[E_{\alpha}, F_{\alpha}\right]$ the corresponding coroot. Overloading notation somewhat, we write $H_{i}:=H_{\alpha_{i}}, i \in I$. Then $\left\{H_{i}\right\}_{i \in I} \cup\left\{E_{\alpha}, F_{\alpha}\right\}_{\alpha \in \Delta+}$ is a Cartan-Weyl basis of $\mathfrak{g}$. We shall assume the Cartan decomposition has been chosen to be compatible with the automorphism $\sigma: \mathfrak{g} \rightarrow \mathfrak{g}$, in the sense that $\sigma(\mathfrak{h})=\mathfrak{h}, \sigma(\mathfrak{n})=\mathfrak{n}$ and $\sigma\left(\mathfrak{n}^{-}\right)=\mathfrak{n}^{-}$. (Such a choice is always possible [Kac83].)

Let $M_{\lambda}$ denote the Verma module over $\mathfrak{g}$ with highest weight $\lambda \in \mathfrak{h}^{*}$, namely

$$
\begin{equation*}
M_{\lambda}:=U(\mathfrak{g}) \otimes_{U(\mathfrak{h} \oplus \mathfrak{n})} \mathbb{C} v_{\lambda} \tag{3.1}
\end{equation*}
$$

where $\mathbb{C} v_{\lambda}$ is the one-dimensional module over $\mathfrak{h} \oplus \mathfrak{n}$ generated by a vector $v_{\lambda}$ with $\mathfrak{n} . v_{\lambda}=0$ and $h . v_{\lambda}=\lambda(h) v_{\lambda}$ for all $h \in \mathfrak{h}$. Similarly, let $M_{\lambda}^{\sigma}$ denote the Verma module over $\mathfrak{g}^{\sigma}$ with highest weight $\lambda \in \mathfrak{h}^{*, \sigma}$,

$$
\begin{equation*}
M_{\lambda}^{\sigma}:=U\left(\mathfrak{g}^{\sigma}\right) \otimes_{U\left(\mathfrak{h}^{\sigma} \oplus \mathfrak{n}^{\sigma}\right)} \mathbb{C} v_{\lambda} \tag{3.2}
\end{equation*}
$$

3.2. The weight function. Let $\lambda_{1}, \ldots, \lambda_{N} \in \mathfrak{h}^{*}$ be $\mathfrak{g}$-weights. Let $\lambda_{0} \in \mathfrak{h}^{*, \sigma}$ be a $\mathfrak{g}^{\sigma}$-weight.

As above, let $z_{1}, \ldots, z_{N}$ be a collection of $N \in \mathbb{Z}_{\geq 0}$ non-zero points in $\mathbb{C}$ such that $\Gamma z_{i} \cap \Gamma z_{j}=\emptyset$ for all $1 \leq i<j \leq N$. In addition, let $w_{1}, \ldots, w_{m}$ be a collection of $m \in \mathbb{Z}_{\geq 0}$ non-zero points in $\mathbb{C}$ such that $\Gamma w_{i} \cap \Gamma w_{j}=\emptyset$ for all $1 \leq i<j \leq m$ and such that $\Gamma w_{i} \cap \Gamma z_{j}=\emptyset$ for all $1 \leq i \leq m$ and $1 \leq j \leq N$. Let $c(1), \ldots, c(m)$ be elements of $I$. We call $c(i)$ the colour of the variable $w_{i}$.

We now define the (cyclotomic) weight function $\psi$ associated to these data.
Recall the projectors $\Pi_{k}$ from (2.2), and in particular $\Pi_{0}$. Define linear maps

$$
\theta_{s}: \bigotimes_{i=1}^{N} M_{\lambda_{i}} \otimes M_{\lambda_{0}}^{\sigma} \otimes \underbrace{\mathfrak{n}_{-} \otimes \ldots \otimes \mathfrak{n}_{-}}_{s} \longrightarrow \bigotimes_{i=1}^{N} M_{\lambda_{i}} \otimes M_{\lambda_{0}}^{\sigma} \otimes \underbrace{\mathfrak{n}_{-} \otimes \ldots \otimes \mathfrak{n}_{-}}_{s-1}
$$

for $s=1,2, \ldots, m$, by

$$
\begin{aligned}
& \theta_{s}\left(x_{1} \otimes \ldots \otimes x_{N} \otimes x_{0} \otimes y_{1} \otimes \ldots \otimes y_{s}\right) \\
& =\frac{\left.x_{1} \otimes \ldots \otimes \ldots \otimes x_{N} \otimes\left(T \Pi_{0} y_{s}\right) x_{0} \otimes y_{1} \otimes \ldots \otimes y_{s-1}\right)}{w_{s}} \\
& +\sum_{i=1}^{N} \sum_{j \in \mathbb{Z}_{T}} \frac{\left.x_{1} \otimes \ldots \otimes x_{i-1} \otimes\left(\sigma^{j} y_{s}\right) x_{i} \otimes x_{i+1} \otimes \ldots \otimes x_{N} \otimes x_{0} \otimes y_{1} \otimes \ldots \otimes y_{s-1}\right)}{w_{s}-\omega^{-j} z_{i}} \\
& \quad+\sum_{i=1}^{s-1} \sum_{j \in \mathbb{Z}_{T}} \frac{x_{1} \otimes \ldots \otimes x_{N} \otimes x_{0} \otimes y_{1} \otimes \ldots \otimes y_{i-1} \otimes\left[\sigma^{j} y_{s}, y_{i}\right] \otimes y_{i+1} \otimes \ldots \otimes y_{s-1}}{w_{s}-\omega^{-j} w_{i}}
\end{aligned}
$$

Then the weight function $\psi$ is by definition the element

$$
\begin{equation*}
\psi:=(-1)^{m}\left(\theta_{1} \circ \cdots \circ \theta_{m}\right)\left(v_{\lambda_{1}} \otimes \ldots \otimes v_{\lambda_{N}} \otimes v_{\lambda_{0}} \otimes F_{\alpha_{c(1)}} \otimes F_{\alpha_{c(2)}} \otimes \ldots \otimes F_{\alpha_{c(m)}}\right) \tag{3.3}
\end{equation*}
$$

For $\lambda \in \mathfrak{h}^{*}$ define $\Pi_{0} \lambda \in \mathfrak{h}^{*, \sigma}$ by $\left(\Pi_{0} \lambda\right)\left(\Pi_{0} h\right):=\lambda\left(\Pi_{0} h\right)$ for $h \in \mathfrak{h}$.
Lemma 3.1. The weight function $\psi$ is an element of $\bigotimes_{i=1}^{N} M_{\lambda_{i}} \otimes M_{\lambda_{0}}^{\sigma}$ of $\mathfrak{g}^{\sigma}$-weight

$$
\begin{equation*}
\lambda_{\infty}:=\lambda_{0}+\sum_{i=1}^{N} \Pi_{0} \lambda_{i}-\sum_{j=1}^{m} \Pi_{0} \alpha_{c(j)} \in \mathfrak{h}^{*, \sigma} . \tag{3.4}
\end{equation*}
$$

Proof. We can regard $\mathfrak{n}_{-}$as a module over $\mathfrak{h}^{\sigma}$, under the adjoint action. Then $\bigotimes_{i=1}^{N} M_{\lambda_{i}} \otimes M_{\lambda_{0}}^{\sigma} \otimes \mathfrak{n}_{-}^{\otimes s}$ is a module over $\mathfrak{h}^{\sigma}$, for every $s \in \mathbb{Z}_{\geq 0}$. Pick any element $\Pi_{0} h \in \mathfrak{h}^{\sigma}$. For each $i \in I$ we have $\left(\Pi_{0} h\right) \cdot F_{\alpha_{i}}:=\left[\Pi_{0} h, F_{\alpha_{i}}\right]=-\alpha_{i}\left(\Pi_{0} h\right) F_{\alpha_{i}}=-\left(\Pi_{0} \alpha_{i}\right)(h) F_{\alpha_{i}}$. That is, $F_{\alpha_{i}}$ has $\mathfrak{g}^{\sigma}$-weight $-\Pi_{0} \alpha_{i}$. Therefore $v_{\lambda_{1}} \otimes \ldots \otimes v_{\lambda_{N}} \otimes v_{\lambda_{0}} \otimes F_{\alpha_{c(1)}} \otimes F_{\alpha_{c(2)}} \otimes \ldots \otimes F_{\alpha_{c(m)}} \in \bigotimes_{i=1}^{N} M_{\lambda_{i}} \otimes M_{\lambda_{0}}^{\sigma} \otimes \mathfrak{n}_{-}^{\otimes m}$ has $\mathfrak{g}^{\sigma}$-weight $\lambda_{\infty}$.

Now, if an element $y \in \mathfrak{n}_{-}$has $\mathfrak{g}^{\sigma}$-weight $\lambda$ then so does $\sigma^{j} y$ for any $j \in \mathbb{Z}_{T}$. Indeed, $\Pi_{0}=\sigma^{j} \Pi_{0}$ so that $\left[\Pi_{0} h, \sigma^{j} y\right]=\left[\sigma^{j} \Pi_{0} h, \sigma^{j} y\right]=\sigma^{j}\left[\Pi_{0} h, y\right]=\sigma^{j} y \lambda\left(\Pi_{0} h\right)$. It follows that the maps $\theta_{s}$ commute with the action of $\mathfrak{h}^{\sigma}$, i.e. preserve $\mathfrak{g}^{\sigma}$-weight. Hence the result.
3.3. The weight $\Lambda_{0}$. Define a weight $\Lambda_{0} \in \mathfrak{h}^{*, \sigma}$ by

$$
\begin{equation*}
\Lambda_{0}(h):=\sum_{r=1}^{T-1} \frac{\operatorname{tr}_{\mathfrak{n}}\left(\sigma^{-r} \mathrm{ad}_{h}\right)}{1-\omega^{r}} \tag{3.5}
\end{equation*}
$$

where $\operatorname{ad}_{h}: \mathfrak{n} \rightarrow \mathfrak{n} ; X \mapsto[h, X]$ is the adjoint action of $\mathfrak{h}$ on $\mathfrak{n}$.

For a more explicit expression for $\Lambda_{0}$, note that [Kac83, §8.6]

$$
\begin{equation*}
\sigma\left(E_{\alpha}\right)=\tau_{\alpha} E_{\sigma(\alpha)}, \quad \sigma\left(H_{i}\right)=H_{\sigma(i)}, \quad \sigma\left(F_{\alpha}\right)=\tau_{\alpha}^{-1} F_{\sigma(\alpha)} . \tag{3.6}
\end{equation*}
$$

Here, by overloading notation, we write $\sigma: \Delta^{+} \rightarrow \Delta^{+}$for the symmetry of the root system, coming in turn from a symmetry $\sigma: I \rightarrow I$ of the Dynkin diagram. The numbers $\tau_{\alpha}, \alpha \in \Delta^{+}$, are certain roots of unity in $\Gamma=\omega^{\mathbb{Z}}$. (So the "inner part" of the automorphism $\sigma: \mathfrak{g} \rightarrow \mathfrak{g}$ is encoded in the choice of $\tau_{\alpha_{i}}, i \in I$; the remaining $\tau_{\alpha}$ are fixed by this choice.)

Then

$$
\begin{equation*}
\Lambda_{0}=\sum_{r=1}^{T-1} \frac{1}{1-\omega^{r}} \sum_{\substack{\alpha \in \Delta+\\ \sigma^{r}(\alpha)=\alpha}}\left(\prod_{p=0}^{r-1} \tau_{\sigma^{p}(\alpha)}^{-1}\right) \alpha . \tag{3.7}
\end{equation*}
$$

3.4. The cyclotomic Bethe equations. Given a complex vector space $A$ on which $\sigma$ acts, for any linear map $\eta: A \rightarrow \mathbb{C}$ we define $L_{\sigma} \eta:=\eta \circ \sigma^{-1}$. Let $\chi \in \mathfrak{h}^{*}$ be such that $L_{\sigma} \chi=\omega \chi$. The cyclotomic Bethe equations (with twisted boundaries) are

$$
\begin{align*}
0=\sum_{r=0}^{T-1} \sum_{i=1}^{N} \frac{\left\langle\alpha_{c(j)}, L_{\sigma}^{r} \lambda_{i}\right\rangle}{w_{j}-\omega^{r} z_{i}}- & \sum_{r=0}^{T-1} \sum_{\substack{k=1 \\
k \neq j}}^{m} \frac{\left\langle\alpha_{c(j)}, L_{\sigma}^{r} \alpha_{c(k)}\right\rangle}{w_{j}-\omega^{r} w_{k}} \\
& +\frac{1}{w_{j}}\left(-\frac{1}{2} \sum_{r=1}^{T-1}\left\langle\alpha_{c(j)}, L_{\sigma}^{r} \alpha_{c(j)}\right\rangle+\left\langle\alpha_{c(j)}, T \lambda_{0}+\Lambda_{0}\right\rangle\right)+\left\langle\alpha_{c(j)}, \chi\right\rangle . \tag{3.8}
\end{align*}
$$

3.5. Eigenvectors of the cyclotomic Gaudin algebra. Let $\chi \in \mathfrak{h}^{*}$ be such that $L_{\sigma} \chi=\omega \chi$. Extend $\chi$ to an element of $\mathfrak{g}^{*}$ by setting $\chi(\mathfrak{n})=\chi\left(\mathfrak{n}_{-}\right)=0$. Then a $\chi$ is a linear map $\Pi_{-1} \mathfrak{g} \rightarrow \mathbb{C}$.

Suppose $\left(w_{1}, \ldots, w_{m} ; c(1), \ldots, c(m)\right)$ are such that cyclotomic Bethe equations (3.8) are satisfied.
Theorem 3.2. The weight function is an eigenvector of the algebra $\mathcal{A}_{\infty, z, 0}^{1,(1), 1}(\mathfrak{g}, \sigma, \chi)^{\Gamma}$. In particular, the eigenvalues $E_{i}$ of the quadratic cyclotomic Gaudin Hamiltonians $\mathcal{H}_{i, 0}$ are given by

$$
E_{i}:=\sum_{\substack{j=1 \\ j \neq i}}^{N} \sum_{s=0}^{T-1} \frac{\left\langle\lambda_{i}, L_{\sigma}^{s} \lambda_{j}\right\rangle}{z_{i}-\omega^{s} z_{j}}-\sum_{j=1}^{m} \sum_{s=0}^{T-1} \frac{\left\langle\lambda_{i}, L_{\sigma}^{s} \alpha_{c(j)}\right\rangle}{z_{i}-\omega^{s} w_{j}}+\frac{1}{z_{i}}\left(\left\langle\lambda_{i}, T \lambda_{0}+\Lambda_{0}\right\rangle+\frac{1}{2} \sum_{s=1}^{T-1}\left\langle\lambda_{i}, L_{\sigma}^{s} \lambda_{i}\right\rangle\right)+\left\langle\lambda_{i}, \chi\right\rangle .
$$

Moreover the eigenvalues of $\mathcal{H}_{0,0}, \mathcal{H}_{0,1}$ and $\mathcal{H}_{\infty, 0}$ are, respectively,

$$
\left\langle\chi, T \lambda_{0}+\Lambda_{0}\right\rangle, \quad \frac{1}{2}\left\langle T \lambda_{0}+\Lambda_{0}, T \lambda_{0}+\Lambda_{0}\right\rangle+\left\langle T \lambda_{0}+\Lambda_{0}, \varrho\right\rangle-K, \quad \frac{1}{2}\langle\chi, \chi\rangle,
$$

where $\varrho:=\frac{1}{2} \sum_{\alpha \in \Delta^{+}} \alpha$.
Remark 3.3. The eigenvalue of $\mathcal{H}_{0,0}$ is zero unless $T=1$, on $\mathbb{Z} / T \mathbb{Z}$-grading grounds, as it should be in view of Remark 2.6. Similarly, the eigenvalue of $\mathcal{H}_{\infty, 0}$ vanishes unless $T=1$ or 2 .

In particular, when $\chi=0$ the weight function is an eigenvector of the algebra $\mathscr{Z}_{\infty, z, 0}^{1,(1), 1}(\mathfrak{g}, \sigma)^{\Gamma}$. Moreover, in that case we have the following.

Theorem 3.4. In the special case $\chi=0$, the weight function $\psi$ belongs to $\left(\otimes_{i=1}^{N} M_{\lambda_{i}} \otimes M_{\lambda_{0}}^{\sigma}\right)_{\lambda_{\infty}}^{\mathfrak{n}^{\sigma}}$, the subspace of $\mathfrak{n}^{\sigma}$-singular vectors of $\mathfrak{g}^{\sigma}$-weight $\lambda_{\infty}$ in $\bigotimes_{i=1}^{N} M_{\lambda_{i}} \otimes M_{\lambda_{0}}^{\sigma}$.

Remark 3.5. We have not shown that the weight function is non-zero. When $\sigma$ is a diagram automorphism, this is proved in [VY].

## 4. Proofs

4.1. Restricted duals and contragredient Verma modules. Given a module $M$ over $\mathfrak{g}$ we write $(M)_{\mu}$ for the subspace of weight $\mu \in \mathfrak{h}^{*}$,

$$
\begin{equation*}
(M)_{\mu}:=\left\{v \in M: \text { there exists } n \in \mathbb{Z}_{\geq 1} \text { such that }(h-\mu(h) 1)^{n} \cdot v=0 \text { for all } h \in \mathfrak{h}\right\} . \tag{4.1}
\end{equation*}
$$

The module $M$ is a weight module if $M=\bigoplus_{\mu \in \mathfrak{h}^{*}}(M)_{\mu}$. In this paper we work with weight modules all of whose weight subspaces are of finite dimension. If $M, N$ are two such modules, then by $\operatorname{Hom}_{\mathbb{C}}(M, N)$ we shall always mean the restricted space of linear maps

$$
\operatorname{Hom}_{\mathbb{C}}(M, N):=\bigoplus_{\mu, \nu \in \mathfrak{h}^{*}} \operatorname{Hom}_{\mathbb{C}}\left((M)_{\mu},(N)_{\nu}\right)
$$

In particular we shall write $M^{*}:=\operatorname{Hom}_{\mathbb{C}}(M, \mathbb{C})=\bigoplus_{\mu \in \mathfrak{h}^{*}}\left((M)_{\mu}\right)^{*}$, i.e. our duals are restricted duals. We have $\operatorname{Hom}_{\mathbb{C}}(M, N)=\operatorname{Hom}_{\mathbb{C}}\left(M \otimes N^{*}, \mathbb{C}\right)=\operatorname{Hom}_{\mathbb{C}}\left(\mathbb{C}, M^{*} \otimes N\right)=M^{*} \otimes N$.

The restricted dual $M_{\lambda}^{*}$ of the Verma module $M_{\lambda}$ is naturally a right $U(\mathfrak{g})$-module. We may twist by any anti-automorphism of $U(\mathfrak{g})$ to obtain a left module. The Cartan anti-automorphism $\varphi: U(\mathfrak{g}) \rightarrow U(\mathfrak{g})$ is defined by

$$
\begin{equation*}
\varphi\left(H_{i}\right)=H_{i}, \quad i \in I, \quad \text { and } \quad \varphi\left(E_{\alpha}\right)=F_{\alpha}, \quad \varphi\left(F_{\alpha}\right)=E_{\alpha}, \quad \alpha \in \Delta^{+} . \tag{4.2}
\end{equation*}
$$

It obeys $\varphi^{2}=\mathrm{id}$. The twist of $M_{\lambda}^{*}$ by $\varphi$ is the left $U(\mathfrak{g})$-module called the contragredient Verma module. Henceforth by $M_{\lambda}^{*}$ we shall always mean the restricted dual equipped with this left $U(\mathfrak{g})$ module structure. That is

$$
(x . f)(v):=f(\varphi(x) . v), \quad f \in M_{\lambda}^{*}, \quad x \in \mathfrak{g}, \quad v \in M_{\lambda} .
$$

See e.g. [Hu08, §3.3].
Let $S: U(\mathfrak{g}) \rightarrow U(\mathfrak{g})$ be the antipode map, i.e. the extension of the map $\mathfrak{g} \rightarrow \mathfrak{g} ; X \rightarrow-X$ to an anti-automorphism of $U(\mathfrak{g})$. We have the automorphism $\varphi \circ S=S \circ \varphi$ of $\mathfrak{g}$. Let us write $\left(M_{\lambda}^{*}\right)^{\varphi \circ S}$ for the left $U(\mathfrak{g})$-module obtained by twisting by this automorphism. In other words $\left(M_{\lambda}^{*}\right)^{\varphi \circ S}$ is the dual of $M_{\lambda}$ in the usual Hopf-algebraic sense. Hence we have

$$
\begin{equation*}
\operatorname{Hom}_{\mathfrak{g}}\left(A, B \otimes M_{\lambda}\right)=\operatorname{Hom}_{\mathfrak{g}}\left(A \otimes\left(M_{\lambda}^{*}\right)^{\varphi \circ S}, B\right) . \tag{4.3}
\end{equation*}
$$

Similarly, one has the notion of weight modules, contragredient Verma modules, etc., over $\mathfrak{g}^{\sigma}$.
Given a module $V$ over a Lie algebra $\mathfrak{a}$, we denote by $V^{\mathfrak{a}}$ the space of invariants $V^{\mathfrak{a}}:=\{x \in V$ : $a . x=0$ for all $a \in \mathfrak{a}\}$.
4.2. Heisenberg algebras at the marked points. Let $z_{1}, \ldots, z_{N}, w_{1}, \ldots, w_{m}$, be as in $\S 3$. For brevity we introduce $p:=N+m$ and $\left(x_{1}, \ldots, x_{p}\right):=\left(z_{1}, \ldots, z_{N}, w_{1}, \ldots, w_{m}\right)$. Let $\boldsymbol{x}:=\left\{x_{1}, \ldots, x_{p}\right\}$.

Let $\mathfrak{n}_{\mathbb{C}}\left(\right.$ resp. $\left.\mathfrak{n}_{\mathbb{C}}^{*}\right)$ denote the vector space $\mathfrak{n}$ (resp. $\left.\mathfrak{n}^{*}\right)$ endowed with the structure of a commutative Lie algebra. On the commutative Lie algebra $\mathfrak{n}_{\mathbb{C}} \oplus \mathfrak{n}_{\mathbb{C}}^{*}$ there is a non-degenerate bilinear
skew-symmetric form $\langle\cdot, \cdot\rangle$ defined by

$$
\begin{equation*}
\langle(X, \eta),(Y, \psi)\rangle=\psi(X)-\eta(Y), \tag{4.4}
\end{equation*}
$$

for any $(X, \eta),(Y, \psi) \in \mathfrak{n}_{\mathbb{C}} \oplus \mathfrak{n}_{\mathbb{C}}^{*}$ and an action by automorphisms of the group $\Gamma$ given by

$$
\begin{equation*}
\omega \cdot(X, \eta):=\left(\sigma(X), L_{\sigma}(\eta)\right) . \tag{4.5}
\end{equation*}
$$

Let $H_{x_{i}}, i=1, \ldots, p$, denote the central extension of the commutative Lie algebra $\left(\mathfrak{n}_{\mathbb{C}} \oplus \mathfrak{n}_{\mathbb{C}}^{*}\right) \otimes$ $\mathbb{C}\left(\left(t-x_{i}\right)\right)$, by a one-dimensional centre $\mathbb{C} \mathbf{1}_{x_{i}}$, defined by the cocycle $\operatorname{res}_{t-x_{i}}\langle f, g\rangle \mathbf{1}_{x_{i}}$. Let $H_{0}^{\Gamma}$ denote the extension of $\left(\mathfrak{n}_{\mathbb{C}} \otimes \mathbb{C}((t))\right)^{\Gamma, 0} \oplus\left(\mathfrak{n}_{\mathbb{C}}^{*} \otimes \mathbb{C}((t))\right)^{\Gamma,-1}$ by a one-dimensional centre $\mathbb{C} \mathbf{1}_{0}$ defined by the cocycle $\operatorname{res}_{t}\langle f, g\rangle \mathbf{1}_{0}$. Let $H_{\infty}^{\Gamma}$ denote the extension of $\left(\mathfrak{n}_{\mathbb{C}} \otimes \mathbb{C}\left(\left(t^{-1}\right)\right)\right)^{\Gamma, 0} \oplus\left(\mathfrak{n}_{\mathbb{C}}^{*} \otimes \mathbb{C}\left(\left(t^{-1}\right)\right)\right)^{\Gamma,-1}$ by a one-dimensional centre $\mathbb{C} \mathbf{1}_{\infty}$ defined by the cocycle $\operatorname{res}_{t^{-1}} t^{2}\langle f, g\rangle \mathbf{1}_{\infty}$.

Let us give a more explicit description of these Lie algebras in terms of generators and relations. To do so, we first construct bases of $\mathfrak{n}$ and $\mathfrak{n}^{*}$ adapted to the automorphism $\sigma$. Recall the projectors $\Pi_{k}, k \in \mathbb{Z} / T \mathbb{Z}$, from (2.2). By the adjoint action, $\mathfrak{g}$ is a module over itself. In particular, it is a module over its Lie subalgebra $\mathfrak{g}^{\sigma}=\Pi_{0} \mathfrak{g}$. As a $\mathfrak{g}^{\sigma}$-module, $\mathfrak{g}=\bigoplus_{k \in \mathbb{Z} / T \mathbb{Z}} \Pi_{k} \mathfrak{g}$. Let $\Delta_{k}^{+}$denote the set of $\mathfrak{g}^{\sigma}$-weights of $\Pi_{k} \mathfrak{n}$ and for $\alpha \in \Delta_{k}^{+}$let $\mathfrak{n}_{(k, \alpha)}$ denote the corresponding weight subspace of $\Pi_{k} \mathfrak{n}$. We may pick a basis of $\mathfrak{n}$ consisting of vectors $E_{(k, \alpha)} \in \mathfrak{n}_{(k, \alpha)}$, where $k \in \mathbb{Z} / T \mathbb{Z}, \alpha \in \Delta_{k}^{+}{ }^{2}$

We now have two bases of $\mathfrak{n}$, namely $E_{\alpha}, \alpha \in \Delta^{+}$, and $E_{(k, \alpha)}, k \in \mathbb{Z} / T \mathbb{Z}, \alpha \in \Delta_{k}^{+}$. We write $E_{\alpha}^{*}$, $\alpha \in \Delta^{+}$, and $E_{(k, \alpha)}^{*}, k \in \mathbb{Z} / T \mathbb{Z}, \alpha \in \Delta_{k}^{+}$for their respective dual bases of $\mathfrak{n}^{*}$.

Then $H_{x_{i}}$ has the following explicit set of generators:

$$
a_{\alpha}[n]_{x_{i}}:=E_{\alpha} \otimes\left(t-x_{i}\right)^{n}, \quad a_{\alpha}^{*}[n]_{x_{i}}:=E_{\alpha}^{*} \otimes\left(t-x_{i}\right)^{n-1},
$$

where $\alpha \in \Delta^{+}$and $n \in \mathbb{Z}$; while explicit sets of generators for $H_{0}^{\Gamma}$ and $H_{\infty}^{\Gamma}$ are

$$
\begin{aligned}
& a_{(k, \alpha)}[n T+k]_{0}:=E_{(k, \alpha)} \otimes t^{n T+k} \quad \in\left(\mathfrak{n}_{\mathbb{C}} \otimes \mathbb{C}((t))\right)^{\Gamma, 0}, \\
& a_{(k, \alpha)}^{*}[n T-k]_{0}:=E_{(k, \alpha)}^{*} \otimes t^{n T-k-1} \in\left(\mathfrak{n}_{\mathbb{C}}^{*} \otimes \mathbb{C}((t))\right)^{\Gamma,-1},
\end{aligned}
$$

and

$$
\begin{aligned}
& a_{(k, \alpha)}[n T+k]_{\infty}:=E_{(k, \alpha)} \otimes t^{n T+k} \quad \in\left(\mathfrak{n}_{\mathbb{C}} \otimes \mathbb{C}\left(\left(t^{-1}\right)\right)\right)^{\Gamma, 0}, \\
& a_{(k, \alpha)}^{*}[n T-k]_{\infty}:=E_{(k, \alpha)}^{*} \otimes t^{n T-k-1} \in\left(\mathfrak{n}_{\mathbb{C}}^{*} \otimes \mathbb{C}\left(\left(t^{-1}\right)\right)\right)^{\Gamma,-1},
\end{aligned}
$$

respectively, where $k \in \mathbb{Z} / T \mathbb{Z}, \alpha \in \Delta_{k}^{+}$, and $n \in \mathbb{Z}$.

Remark 4.1. The above notation is in accordance with the notations for the modes $a[n], n \in \mathbb{Z}$ of an element $a \in \mathscr{L}$ in a vertex Lie algebra $\mathscr{L}$ used in [VY16b] and recalled in appendix B below. In particular, the shift by 1 in the power of $t$ for the modes of $a_{\alpha}^{*}$ has to do with the fact that this is an element of degree 0 in the vertex Lie algebra associated to the Heisenberg algebra, cf. [VY16b, Example 2.6]. Note also our conventions for the modes at $\infty$.

[^1]These generators obey the relations:

$$
\begin{aligned}
{\left[a_{\alpha}[n]_{x_{i}}, a_{\beta}^{*}[m]_{x_{j}}\right] } & =\delta_{i j} \delta_{\alpha \beta} \delta_{n,-m} \mathbf{1}_{x_{i}}, \\
{\left[a_{(i, \alpha)}[n]_{0}, a_{(j, \beta)}^{*}[m]_{0}\right] } & =\delta_{i j} \delta_{\alpha \beta} \delta_{n,-m} \mathbf{1}_{0}, \\
{\left[a_{(i, \alpha)}[n]_{\infty}, a_{(j, \beta)}^{*}[m]_{\infty}\right] } & =\delta_{i j} \delta_{\alpha \beta} \delta_{n,-m} \mathbf{1}_{\infty},
\end{aligned}
$$

with all other commutators vanishing.
Each $H_{x_{i}}$ is isomorphic to the Heisenberg Lie algebra $H(\mathfrak{g})$, while $H_{0}^{\Gamma}$ and $H_{\infty}^{\Gamma}$ are isomorphic to a subalgebra $H(\mathfrak{g})^{\Gamma}$. Note that the opposite Lie algebra $H_{\infty}^{\Gamma, o p}$ differs from $H_{\infty}^{\Gamma}$ only in the sign of the central extension.

Let $H_{\infty, p, 0}$ denote the extension of

$$
\begin{align*}
\left(\mathfrak{n}_{\mathbb{C}} \otimes \mathbb{C}\left(\left(t^{-1}\right)\right)\right)^{\Gamma, 0} \oplus\left(\mathfrak{n}_{\mathbb{C}}^{*} \otimes \mathbb{C}\left(\left(t^{-1}\right)\right)\right)^{\Gamma,-1} & \\
& \oplus \bigoplus_{i=1}^{p}\left(\mathfrak{n}_{\mathbb{C}} \oplus \mathfrak{n}_{\mathbb{C}}^{*}\right) \otimes \mathbb{C}\left(\left(t-x_{i}\right)\right) \\
& \oplus\left(\mathfrak{n}_{\mathbb{C}} \otimes \mathbb{C}((t))\right)^{\Gamma, 0} \oplus\left(\mathfrak{n}_{\mathbb{C}}^{*} \otimes \mathbb{C}((t))\right)^{\Gamma,-1} \tag{4.6}
\end{align*}
$$

by a one-dimensional centre $\mathbb{C} \mathbf{1}$, defined by the cocycle

$$
\begin{equation*}
\Omega(f, g):=\left(\frac{1}{T} \operatorname{res}_{t}\left\langle f_{0}, g_{0}\right\rangle+\sum_{i=1}^{p} \operatorname{res}_{t-x_{i}}\left\langle f_{x_{i}}, g_{x_{i}}\right\rangle-\frac{1}{T} \operatorname{res}_{t^{-1}} t^{2}\left\langle f_{\infty}, g_{\infty}\right\rangle\right) \mathbf{1} \tag{4.7}
\end{equation*}
$$

where $f=\left(f_{\infty} ; f_{x_{1}}, \ldots, f_{x_{p}} ; f_{0}\right)$ and $g=\left(g_{\infty} ; g_{x_{1}}, \ldots, g_{x_{p}} ; g_{0}\right)$ are elements of the Lie algebra (4.6). Equivalently, $H_{\infty, p, 0}$ is the quotient of the direct sum $H_{\infty}^{\Gamma, 0 \mathbf{p}} \oplus \bigoplus_{i=1}^{p} H_{x_{i}} \oplus H_{0}^{\Gamma}$ by the ideal generated by $\mathbf{1}_{x_{i}}-T \mathbf{1}_{0}, i=1, \ldots, p$, and $\mathbf{1}_{\infty}-\mathbf{1}_{0}$. Then $\mathbf{1}:=\mathbf{1}_{x_{i}}$.

Define also

$$
\mathfrak{h}_{\infty, p, 0}:=\left(\mathfrak{h} \otimes \mathbb{C}\left(\left(t^{-1}\right)\right)\right)^{\Gamma, 0} \oplus \bigoplus_{i=1}^{p} \mathfrak{h} \otimes \mathbb{C}\left(\left(t-x_{i}\right)\right) \oplus(\mathfrak{h} \otimes \mathbb{C}((t)))^{\Gamma, 0}
$$

Let us give a set of explicit generators for this commutative Lie algebra $\mathfrak{h}_{\infty, p, 0}$. Let $H_{(k, a)}, a=$ $1, \ldots, \operatorname{dim}\left(\Pi_{k} \mathfrak{h}\right)$, be a basis of $\Pi_{k} \mathfrak{h}$ for each $k \in \mathbb{Z} / T \mathbb{Z}$. Then $\mathfrak{h}_{0, p, \infty}$ is the commutative Lie algebra with basis

$$
\begin{align*}
b_{j}[n]_{x_{i}} & :=H_{j} \otimes\left(t-x_{i}\right)^{n} \in \mathfrak{h} \otimes \mathbb{C}\left(\left(t-x_{i}\right)\right), \\
b_{(k, a)}[n T+k]_{0} & :=H_{(k, a)} \otimes t^{n T+k} \in(\mathfrak{h} \otimes \mathbb{C}((t)))^{\Gamma, 0},  \tag{4.8}\\
b_{(k, a)}[n T+k]_{\infty} & :=H_{(k, a)} \otimes t^{n T+k} \in\left(\mathfrak{h} \otimes \mathbb{C}\left(\left(t^{-1}\right)\right)\right)^{\Gamma, 0},
\end{align*}
$$

for $j \in I, k \in \mathbb{Z} / T \mathbb{Z}, a=1, \ldots, \operatorname{dim}\left(\Pi_{k} \mathfrak{h}\right)$ and $n \in \mathbb{Z}$.
4.3. Wakimoto modules at the marked points. For each $i=1, \ldots, p$, let $\mathbb{C}\left\rangle_{x_{i}}\right.$ denote the one-dimensional left module over $U\left(\left(\mathfrak{n}_{\mathbb{C}} \oplus \mathfrak{n}_{\mathbb{C}}^{*}\right) \otimes \mathbb{C}\left[\left[t-x_{i}\right]\right] \oplus \mathbb{C} \mathbf{1}_{x_{i}}\right)$ on which $\mathbf{1}_{x_{i}}$ acts as 1 and $\left(\mathfrak{n}_{\mathbb{C}} \oplus \mathfrak{n}_{\mathbb{C}}^{*}\right) \otimes \mathbb{C}[[t]]$ acts as zero. Define $\mathrm{M}_{x_{i}}$ to be the induced module over $H_{x_{i}}$,

$$
\mathbf{M}_{x_{i}}:=U\left(H_{x_{i}}\right) \otimes_{\left.U\left(\left(\mathbf{n}_{\mathbb{C}} \oplus \mathfrak{n}_{\mathbb{C}}^{*}\right) \otimes \mathbb{C}\left[t-x_{i}\right]\right] \oplus \mathbb{C} 1_{x_{i}}\right)} \mathbb{C}| \rangle_{x_{i}} .
$$

Suppose we are given an $\mathfrak{h}^{*}$-valued Laurent series $\nu_{i} \in \mathfrak{h}^{*} \otimes \mathbb{C}\left(\left(t-x_{i}\right)\right)$ in the local coordinate $t-x_{i}$ about the point $x_{i}$. Let $\mathbb{C} v_{\nu_{i}}$ denotes the one-dimensional module over $\mathfrak{h} \otimes \mathbb{C}\left(\left(t-x_{i}\right)\right)$ on which $f . v_{\nu_{i}}=v_{\nu_{i}} \operatorname{res}_{t-x_{i}} \nu_{i}(f)$, for any $f \in \mathfrak{h} \otimes \mathbb{C}\left(\left(t-x_{i}\right)\right)$. Then the Wakimoto module $W_{\nu_{i}}$ is the module over $H_{x_{i}} \oplus \mathfrak{h} \otimes \mathbb{C}\left(\left(t-x_{i}\right)\right)$ given by

$$
W_{\nu_{i}}:=\mathrm{M}_{x_{i}} \otimes \mathbb{C} v_{\nu_{i}}
$$

Explicitly, $W_{\nu_{i}}$ is the Fock module generated by a vacuum vector $\left\rangle_{x_{i}}\right.$ such that $\left.\left.\mathbf{1}_{x_{i} i}\right|\right\rangle_{x_{i}}=| \rangle_{x_{i}}$,

$$
a_{\alpha}[n]| \rangle_{x_{i}}=0, \quad n \geq 0, \quad a_{\alpha}^{*}[n]| \rangle_{x_{i}}=0, \quad n \geq 1,
$$

and

$$
b_{k}[n]| \rangle_{x_{i}}=| \rangle_{x_{i}} \nu_{i,-n-1}\left(H_{k}\right)
$$

where $\nu_{i}\left(t-x_{i}\right)=: \sum_{s=-S}^{\infty} \nu_{i, s}\left(t-x_{i}\right)^{s}$ for some $S \in \mathbb{Z}$ and $\nu_{i, s} \in \mathfrak{h}^{*}$. Here and in what follows we use the obvious shorthand $a_{\alpha}[n]| \rangle_{x_{i}}$ to denote $a_{\alpha}[n]_{x_{i}}| \rangle_{x_{i}}$, etc.

Similarly, let $\mathbb{C}\left\rangle_{0}\right.$ denote the one-dimensional module over $\left(\mathfrak{n}_{\mathbb{C}} \otimes \mathbb{C}[[t]]\right)^{\Gamma, 0} \oplus\left(\mathfrak{n}_{\mathbb{C}}^{*} \otimes \mathbb{C}[[t]]\right)^{\Gamma,-1} \oplus \mathbb{C} \mathbf{1}_{0}$ on which $\mathbf{1}_{0}$ acts as $\frac{1}{T}$ and the first two summands act as zero. Define $\mathrm{M}_{0}^{\Gamma}$ to be the induced module over $H_{0}^{\Gamma}$ :

$$
\mathrm{M}_{0}^{\Gamma}:=U\left(H_{0}^{\Gamma}\right) \otimes_{\left.\left.U\left(\left(\mathfrak{n}_{\mathbb{C}} \otimes \mathbb{C}[t t]\right]\right)^{\Gamma, 0} \oplus\left(\mathfrak{n}_{\mathbb{C}}^{*} \otimes \mathbb{C}[t t]\right)\right)^{\Gamma,-1} \oplus \mathbb{C} \mathbf{1}_{0}\right)} \mathbb{C}| \rangle_{0} .
$$

Suppose we are given an element $\nu_{0} \in\left(\mathfrak{h}^{*} \otimes \mathbb{C}((t))\right)^{\Gamma,-1}$. That is, $\nu_{0}$ is a $\mathfrak{h}^{*}$-valued Laurent series in $t$ such that $\nu_{0}(\omega t)=\omega^{-1} \sigma \nu_{0}(t)$. Let $\mathbb{C} v_{\nu_{0}}$ denote the one-dimensional module over $(\mathfrak{h} \otimes \mathbb{C}((t)))^{\Gamma, 0}$ given by

$$
\begin{equation*}
f \cdot v_{\nu_{0}}=v_{\nu_{0}} \operatorname{res}_{t} \nu_{0}(f), \tag{4.9}
\end{equation*}
$$

for any $f \in(\mathfrak{h} \otimes \mathbb{C}((t)))^{\Gamma, 0}$. We may then consider the twisted Wakimoto module $W_{\nu_{0}}^{\Gamma}$ defined as the $H_{0}^{\Gamma} \oplus(\mathfrak{h} \otimes \mathbb{C}((t)))^{\Gamma, 0}$-module

$$
W_{\nu_{0}}^{\Gamma}:=\mathrm{M}_{0}^{\Gamma} \otimes \mathbb{C} v_{\nu_{0}} .
$$

Explicitly, it is the Fock module generated by a vacuum vector $\left\rangle_{0}\right.$ such that $\left.\left.\mathbf{1}_{0}\right|\right\rangle_{0}=| \rangle_{0} \frac{1}{T}$,

$$
\begin{equation*}
a_{(k, \alpha)}[n]| \rangle_{0}=0, \quad n \geq 0, \quad a_{(k, \alpha)}^{*}[n]| \rangle_{0}=0, \quad n \geq 1, \tag{4.10}
\end{equation*}
$$

and

$$
b_{(k, a)}[n]| \rangle_{0}=| \rangle_{0} \nu_{0,-n-1}\left(H_{(k, a)}\right)
$$

where $\nu_{0}(t)=: \sum_{s=-S}^{\infty} \nu_{0, s} t^{s}$ for some $S \in \mathbb{Z}$ and $\nu_{0, s} \in \mathfrak{h}^{*}$.
Finally, define

$$
H_{\infty}^{\Gamma,+}:=\left(\mathfrak{n}_{\mathbb{C}} \otimes t^{-1} \mathbb{C}\left[\left[t^{-1}\right]\right]\right)^{\Gamma, 0} \oplus\left(\mathfrak{n}_{\mathbb{C}}^{*} \otimes t^{-1} \mathbb{C}\left[\left[t^{-1}\right]\right]\right)^{\Gamma,-1} \oplus \mathbb{C} \mathbf{1}_{\infty},
$$

and let $\mathrm{M}_{\infty}^{\Gamma, \vee}$ denote the right module over $U\left(H_{\infty}^{\Gamma}\right)$ induced from the trivial one-dimensional right module $\mathbb{C}\langle |$ over $U\left(H_{\infty}^{\Gamma,+}\right)$ on which $\mathbf{1}_{\infty}$ acts as $\frac{1}{T}$ and the first two summands act as zero:

$$
\mathrm{M}_{\infty}^{\Gamma, V}:=\mathbb{C}\langle | \otimes_{U\left(H_{\infty}^{\Gamma,+}\right)} U\left(H_{\infty}^{\Gamma}\right)
$$

Suppose $\nu_{\infty} \in\left(\mathfrak{h}^{*} \otimes \mathbb{C}\left(\left(t^{-1}\right)\right)\right)^{\Gamma,-1}$. That is, $\nu_{\infty}$ is a $\mathfrak{h}^{*}$-valued Laurent series in $t^{-1}$ such that $\nu_{\infty}\left(\omega^{-1} t^{-1}\right)=\omega^{-1} \sigma \nu_{\infty}\left(t^{-1}\right)$. Let $\mathbb{C} v_{\nu_{\infty}}^{*}$ denote the one-dimensional module over $\left(\mathfrak{h} \otimes \mathbb{C}\left(\left(t^{-1}\right)\right)\right)^{\Gamma, 0}$ given by

$$
\begin{equation*}
f \cdot v_{\nu_{\infty}}^{*}=-v_{\nu_{\infty}}^{*} \operatorname{res}_{t^{-1}} t^{2} \nu_{\infty}(f), \tag{4.11}
\end{equation*}
$$

for $f \in\left(\mathfrak{h} \otimes \mathbb{C}\left(\left(t^{-1}\right)\right)\right)^{\Gamma, 0}$. Then we have the right $U\left(H_{\infty}^{\Gamma} \oplus\left(\mathfrak{h} \otimes \mathbb{C}\left(\left(t^{-1}\right)\right)\right)^{\Gamma, 0}\right)$-module

$$
W_{\nu_{\infty}}^{\Gamma, V}:=\mathrm{M}_{\infty}^{\Gamma, \vee} \otimes \mathbb{C} v_{\nu_{\infty}}^{*}
$$

Explicitly, $W_{\nu_{\infty}}^{\Gamma, V}$ is the Fock module generated by a vacuum vector $\langle |$ such that $\langle | \mathbf{1}_{\infty}=\frac{1}{T}\langle |$,

$$
\begin{equation*}
\langle | a_{(k, \alpha)}[n]_{\infty}=0, \quad n<0, \quad\langle | a_{(k, \alpha)}^{*}[n]_{\infty}=0, \quad n \leq 0, \tag{4.12}
\end{equation*}
$$

and

$$
\begin{equation*}
\langle | b_{(k, a)}[n]_{\infty}=-\nu_{\infty,-n-1}\left(H_{(k, a)}\right)\langle | \tag{4.13}
\end{equation*}
$$

where $\nu_{\infty}(t)=: \sum_{s=-\infty}^{S} \nu_{\infty, s} t^{s}$ for some $S \in \mathbb{Z}$ and $\nu_{\infty, s} \in \mathfrak{h}^{*}$.
4.4. Free field realization. The modules $W_{\nu_{i}}$ are smooth. That means, by definition, that for each $v \in W_{\nu_{i}}$

$$
0=a_{\alpha}[n]_{x_{i}} v=a_{\alpha}^{*}[n]_{x_{i}} v=b_{k}[n]_{x_{i}} v \quad \text { for all } \quad n \gg 0,
$$

for all $\alpha \in \Delta^{+}$and $k \in I .^{3}$ Similarly $W_{\nu_{0}}^{\Gamma}$ is smooth. The module $W_{\nu_{\infty}}^{\Gamma, V}$ is co-smooth. By that we mean that for each $v \in W_{\nu_{\infty}}^{\Gamma, V}$,

$$
0=v a_{(k, \alpha)}[n]_{\infty}=v a_{(k, \alpha)}^{*}[n]_{\infty}=v b_{(k, a)}[n]_{\infty} \quad \text { for all } \quad n \ll 0,
$$

for all $k \in \mathbb{Z} / T \mathbb{Z}, \alpha \in \Delta_{k}^{+}$and $a \in\left\{1, \ldots, \operatorname{dim} \Pi_{k} \mathfrak{h}\right\}$.
Let now $\pi_{0} \simeq \mathbb{C}\left[b_{i}[n]\right]_{i \in I ; n \leq-1}$ be the induced representation of $\mathfrak{h} \otimes \mathbb{C}((t))$ in which $b_{i}[n]$ acts as 0 for all $i \in I$ and all $n \in \mathbb{Z}_{\geq 0}$. Let M be the induced module over the Heisenberg Lie algebra with generators $a_{\alpha}[n], a_{\alpha}^{*}[n]$, and $\mathbf{1}$, which we denote $H(\mathfrak{g})$ as in [VY16a, $\left.\S 3.5\right]$. Then

$$
\begin{equation*}
\mathbb{W}_{0}:=\mathrm{M} \otimes \pi_{0}, \tag{4.14}
\end{equation*}
$$

is an induced representation of $H(\mathfrak{g}) \oplus \mathfrak{h} \otimes \mathbb{C}((t))$. Explicitly, $\mathbb{W}_{0}$ is the Fock module generated by a vacuum vector $\rangle$ such that $\mathbf{1}|\rangle=| \rangle$,

$$
a_{\alpha}[n]| \rangle=0, \quad n \geq 0, \quad a_{\alpha}^{*}[n]| \rangle=0, \quad n \geq 1,
$$

and

$$
b_{k}[n]| \rangle=0, \quad n \geq 0 .
$$

Let us recall some facts about the free-field realization of $\widehat{\mathfrak{g}}$.
The $H(\mathfrak{g}) \oplus \mathfrak{h} \otimes \mathbb{C}((t))$-module $\mathbb{W}_{0}$ is endowed with the structure of a vertex algebra, see [FFR94] or e.g. [VY16a, equation (3.24)]. In the notation of [VY16b], see also appendix B, it corresponds to the vertex algebra $\mathbb{V}(\mathscr{M})$ where $\mathscr{M}$ is the vertex Lie algebra underlying the Lie algebra $H(\mathfrak{g}) \oplus \mathfrak{h} \otimes \mathbb{C}((t))$, given for instance in [VY16b, Example 2.6]. Specifically, $\mathscr{M}$ is generated by the finite dimensional vector space $M^{o}=\mathfrak{n}_{\mathbb{C}} \oplus \mathfrak{n}_{\mathbb{C}}^{*} \oplus \mathfrak{h}$ with non-trivial $n^{\text {th }}$-products given by

$$
a_{(0)} b=\langle a, b\rangle \mathbf{1}
$$

for any $a, b \in \mathfrak{n}_{\mathbb{C}} \oplus \mathfrak{n}_{\mathbb{C}}^{*}$, where the skew-symmetric form $\langle\cdot, \cdot\rangle$ on the right hand side was defined in (4.4). As a vertex algebra, $\mathbb{W}_{0}$ is in particular also a vertex Lie algebra. We may therefore consider the associated "big" Lie algebra $L\left(\mathbb{W}_{0}\right)$, cf. §B.1, consisting of all formal modes of states in $\mathbb{W}_{0}$.

[^2]For each of the marked points $x_{i}, i=1, \ldots, p$, there is a "local" copy $\mathrm{L}\left(\mathbb{W}_{0}\right)_{x_{i}}$ of $\mathrm{L}\left(\mathbb{W}_{0}\right)$, which contains $\mathrm{L}(\mathscr{M})_{x_{i}}=H_{x_{i}} \oplus \mathfrak{h} \otimes \mathbb{C}\left(\left(t-x_{i}\right)\right)$ as a subalgebra. Moreover, every smooth module over $H_{x_{i}} \oplus \mathfrak{h} \otimes \mathbb{C}\left(\left(t-x_{i}\right)\right)$ on which $\mathbf{1}_{x_{i}}$ acts as 1 becomes a smooth $\mathrm{L}\left(\mathbb{W}_{0}\right)_{x_{i}}$-module in a canonical way. See e.g. [VY16b, Proposition 5.9].

At the fixed-points 0 and $\infty$ we have the local copies $L\left(\mathbb{W}_{0}\right)_{0}$ and $\mathrm{L}\left(\mathbb{W}_{0}\right)_{\infty}$ of the big Lie algebra $\mathrm{L}\left(\mathbb{W}_{0}\right)$. The automorphism $\sigma$ of the Lie algebra $\mathrm{L}(\mathscr{M})=H(\mathfrak{g}) \oplus \mathfrak{h} \otimes \mathbb{C}((t))$ extends in a unique way to an automorphism of $\mathbb{W}_{0}$ as a vertex algebra. So we get an action by automorphisms of the group $\Gamma$ on both local Lie algebras $\mathrm{L}\left(\mathbb{W}_{0}\right)_{0}$ and $\mathrm{L}\left(\mathbb{W}_{0}\right)_{\infty}$. Let $\mathrm{L}\left(\mathbb{W}_{0}\right)_{0}^{\Gamma}$ and $\mathrm{L}\left(\mathbb{W}_{0}\right)_{\infty}^{\Gamma}$ denote the respective fixed-point subalgebras. They contain the local Lie algebras $\mathrm{L}(\mathscr{M})_{0}^{\Gamma}=H_{0}^{\Gamma} \oplus(\mathfrak{h} \otimes \mathbb{C}((t)))^{\Gamma, 0}$ and $\mathrm{L}(\mathscr{M})_{\infty}^{\Gamma}=H_{\infty}^{\Gamma} \oplus\left(\mathfrak{h} \otimes \mathbb{C}\left(\left(t^{-1}\right)\right)\right)^{\Gamma, 0}$ as subalgebras, respectively, using the notation of appendix B.2.

Every smooth module over $\mathrm{L}(\mathscr{M})_{0}^{\Gamma}$ on which $\mathbf{1}_{0}$ acts as $\frac{1}{T}$ becomes a smooth module over $\mathrm{L}\left(\mathbb{W}_{0}\right)_{0}^{\Gamma}$. See [VY16b, Proposition 5.8]. Likewise, every co-smooth module over $L(\mathscr{M})_{\infty}^{\Gamma}$ on which $\mathbf{1}_{\infty}$ acts as $\frac{1}{T}$ becomes a co-smooth module over $\mathrm{L}\left(\mathbb{W}_{0}\right)_{\infty}^{\Gamma}$. See Proposition B.3.

Now the vacuum Verma module $\mathbb{V}_{0}^{-h^{\vee}}$ also has the structure of a vertex algebra. In the notation of [VY16b] recalled in appendix B, it corresponds to the vertex algebra $\mathbb{V}(\mathscr{L})$ where $\mathscr{L}$ is the vertex Lie algebra underlying the affine Lie algebra $\widehat{\mathfrak{g}}$, generated by $L^{o}=\mathfrak{g}$ with non-trivial $n^{\text {th }}$-products given by (see e.g. [VY16b, Example 2.5])

$$
a_{(0)} b=[a, b], \quad a_{(1)} b=-h^{\vee}\langle a, b\rangle K,
$$

for any $a, b \in \mathfrak{g}$. The normalisation of the $1^{\text {st }}$-product by $-h^{\vee}$ is chosen so as to conform with the convention adopted in [VY16b] that the central element $K \in \mathrm{~L}(\mathscr{L})$ should act as 1 on all modules over $\mathrm{L}(\mathscr{L})$, including $\mathbb{V}(\mathscr{L})$.

The associated "big" Lie algebra $\mathrm{L}\left(\mathbb{V}_{0}^{-h^{\vee}}\right)$ contains $\widehat{\mathfrak{g}}=\mathrm{L}(\mathscr{L})$ as a subalgebra, via the embedding sending $K \mapsto v_{0}(-1)$ and $A[n] \mapsto\left(A[-1] v_{0}\right)(n)$ for any $A \in \mathfrak{g}$ and $n \in \mathbb{Z}$. In this way, every smooth module over $\mathrm{L}\left(\mathbb{V}_{0}^{-h^{\vee}}\right)$ pull back to a smooth module over $\widehat{\mathfrak{g}}$ of level $-h^{\vee}$. In the same way, $\mathrm{L}\left(\mathbb{V}_{0}^{-h^{\vee}}\right)^{\Gamma}$ contains the twisted affine algebra $\hat{\mathfrak{g}}^{\Gamma}$ as a subalgebra, so that smooth (resp. co-smooth right) modules over $L\left(\mathbb{V}_{0}^{-h^{\vee}}\right)^{\Gamma}$ pulls back to smooth (resp. co-smooth right) module over $\hat{\mathfrak{g}}^{\Gamma}$ of level $-h^{\vee} / T$.

The Feigin-Frenkel homomorphism, or free field realisation of $\mathbb{V}_{0}^{-h^{\vee}},[$ FF90] (see also [Fre07]) is a homomorphism of vertex algebras

$$
\begin{equation*}
\rho: \mathbb{V}_{0}^{-h^{\vee}} \longrightarrow \mathbb{W}_{0} \tag{4.15}
\end{equation*}
$$

There is a $\mathbb{Z}$-grading on $\mathbb{W}_{0}$ defined by $\operatorname{deg}\left\rangle=0\right.$ and $\operatorname{deg} a_{\alpha}[n]=\operatorname{deg} a_{\alpha}^{*}[n]=\operatorname{deg} b_{i}[n]=n$. There is a $\mathbb{Z}$-grading on $\mathbb{V}_{0}^{-h^{\vee}}$ defined by $\operatorname{deg} v_{0}=0$ and $\operatorname{deg} X[n]=n$ for $X \in \mathfrak{g}$. The homomorphism $\rho$ respects these $\mathbb{Z}$-gradations (see e.g. [Fre07, $\S 6.2 .4]$ ).

The homomorphism $\rho$ induces a homomorphism of big Lie algebras $\mathrm{L}\left(\mathbb{V}_{0}^{-h^{\vee}}\right)_{x_{i}} \rightarrow \mathrm{~L}\left(\mathbb{W}_{0}\right)_{x_{i}}$ for every $i=1, \ldots, p$. Moreover, by the equivariance of the homomorphism with respect to the action of $\sigma$ on both vertex algebras [Sz02], it also induces homomorphisms of $\Gamma$-invariant subalgebras $\mathrm{L}\left(\mathbb{V}_{0}^{-h^{\vee}}\right)_{0}^{\Gamma} \rightarrow \mathrm{L}\left(\mathbb{W}_{0}\right)_{0}^{\Gamma}$ and $\mathrm{L}\left(\mathbb{V}_{0}^{-h^{\vee}}\right)_{\infty}^{\Gamma} \rightarrow \mathrm{L}\left(\mathbb{W}_{0}\right)_{\infty}^{\Gamma}$.

By means of this homomorphism, every smooth module over $\mathrm{L}(\mathscr{M})_{x_{i}}=H_{x_{i}} \oplus \mathfrak{h} \otimes \mathbb{C}\left(\left(t-x_{i}\right)\right)$ on which $\mathbf{1}_{x_{i}}$ acts as 1 becomes a smooth module of level $-h^{\vee}$ over the local copy $\widehat{\mathfrak{g}}_{x_{i}}$ of $\widehat{\mathfrak{g}}$ at $x_{i}$, for each $i=1, \ldots, p$. Likewise, every smooth module over $\mathrm{L}(\mathscr{M})_{0}^{\Gamma}=H_{0}^{\Gamma} \oplus(\mathfrak{h} \otimes \mathbb{C}((t)))^{\Gamma, 0}$ on which
$\mathbf{1}_{0}$ acts as $\frac{1}{T}$ becomes a smooth module of level $-h^{\vee} / T$ over $\hat{\mathfrak{g}}_{0}^{\Gamma}$, and every co-smooth module over $\mathrm{L}\left(\mathscr{M}^{\mathrm{op}}\right)_{\infty}^{\Gamma}=H_{\infty}^{\mathrm{op}, \Gamma} \oplus\left(\mathfrak{h} \otimes \mathbb{C}\left(\left(t^{-1}\right)\right)\right)^{\Gamma, 0}$ on which $\mathbf{1}_{\infty}$ acts as $\frac{1}{T}$ becomes a co-smooth module over $\widehat{\mathfrak{g}}_{\infty}^{\mathrm{op}, \Gamma}$ of level $-h^{\vee} / T$.

In particular, these statements apply to the Wakimoto modules $W_{\nu_{i}}, i=1, \ldots, p, W_{\nu_{0}}^{\Gamma}$ and $W_{\nu_{\infty}}^{\Gamma, \vee}$. We shall need the following facts about the structure of these modules.

For each $i=1, \ldots, p$, let $\widetilde{W}_{\nu_{i}}$ denote the linear span of states of the form $a_{\alpha_{1}}^{*}[0] \ldots a_{\alpha_{k}}^{*}[0]| \rangle_{x_{i}}$, $k \in \mathbb{Z}_{\geq 0}$.

Lemma 4.2. Suppose $\nu_{i}$ has at most a simple pole. Then the subspace $\widetilde{W}_{\nu_{i}}$ is stable under the action of the subalgebra $\mathfrak{g} \subset \widehat{\mathfrak{g}}_{x_{i}}$ generated by zero-modes, $X[0]_{x_{i}}$ with $X \in \mathfrak{g}$, and there is an isomorphism of left $U(\mathfrak{g})$-modules

$$
\begin{equation*}
\widetilde{W}_{\nu_{i}} \cong M_{\mathrm{res}_{t-x_{i}} \nu_{i}}^{*} \tag{4.16}
\end{equation*}
$$

Moreover the subspace $\widetilde{W}_{\nu_{i}}$ is annihilated by all strictly positive modes $X[n]_{x_{i}}, n>0, X \in \mathfrak{g}$.
Proof. See [FFR94, $\S 5]$ for the isomorphism (4.16). Note that $a_{\alpha}[n]$ and $a_{\alpha}^{*}[n]$ act as zero on $\widetilde{W}_{\nu_{i}}$ for all $n>0$. The fact that $\nu_{i}$ has at most a simple pole means that $b_{k}[n]$ acts as zero on $\widetilde{W}_{\nu_{i}}$ for all $n>0$. The "moreover" part follows since Feigin-Frenkel homomorphism respects the $\mathbb{Z}$-grading of $\mathbb{V}_{0}^{-h^{\vee}}$ and $\mathbb{W}_{0}$ and their big Lie algebras.

Let $\widetilde{W}_{\nu_{0}}^{\Gamma} \subset W_{\nu_{0}}^{\Gamma}$ denote the linear span of states of the form $a_{\left(0, \alpha_{1}\right)}^{*}[0] \ldots a_{\left(0, \alpha_{k}\right)}^{*}[0]| \rangle_{0}, k \in \mathbb{Z}_{\geq 0}$, $\alpha_{1}, \ldots, \alpha_{k} \in \Delta_{0}^{+}$.

Lemma 4.3. Suppose $\nu_{0}$ has at most a simple pole. Then the subspace $\widetilde{W}_{\nu_{0}}^{\Gamma}$ is stable under the action of the subalgebra $\mathfrak{g}^{\sigma} \subset \widehat{\mathfrak{g}}_{0}^{\Gamma}$ generated by zero-modes, $X[0]_{0}$ with $X \in \mathfrak{g}^{\sigma}$, and there is an isomorphism of left $U\left(\mathfrak{g}^{\sigma}\right)$-modules

$$
\begin{equation*}
\widetilde{W}_{\nu_{0}}^{\Gamma} \cong M_{\frac{1}{T}\left(\operatorname{res}_{t}\left(\nu_{0}\right)-\Lambda_{0}\right)}^{*, \sigma} \tag{4.17}
\end{equation*}
$$

where for $\lambda \in \mathfrak{h}^{*, \sigma}$ we denote by $M_{\lambda}^{*, \sigma}$ the contragredient Verma module over $\mathfrak{g}^{\sigma}$ of highest weight $\lambda$, and where $\Lambda_{0}$ is the weight given in (3.7).

Moreover the subspace $\widetilde{W}_{\nu_{0}}^{\Gamma}$ is annihilated by all strictly positive modes $X[n]_{0}, n>0, X \in \mathfrak{g}$.
Proof. The first part is [VY16a, Proposition 4.4], and the "moreover" part is again on $\mathbb{Z}$-grading grounds.

Concerning the Wakimoto module at infinity, we shall need two results. Recall the definition of $\chi$ in $\S 3.5$. The first result applies only to the case $\chi=0$, so that $t^{2} \nu_{\infty}$ has at most a simple pole in $t^{-1}$.

Let $\widetilde{W}_{\nu \infty}^{\Gamma, V} \subset W_{\nu \infty}^{\Gamma, V}$ denote the span of vectors of the form $\langle | a_{\left(0, \alpha_{1}\right)}[0] \ldots a_{\left(0, \alpha_{k}\right)}[0], k \in \mathbb{Z}_{\geq 0}$, $\alpha_{1}, \ldots, \alpha_{k} \in \Delta_{0}^{+}$.

Lemma 4.4. Suppose $t^{2} \nu_{\infty}$ has at most a simple pole in $t^{-1}$. Then the subspace $\widetilde{W}_{\nu_{\infty}}^{\Gamma, V}$ is stable under the action of the subalgebra $\mathfrak{g}^{\sigma} \subset \hat{\mathfrak{g}}_{\infty}^{\Gamma}$ generated by zero-modes, $X[0]_{\infty}$ with $X \in \mathfrak{g}^{\sigma}$, and there is an isomorphism of right $U\left(\mathfrak{g}^{\sigma}\right)$-modules (or equivalently left $U\left(\mathfrak{g}^{\sigma, \mathrm{op}}\right)$ modules)

$$
\begin{equation*}
\widetilde{W}_{\nu_{\infty}}^{\Gamma, \vee} \cong \operatorname{Hom}_{\mathbb{C}}\left(M_{\frac{1}{T}\left(-\operatorname{res}_{t^{-1}}\left(t^{2} \nu_{\infty}\right)-\Lambda_{0}\right)}^{*, \sigma}, \mathbb{C}\right) \cong M_{\frac{1}{T}\left(-\operatorname{res}_{t^{-1}}\left(t^{2} \nu_{\infty}\right)-\Lambda_{0}\right)}^{\sigma, \varphi} . \tag{4.18}
\end{equation*}
$$

Moreover the subspace $\widetilde{W}_{\nu_{\infty}}^{\Gamma, \vee}$ is annihilated by all strictly negative modes $X[n]_{\infty}, n<0, X \in \mathfrak{g}^{\sigma}$.
Proof. The condition that $t^{2} \nu_{\infty}$ has at most a simple pole in $t^{-1}$ means that $\langle | b_{(k, \alpha)}[n]=0$ for all $n<0$. We have $\langle | a_{(k, \alpha)}[n]=0$ and $\langle | a_{(k, \alpha)}^{*}[n]=0$ for all $n<0$. Therefore on $\mathbb{Z}$-grading grounds, for all $X \in \mathfrak{g}, X[n]$ acts as zero on $\widetilde{W}_{\nu \infty}^{\Gamma, v}$ for all $n<0$, and the only terms in the free-field expressions for the zero-modes $X[0]$ which survive on vectors in $\widetilde{W}_{\nu_{\infty}}^{\Gamma, \vee}$ are those involving only the zero-modes of $a$ 's, $a^{*}$ 's, and $b$ 's. In particular since only the zero-modes of $b$ 's contribute,

$$
\widetilde{W}_{\nu_{\infty}}^{\Gamma, \vee} \cong \widetilde{W}_{-\lambda / t}^{\Gamma, v} \quad \text { where } \quad \lambda:=-\operatorname{res}_{t^{-1}}\left(t^{2} \nu_{\infty}\right)
$$

as modules over the copy of algebra $\mathfrak{g}^{\sigma}$ generated by zero-modes.
Also on $\mathbb{Z}$-grading grounds, the only terms in the free-field expressions for the zero modes $X[0]$ which survive on vectors in $\widetilde{W}_{\lambda / t}^{\Gamma}$ are those involving only the zero modes of $a$ 's, $a^{*}$ 's, and $b$ 's.

Now $W_{-\lambda / t}^{\Gamma, V}$ and $W_{\lambda / t}^{\Gamma}$ are Fock modules over a copy of $H(\mathfrak{g})^{\Gamma}$ as in (4.12) and (4.10). We have the natural pairing between them, namely the bilinear map $B: W_{-\lambda / t}^{\Gamma, V} \times W_{\lambda / t}^{\Gamma} \rightarrow \mathbb{C}$ defined by $B(\langle |,| \rangle)=1$ and $B(\langle | X, Y| \rangle)=B(\langle |, X Y| \rangle)=B(\langle | X Y,| \rangle)$ for $X, Y \in U\left(H(\mathfrak{g})^{\Gamma}\right)$. This pairing is non-degenerate. In this way the right $U\left(H(\mathfrak{g})^{\Gamma}\right)$-module $W_{-\lambda / t}^{\Gamma, V}$ is the (suitably restricted) dual of the left $U\left(H(\mathfrak{g})^{\Gamma}\right)$-module $W_{\lambda / t}^{\Gamma}$. The pairing restricts to a non-degenerate pairing of the grade zero subspaces, $\widetilde{W}_{-\lambda / t}^{\Gamma, \vee}$ and $\widetilde{W}_{\lambda / t}^{\Gamma}$. So $\widetilde{W}_{-\lambda / t}^{\Gamma, \vee}$ is the dual of $\widetilde{W}_{\lambda / t}^{\Gamma}$ as a module over the algebra of zero-modes of $a$ 's and $a^{*}$ 's, and also of $b$ 's since $\langle | \mid b_{(k, a)}[0]=\lambda\left(H_{(k, a)}\right)\langle |$ and $b_{(k, a)}[0]| \rangle=| \rangle \lambda\left(H_{(k, a)}\right)$.

We conclude that $\widetilde{W}_{-\lambda / t}^{\Gamma, \vee}$ is the dual $\operatorname{Hom}_{\mathbb{C}}\left(\widetilde{W}_{\lambda / t}^{\Gamma}, \mathbb{C}\right)$ of $\widetilde{W}_{\lambda / t}^{\Gamma}$ as a module over the algebra of zero-modes $X[0], X \in \mathfrak{g}^{\sigma}$. And Lemma 4.3 shows that $\widetilde{W}_{\lambda / t}^{\Gamma} \cong M_{\frac{1}{T}\left(\lambda-\Lambda_{0}\right)}^{*, \sigma}$. The result follows.

Second, we need a result valid for all $\chi$.
Let $\mathbb{C}_{\chi}$ denote the one-dimensional representation of $\left(\mathfrak{g}^{\text {op }} \otimes t^{-1} \mathbb{C}\left[\left[t^{-1}\right]\right]\right)^{\Gamma}$ defined by

$$
\left(\mathfrak{g}^{\mathrm{op}} \otimes t^{-1} \mathbb{C}\left[\left[t^{-1}\right]\right]\right)^{\Gamma} \longrightarrow\left(\mathfrak{g}^{\mathrm{op}} \otimes t^{-1} \mathbb{C}\left[\left[t^{-1}\right]\right]\right)^{\Gamma} /\left(\mathfrak{g}^{\mathrm{op}} \otimes t^{-2} \mathbb{C}\left[\left[t^{-1}\right]\right]\right)^{\Gamma} \simeq_{\mathbb{C}} \Pi_{-1} \mathfrak{g} \xrightarrow{\chi} \mathbb{C} .
$$

Lemma 4.5. Suppose that $t^{2} \nu_{\infty}$ has at most a double pole in $t^{-1}$, with 2-residue $-\chi$. That is, suppose that

$$
\nu_{\infty}(t)=-\chi+\mathcal{O}(1 / t)
$$

Then there is an isomorphism of $\left(\mathfrak{g}^{\text {op }} \otimes t^{-1} \mathbb{C}\left[\left[t^{-1}\right]\right]\right)^{\Gamma}$-modules

$$
\begin{equation*}
\mathbb{C}\langle | \cong \mathbb{C}_{\chi} . \tag{4.19}
\end{equation*}
$$

Proof. It follows from the explicit form of the Feigin-Frenkel homomorphism and the definition of the quasi-module map $Y_{W}$, cf. appendix B and in particular Proposition B.3, that the vector $\langle |$ of (4.13) obeys

$$
\langle |\left(\mathfrak{g} \otimes t^{-2} \mathbb{C}\left[\left[t^{-1}\right]\right]\right)^{\Gamma}=0, \quad\langle |\left(\Pi_{-1} A\right)[-1]=\frac{1}{T} \chi(A)\langle | .
$$

4.5. The generators $G_{\alpha}[n]$. The Verma module $M_{\lambda}=U(\mathfrak{g}) \otimes_{U(\mathfrak{h} \oplus \mathfrak{n})} \mathbb{C} v_{\lambda} \cong_{\mathbb{C}} U\left(\mathfrak{n}_{-}\right) \otimes_{\mathbb{C}} \mathbb{C} v_{\lambda}$, $\S 3.1$, is by definition a left module over $U(\mathfrak{g})$. In particular it is a left module over the subalgebra $U\left(\mathfrak{n}_{-}\right) \subset U(\mathfrak{g})$. But $M_{\lambda}$ also admits a second left action, call it $\triangleright$, of this subalgebra $U\left(\mathfrak{n}_{-}\right)$. These
two left actions of $U\left(\mathfrak{n}_{-}\right)$are mutually commuting. They are given by

$$
X .\left(n \otimes v_{\lambda}\right)=X n \otimes v_{\lambda} \quad \text { and } \quad X \triangleright\left(n \otimes v_{\lambda}\right)=-n X \otimes v_{\lambda}, \quad X \in \mathfrak{n}_{-}, n \in U\left(\mathfrak{n}_{-}\right) .
$$

In particular

$$
\begin{equation*}
X .\left(1 \otimes v_{\lambda}\right)=-X \triangleright\left(1 \otimes v_{\lambda}\right) . \tag{4.20}
\end{equation*}
$$

Correspondingly the contragredient Verma module admits a second left action of $U(\mathfrak{n})$. We write $G_{\alpha}, \alpha \in \Delta^{+}$, for the generators of this second copy of $\mathfrak{n}$, which we denote $\mathfrak{n}_{\langle G\rangle}$. So we have $\left[G_{\alpha}, E_{\beta}\right]=0$ for all $\alpha, \beta \in \Delta^{+}$, where $E_{\alpha}$ are the generators of $\mathfrak{n}$ from $\S 3.1$.

We have the vacuum Verma module of $\mathfrak{n}_{\langle G\rangle} \otimes \mathbb{C}((t))$,

$$
\begin{equation*}
U\left(\mathfrak{n}_{\langle G\rangle} \otimes \mathbb{C}((t))\right) \otimes_{U\left(\mathfrak{n}_{\langle G\rangle} \otimes \mathbb{C}[t t]\right)} \mathbb{C} v_{0} \tag{4.21}
\end{equation*}
$$

where $v_{0}$ is a nonzero vector such that $\mathfrak{n}_{\langle G\rangle} \otimes \mathbb{C}[[t]] v_{0}=0$. It has a natural vertex algebra structure, in which $Y\left(A[-1] v_{0}, x\right)=\sum_{n \in \mathbb{Z}} A[n] x^{-n-1}$ for all $A \in \mathfrak{n}_{\langle G\rangle}$. There is a free-field realization of this vertex algebra, i.e. there is an injective homomorphism from this vertex algebra into $\mathbb{W}_{0}$. It is defined by

$$
G_{\alpha}[-1]| \rangle=\sum_{\beta \in \Delta^{+}} R_{\alpha}^{\beta}\left(a^{*}[0]\right) a_{\beta}[-1]| \rangle, \quad \alpha \in \Delta^{+},
$$

for certain polynomials $R_{\alpha}^{\beta}$ homogeneous of grade $\beta-\alpha$.
By means of this free-field realization, the Wakimoto module $W_{\nu_{i}}$ becomes a module not only over $\widehat{\mathfrak{g}}_{x_{i}}$ but also over a copy of $\mathfrak{n}_{\langle G\rangle} \otimes \mathbb{C}\left(\left(t-x_{i}\right)\right)$ whose generators we denote $G_{\alpha}[n]_{x_{i}}:=G_{\alpha} \otimes\left(t-x_{i}\right)^{n}$. With the obvious modifications, the statement of Lemma 4.2 holds for this copy of $\mathfrak{n}_{\langle G\rangle} \otimes \mathbb{C}\left(\left(t-x_{i}\right)\right)$. Namely, the subspace $\widetilde{W}_{\nu}$ is stable under the subalgebra $U\left(\mathfrak{n}_{\langle G\rangle}\right) \subset U\left(\mathfrak{n}_{\langle G\rangle} \otimes \mathbb{C}((t))\right)$ generated by zero modes, and (4.16) is an isomorphism of modules over $\mathfrak{n}_{\langle G\rangle}$, i.e. $G_{\alpha}[0]$ acts as $G_{\alpha}$.

The analogous statements hold for $\mathfrak{n}_{\langle G\rangle}^{\sigma}$ too. Namely we have generators $G_{(0, \alpha)}, \alpha \in \Delta_{0}^{+}$which act from the left on contragredient Verma modules $M_{\lambda}^{*, \sigma}$ over $\mathfrak{n}^{\sigma}$. These generators commute with the generators $E_{(0, \alpha)}$ of the standard left action of $U\left(\mathfrak{n}^{\sigma}\right)$.

Lemma 4.6. The subspace $\widetilde{W}_{\nu_{0}}^{\Gamma}$ is stable under the action of the subalgebra $\mathfrak{n}_{\langle G\rangle}^{\sigma} \subset\left(\mathfrak{n}_{\langle G\rangle} \otimes \mathbb{C}((t))\right)^{\Gamma}$ generated by zero-modes, $X[0]_{0}$ with $X \in \mathfrak{n}_{\langle G\rangle}^{\sigma}$, and (4.17) is an isomorphism of left $U\left(\mathfrak{n}_{\langle G\rangle}^{\sigma}\right)$ modules. Moreover the subspace $\widetilde{W}_{\nu_{0}}^{\Gamma}$ is annihilated by all strictly positive modes $X[n]_{0}, n>0$, $X \in \mathfrak{n}_{\langle G\rangle}$.

Proof. The first part of the proof of [VY16a, Proposition 4.4] shows that (4.17) is an isomorphism of left $U\left(\mathfrak{n}^{\sigma}\right)$ modules. The argument that it is an isomorphism of $U\left(\mathfrak{n}_{\langle G\rangle}^{\sigma}\right)$-modules is line-by-line identical but with the polynomials $P_{\alpha}^{\beta}$ (which define the free-field realization of $E$ 's) replaced by $R_{\alpha}^{\beta}$.
4.6. Global Heisenberg algebra and coinvariants. Let $H_{\infty, \boldsymbol{x}, 0}^{\Gamma}$ be the commutative Lie algebra

$$
H_{\infty, \boldsymbol{x}, 0}^{\Gamma}:=\left(\mathfrak{n}_{\mathbb{C}} \otimes \mathbb{C}_{\infty, \Gamma \boldsymbol{x}, 0}(t)\right)^{\Gamma, 0} \oplus\left(\mathfrak{n}_{\mathbb{C}}^{*} \otimes \mathbb{C}_{\infty, \Gamma \boldsymbol{x}, 0}(t)\right)^{\Gamma,-1} .
$$

That is, an element of $H_{\infty, \boldsymbol{x}, 0}^{\Gamma}$ is a pair $(f(t), g(t))$ where $f(t)$ is a rational function valued in $\mathfrak{n}_{\mathbb{C}}$ with poles at most at the points $0, x_{1}, \ldots, x_{p}, \infty$ and obeying the equivariance condition $f(\omega t)=\sigma f(t)$, and where $g(t)$ is a rational function valued in $\mathfrak{n}_{\mathbb{C}}^{*}$ with poles at most at the points $0, x_{1}, \ldots, x_{p}, \infty$
and obeying the equivariance condition $g(\omega t)=\omega^{-1} L_{\sigma} g(t)$. By virtue of the residue theorem there is an embedding of Lie algebras,

$$
H_{\infty, x, 0}^{\Gamma} \longleftrightarrow H_{\infty, p, 0}
$$

by taking Laurent expansions, $f(t) \mapsto\left(-\iota_{t^{-1}} ; \iota_{t-x_{1}}, \ldots, \iota_{t-x_{p}} ; \iota_{t}\right) f(t)$. Cf. (2.8).
Let us write

$$
\mathfrak{h}_{\infty, \boldsymbol{x}, 0}^{\Gamma}:=\left(\mathfrak{h} \otimes \mathbb{C}_{0, \Gamma \boldsymbol{x}, \infty}(t)\right)^{\Gamma, 0} \quad \text { and } \quad \mathfrak{h}_{\infty, \boldsymbol{x}, 0}^{*, \Gamma}:=\left(\mathfrak{h}^{*} \otimes \mathbb{C}_{0, \Gamma \boldsymbol{x}, \infty}(t)\right)^{\Gamma,-1}
$$

There is an embedding of commutative Lie algebras

$$
\mathfrak{h}_{\infty, \boldsymbol{x}, 0}^{\Gamma} \longleftrightarrow \mathfrak{h}_{\infty, p, 0}
$$

and we have the following, which is [VY16a, Proposition 4.3] but now including the pole at $\infty$.
Lemma 4.7. The space of coinvariants $\mathbb{C} v_{\nu_{\infty}} \otimes \bigotimes_{i=1}^{p} \mathbb{C} v_{\nu_{i}} \otimes \mathbb{C} v_{\nu_{0}} / \mathfrak{h}_{\infty, \boldsymbol{x}, 0}^{\Gamma}$ is one-dimensional if and only if there exists a $\nu(t) \in \mathfrak{h}_{\infty, \boldsymbol{x}, 0}^{*, \Gamma}$ such that

$$
\left(\nu_{\infty} ; \nu_{1}, \ldots, \nu_{p} ; \nu_{0}\right)=\left(-\iota_{t^{-1}} ; \iota_{t-x_{1}}, \ldots, \iota_{t-x_{p}} ; \iota_{t}\right) \nu(t)
$$

Otherwise it is zero-dimensional.
Define

$$
\begin{aligned}
H_{\infty, p, 0}^{+}:=\left(\mathfrak{n}_{\mathbb{C}} \otimes \mathbb{C}\left[\left[t^{-1}\right]\right]\right)^{\Gamma, 0} \oplus\left(\mathfrak{n}_{\mathbb{C}}^{*}\right. & \left.\otimes \mathbb{C}\left[\left[t^{-1}\right]\right]\right)^{\Gamma,-1} \\
& \oplus \bigoplus_{i=1}^{p}\left(\mathfrak{n}_{\mathbb{C}} \oplus \mathfrak{n}_{\mathbb{C}}^{*}\right) \otimes \mathbb{C}\left[\left[t-x_{i}\right]\right] \\
& \oplus\left(\mathfrak{n}_{\mathbb{C}} \otimes \mathbb{C}[[t]]\right)^{\Gamma, 0} \oplus\left(\mathfrak{n}_{\mathbb{C}}^{*} \otimes \mathbb{C}[[t]]\right)^{\Gamma,-1} \oplus \mathbb{C} 1
\end{aligned}
$$

Then

$$
H_{\infty, p, 0}=H_{\infty, p, 0}^{+}+H_{\infty, \boldsymbol{x}, 0}^{\Gamma} \quad \text { and } \quad H_{\infty, p, 0}^{+} \cap H_{\infty, \boldsymbol{x}, 0}^{\Gamma}=\{0\}
$$

We may regard $M_{\infty}^{\Gamma, \vee}$ as a left module over $U\left(H_{\infty}^{\Gamma, o p}\right)$ and then

$$
\begin{equation*}
\left(\mathrm{M}_{\infty}^{\Gamma, \vee} \otimes \bigotimes_{i=1}^{p} \mathrm{M}_{x_{i}} \otimes \mathrm{M}_{0}^{\Gamma}\right)=U\left(H_{\infty, p, 0}\right) \otimes_{U\left(H_{\infty, p, 0}^{+}\right)} \mathbb{C}\langle | \otimes| \rangle_{x_{1}} \otimes \ldots \otimes| \rangle_{x_{p}} \otimes| \rangle_{0} \tag{4.22}
\end{equation*}
$$

Hence

$$
\left(\mathrm{M}_{\infty}^{\Gamma, \vee} \otimes \bigotimes_{i=1}^{p} \mathrm{M}_{x_{i}} \otimes \mathrm{M}_{0}^{\Gamma}\right) / H_{\infty, \boldsymbol{x}, 0}^{\Gamma} \cong_{\mathbb{C}} \mathbb{C}\langle | \otimes| \rangle_{x_{1}} \otimes \ldots \otimes| \rangle_{x_{p}} \otimes| \rangle_{0} \cong_{\mathbb{C}} \mathbb{C}
$$

Therefore for any $\nu(t) \in \mathfrak{h}_{\infty, \boldsymbol{x}, 0}^{*, \Gamma}$ the space of coinvariants

$$
\begin{equation*}
\left(W_{-\iota_{t}-1 \nu(t)}^{\Gamma, \vee} \otimes \bigotimes_{i=1}^{p} W_{\iota_{t-x_{i}} \nu(t)} \otimes W_{\iota_{t} \nu(t)}^{\Gamma}\right) /\left(H_{\infty, \boldsymbol{x}, 0}^{\Gamma} \oplus \mathfrak{h}_{\infty, \boldsymbol{x}, 0}^{\Gamma}\right) \tag{4.23}
\end{equation*}
$$

has dimension one. That means there is a unique $\left(H_{\infty, \boldsymbol{x}, 0}^{\Gamma} \oplus \mathfrak{h}_{\infty, \boldsymbol{x}, 0}^{\Gamma}\right)$-invariant linear functional, call it $\tau_{\nu(t)}$,

$$
\begin{equation*}
\tau_{\nu(t)}: W_{-\iota_{t}{ }^{-1} \nu(t)}^{\Gamma, \vee} \otimes \bigotimes_{i=1}^{p} W_{\iota_{t-x_{i}} \nu(t)} \otimes W_{\iota_{t} \nu(t)}^{\Gamma} \longrightarrow \mathbb{C} \tag{4.24}
\end{equation*}
$$

normalised such that $\tau_{\nu(t)}\left(\langle | \otimes| \rangle_{x_{1}} \otimes \ldots \otimes| \rangle_{x_{p}} \otimes| \rangle_{0}\right)=1$.

By functoriality, i.e. [VY16b, Theorem 6.2] and [VY16b, Corollary 6.6] generalised to the situation of Appendix B for coinvariants of a tensor product of modules including one attached to infinity, the functional $\tau_{\nu(t)}$ is also invariant under $\mathfrak{g}_{\infty, \boldsymbol{x}, 0}^{\Gamma}$.
4.7. Proof of Theorem 3.2. Let $c(1), \ldots, c(m) \in I, \lambda_{1}, \ldots, \lambda_{N} \in \mathfrak{h}^{*}, \lambda_{0} \in \mathfrak{h}^{*, \sigma}$ and $\chi$ be as in $\S 3$. Recall

$$
\left(x_{1}, \ldots, x_{p}\right)=\left(z_{1}, \ldots, z_{N}, w_{1}, \ldots, w_{m}\right) .
$$

We now fix

$$
\begin{equation*}
\nu(t):=\chi+\sum_{r \in \mathbb{Z}_{T}}\left(\sum_{i=1}^{N} \frac{L_{\sigma}^{r} \lambda_{i}}{t-\omega^{r} z_{i}}-\sum_{j=1}^{m} \frac{L_{\sigma}^{r} \alpha_{c(j)}}{t-\omega^{r} w_{j}}\right)+\frac{T \lambda_{0}+\Lambda_{0}}{t} \in \mathfrak{h}_{\infty, \boldsymbol{x}, 0}^{*, \Gamma}, \tag{4.25}
\end{equation*}
$$

where $\Lambda_{0}$ is as in (3.5).
The following is [FFR94, Lemma 2]. We write $G_{i}:=G_{\alpha_{i}}$. (Recall $G_{\alpha}$ for $\alpha \in \Delta^{+}$a positive root was defined in §4.5.)

Lemma 4.8. Let $\mu(t)$ be a highest weight of the form

$$
\mu(t)=-\frac{\alpha_{i}}{t}+\sum_{n=0}^{\infty} \mu^{(n)} t^{n}
$$

$\mu^{(n)} \in \mathfrak{h}^{*}$. Then the vector $G_{i}[-1]| \rangle \in W_{\mu(t)}$ is singular for $\widehat{\mathfrak{g}}$ (i.e. $X[n] G_{i}[-1]| \rangle=0$ for all $n \geq 0$ and all $X \in \mathfrak{g}$ ) if and only if

$$
\left\langle\alpha_{i}, \mu^{(0)}\right\rangle=0
$$

Let $\nu_{j}^{(0)}$ be the constant term in the Laurent expansion of $\nu(t)$ about $w_{j}$, for $j=1, \ldots, m$. Note that the Bethe equations (3.8) are equivalent to the statement that

$$
\begin{equation*}
\left\langle\alpha_{c(j)}, \nu_{j}^{(0)}\right\rangle=0, \quad j=1, \ldots, m . \tag{4.26}
\end{equation*}
$$

By Lemma 4.8, under this condition the vector $G_{c(j)}[-1]| \rangle_{w_{j}} \in W_{\iota t-w_{j}} \nu(t)$ is singular for $\widehat{\mathfrak{g}}_{\left(w_{j}\right)}$.
Henceforth, suppose that the Bethe equations are indeed satisfied. It follows that the linear functional

$$
\begin{equation*}
\psi_{\nu(t)}: W_{-\iota_{t}-1 \nu(t)}^{\Gamma, V} \otimes \bigotimes_{i=1}^{N} W_{\iota t-z_{i} \nu(t)} \otimes W_{\iota t}^{\Gamma} \quad \rightarrow \mathbb{C} \tag{4.27}
\end{equation*}
$$

defined by

$$
\psi_{\nu(t)}\left(v_{\infty}, v_{1}, \ldots, v_{N}, v_{0}\right)=\tau_{\nu(t)}\left(v_{\infty}, v_{1}, \ldots, v_{N}, G_{c(1)}[-1]| \rangle_{w_{1}}, \ldots, G_{c(m)}[-1]| \rangle_{w_{m}}, v_{0}\right)
$$

is invariant under the Lie subalgebra $\mathfrak{g}_{\infty, z, 0}^{\Gamma} \subset \mathfrak{g}_{\infty, x, 0}^{\Gamma}$.
Now we use the set-up of $\S 2.4$ and specialize by setting

$$
\begin{align*}
& \mathcal{M}_{z_{i}}=M_{\lambda_{i}}^{*} \cong \widetilde{W}_{\iota t-z_{i} \nu(t)} \\
& \mathcal{M}_{0} \subset W_{\iota_{t-z_{i}} \nu(t)}^{*, \sigma}, \quad i=1, \ldots, N, \\
& \mathcal{M}_{\lambda_{0}}=\mathbb{C}_{\chi} \cong \mathbb{C}\langle |  \tag{4.28}\\
& \cong W_{\iota t \nu(t)}^{\Gamma} \quad \subset W_{-\iota_{t}-1 \nu(t)}^{\Gamma},
\end{align*}
$$

Here the isomorphisms are as in Lemmas 4.2, 4.3 and 4.5, given that we have $\operatorname{res}_{t-z_{i}} \nu(t)=\lambda_{i}$ and $\frac{1}{T}\left(\operatorname{res}_{t} \nu(t)-\Lambda_{0}\right)=\lambda_{0}$.

Let $\psi$ denote the restriction of the linear functional $\psi_{\nu(t)}$ of (4.27) to the subspace

$$
\begin{equation*}
\mathbb{C}_{\chi} \otimes \bigotimes_{i=1}^{N} M_{\lambda_{i}}^{*} \otimes M_{\lambda_{0}}^{*, \sigma} \cong \mathbb{C}\langle | \otimes \bigotimes_{i=1}^{N} \widetilde{W}_{\iota_{t-z_{i}} \nu(t)} \otimes \widetilde{W}_{\iota t \nu(t)}^{\Gamma} \tag{4.29}
\end{equation*}
$$

Such a functional is the same thing as a vector

$$
\begin{equation*}
\psi \in \bigotimes_{i=1}^{N} M_{\lambda_{i}} \otimes M_{\lambda_{0}}^{\sigma} \tag{4.30}
\end{equation*}
$$

To complete the proof of Theorem 3.2, we shall now establish that
(i) $\psi$ is the Bethe vector defined in (3.3), and
(ii) it is an eigenstate of the Gaudin algebra with the given eigenvalue for the quadratic Hamiltonians.

We consider (i) first.
For notational convenience let us reorder tensor factors in the argument of $\tau_{\nu(t)}$ in such a way that

$$
\begin{equation*}
\tau_{\nu(t)}: W_{-\iota_{t}-1 \nu(t)}^{\Gamma, V} \otimes \bigotimes_{i=1}^{N} W_{\iota t-z_{i} \nu} \nu(t) \otimes W_{\iota t \nu(t)}^{\Gamma} \otimes \bigotimes_{j=1}^{m} W_{\iota t-w_{j} \nu(t)} \longrightarrow \mathbb{C} \tag{4.31}
\end{equation*}
$$

Lemma 4.9. For any $\boldsymbol{v} \otimes v_{0} \in \bigotimes_{i=1}^{N} \widetilde{W}_{\iota t-z_{i}} \nu(t) \otimes \widetilde{W}_{\iota t \nu(t)}^{\Gamma}$ and any $X_{1}, \ldots, X_{s} \in \mathfrak{n}_{\langle G\rangle}$, we have

$$
\begin{aligned}
& \tau_{\nu(t)}\left(\langle |, \boldsymbol{v}, v_{0}, X_{1}[-1]| \rangle_{w_{1}}, \ldots, X_{s-1}[-1]| \rangle_{w_{s-1}}, X_{s}[-1]| \rangle_{w_{s}},| \rangle_{w_{s+1}}, \ldots,| \rangle_{w_{m}}\right) \\
& =\sum_{i=1}^{N} \sum_{k \in \mathbb{Z}_{T}} \frac{\tau_{\nu(t)}\left(\langle |,\left(\sigma^{k} X_{s}\right)[0]_{z_{i}} \boldsymbol{v}, v_{0}, X_{1}[-1]| \rangle_{w_{1}}, \ldots, X_{s-1}[-1]| \rangle_{w_{s-1}},| \rangle_{w_{s}}, \ldots,| \rangle_{w_{m}}\right)}{w_{s}-\omega^{-k} z_{i}} \\
& +\frac{\tau_{\nu(t)}\left(\langle |, \boldsymbol{v},\left(T \Pi_{0} X_{s}\right)_{(0)}^{W} v_{0}, X_{1}[-1]| \rangle_{w_{1}}, \ldots, X_{s-1}[-1]| \rangle_{w_{s-1}},| \rangle_{w_{s}}, \ldots,| \rangle_{w_{m}}\right)}{w_{s}} \\
& \quad+\sum_{j=1}^{s-1} \sum_{k \in \mathbb{Z}_{T}} \frac{1}{w_{s}-\omega^{-k} w_{j}} \tau_{\nu(t)}\left(\langle |, \boldsymbol{v}, v_{0}, X_{1}[-1]| \rangle_{w_{1}}, \ldots, X_{j-1}[-1]| \rangle_{w_{j-1}},\left[\sigma^{k}\left(X_{s}\right), X_{j}\right][-1]| \rangle_{w_{j}},\right. \\
& \left.X_{j+1}[-1]| \rangle_{w_{j+1}}, \ldots, X_{s-1}[-1]| \rangle_{w_{s-1}},| \rangle_{w_{s}}, \ldots,| \rangle_{w_{m}}\right) .
\end{aligned}
$$

Proof. Let us write $\boldsymbol{y}=X_{1}[-1]| \rangle_{w_{1}} \otimes \ldots \otimes X_{s-1}[-1]| \rangle_{w_{s-1}} \otimes| \rangle_{w_{s}} \otimes \ldots \otimes| \rangle_{w_{m}} \in \bigotimes_{j=1}^{m} W_{\iota_{t-w_{j}}} \nu(t)$ for brevity. We have

$$
\begin{equation*}
0=\sum_{k \in \mathbb{Z}_{T}}\left[\frac{\sigma^{k} X_{s}[-1]| \rangle}{\omega^{-k} t-w_{s}} \cdot\left(\langle | \otimes \boldsymbol{v} \otimes v_{0} \otimes \boldsymbol{y}\right)\right] . \tag{4.32}
\end{equation*}
$$

The square brackets [.] on the right hand side denotes the class in the space of coinvariants

$$
\left(W_{-\iota_{t}-1}^{\Gamma, v} \nu(t) \otimes \bigotimes_{i=1}^{N} W_{\iota t-z_{i}} \nu(t) \otimes W_{\iota t}^{\Gamma} \nu(t) \otimes \bigotimes_{j=1}^{m} W_{\iota t-w_{j}} \nu(t)\right) /\left(H_{\infty, \boldsymbol{z}, 0, \boldsymbol{w}}^{\Gamma} \oplus \mathfrak{h}_{\infty, z, 0, \boldsymbol{w}}^{\Gamma}\right)
$$

and the equality in (4.32) follows from the straightforward generalisation of the isomorphism in [VY16b, Theorem 6.2] including the point at infinity. Now

$$
\begin{align*}
& \sum_{k \in \mathbb{Z}_{T}} \frac{\sigma^{k} X_{s}[-1]| \rangle}{\omega^{-k} t-w_{s}} \cdot\left(\langle | \otimes \boldsymbol{v} \otimes v_{0} \otimes \boldsymbol{y}\right) \\
&= \sum_{k \in \mathbb{Z}_{T}}\left(-\operatorname{res}_{t^{-1}} t^{2} \iota_{t^{-1}} \frac{1}{\omega^{-k} t-w_{s}}\langle | Y_{W}\left(\sigma^{k} X_{s}[-1]| \rangle, t\right) \otimes \boldsymbol{v} \otimes v_{0} \otimes \boldsymbol{y}\right. \\
&+\langle | \otimes \sum_{i=1}^{N} \operatorname{res}_{t-z_{i}} \iota_{t-z_{i}} \frac{1}{\omega^{-k} t-w_{s}} Y_{M}\left(\sigma^{k} X_{s}[-1]| \rangle, t-z_{i}\right)_{z_{i}} \boldsymbol{v} \otimes v_{0} \otimes \boldsymbol{y} \\
& \quad+\langle | \otimes \boldsymbol{v} \otimes \operatorname{res}_{t} \iota_{t} \frac{1}{\omega^{-k} t-w_{s}} Y_{W}\left(\sigma^{k} X_{s}[-1]| \rangle, t\right) v_{0} \otimes \boldsymbol{y} \\
&\left.+\langle | \otimes \boldsymbol{v} \otimes v_{0} \otimes \sum_{j=1}^{m} \operatorname{res}_{t-w_{j}} \iota_{t-w_{j}} \frac{1}{\omega^{-k} t-w_{s}} Y_{M}\left(\sigma^{k} X_{s}[-1]| \rangle, t-w_{j}\right)_{w_{j}} \boldsymbol{y}\right) . \tag{4.33}
\end{align*}
$$

The $\operatorname{res}_{t-w_{s}}$ term here is $\langle | \otimes \boldsymbol{v} \otimes v_{0} \otimes X_{1}[-1]| \rangle_{w_{1}} \otimes \ldots \otimes X_{s-1}[-1]| \rangle_{w_{s-1}} \otimes X_{s}[-1]| \rangle_{w_{s}} \otimes| \rangle_{w_{s+1}} \otimes \ldots \otimes| \rangle_{w_{m}}$. (The expansion of $1 /\left(\omega^{-k} t-w_{s}\right)$ at $t=w_{s}$ has a regular part but it does not contribute because $\left.\left(X_{s}[-1]| \rangle\right)(n)\left\rangle=X_{s}[n]\right|\right\rangle$ is zero for all $n \geq 0$.)

Consider the first term on the right. We have

$$
\begin{equation*}
\operatorname{res}_{t^{-1}} t^{2} \iota_{t^{-1}} \frac{1}{\omega^{-k} t-w_{s}}\langle | Y_{W}\left(\sigma^{k} X_{s}[-1]| \rangle, t\right)=\sum_{n \geq 0} w_{s}^{n} \omega^{k(n+1)}\langle |\left(\sigma^{k} X_{s}[-1]| \rangle\right)_{(-n-1)}^{W} \tag{4.34}
\end{equation*}
$$

and this vanishes on $\mathbb{Z}$-grading grounds. (The free-field generators of $\mathfrak{n}_{\langle G\rangle}$ involve $a, a^{*}$ but not $b$.) Consider the third term. We have

$$
\begin{equation*}
\left.\operatorname{res}_{t} \iota_{t} \frac{1}{\omega^{-k} t-w_{s}} Y_{W}\left(\sigma^{k} X_{s}[-1]| \rangle, t\right) v_{0}=-\sum_{n \geq 0} w_{s}^{-n-1} \omega^{-k n}\left(\sigma^{k} X_{s}[-1]| \rangle\right)\right)_{(n)}^{W} v_{0} \tag{4.35}
\end{equation*}
$$

and on grading grounds only the term $n=0$ in this sum contributes, since $v_{0} \in \widetilde{W}_{\iota t \nu(t)}$. The remaining terms are similar.

Recall that we can regard an element of a Verma module $M_{\lambda}$ as a linear map $M_{\lambda}^{*} \rightarrow \mathbb{C}$.

Lemma 4.10. The restriction of the linear functional

$$
\tau_{\nu(t)}\left(\langle |, \cdot, \ldots, \cdot, \cdot,| \rangle_{w_{1}}, \ldots,| \rangle_{w_{m}}\right): \bigotimes_{i=1}^{N} W_{\iota t-z_{i} \nu(t)} \otimes W_{\iota t \nu(t)}^{\Gamma} \rightarrow \mathbb{C}
$$

to the subspace

$$
\bigotimes_{i=1}^{N} \widetilde{W}_{\iota t-z_{i} \nu(t)} \otimes \widetilde{W}_{\iota t \nu(t)}^{\Gamma} \cong_{\mathfrak{g}^{\sigma}} \bigotimes_{i=1}^{N} M_{\lambda_{i}}^{*} \otimes M_{\lambda_{0}}^{*, \sigma} .
$$

is equal to the tensor product

$$
v_{\lambda_{1}} \otimes \ldots v_{\lambda_{N}} \otimes v_{\lambda_{0}}
$$

of the highest weight vectors $v_{\lambda_{i}} \in M_{\lambda_{i}}, i=1, \ldots, N$, and $v_{\lambda_{0}} \in M_{\lambda_{0}}^{\sigma}$.

Proof. The functional $\tau_{\nu(t)}\left(\langle |, \cdot, \ldots, \cdot, \cdot,| \rangle_{w_{1}}, \ldots,| \rangle_{w_{m}}\right)$ agrees with $v_{\lambda_{1}} \otimes \ldots v_{\lambda_{N}} \otimes v_{\lambda_{0}}$ on the vector $\left.\left\rangle_{z_{1}} \otimes \ldots \otimes\right|\right\rangle_{z_{N}} \otimes| \rangle_{z_{0}}$. Indeed, $\tau_{\nu(t)}$ was normalized such that

$$
\tau_{\nu(t)}\left(\langle |,| \rangle_{z_{1}}, \ldots,| \rangle_{z_{N}},| \rangle_{z_{0}},| \rangle_{w_{1}}, \ldots,| \rangle_{w_{m}}\right)=1
$$

and by definition

$$
\left.\left(v_{\lambda_{1}} \otimes \ldots v_{\lambda_{N}} \otimes v_{\lambda_{0}}\right)\left(\left\rangle_{z_{1}}, \ldots,\right|\right\rangle_{z_{N}},| \rangle_{z_{0}}\right)=1 .
$$

It remains to show that the two agree on all other vectors in $\bigotimes_{i=1}^{N} \widetilde{W}_{\iota t-z_{i} \nu} \nu(t) \otimes \widetilde{W}_{\iota t \nu(t)}^{\Gamma}$. For that it is enough to show that, for any $\boldsymbol{v} \otimes v_{0} \in \bigotimes_{i=1}^{N} \widetilde{W}_{\iota t-z_{i}} \nu(t) \otimes \widetilde{W}_{\iota_{t} \nu(t)}^{\Gamma}$,

$$
\begin{equation*}
\tau_{\nu(t)}\left(\langle |, \boldsymbol{v}, a_{(0, \alpha)}^{*}[0] v_{0},| \rangle_{w_{1}}, \ldots,| \rangle_{w_{m}}\right)=0 \tag{4.36}
\end{equation*}
$$

for any $\alpha \in \Delta_{0}^{+}$, and

$$
\begin{equation*}
\tau_{\nu(t)}\left(\langle |, a_{\alpha}^{*}[0]_{\left(z_{i}\right)} \boldsymbol{v}, v_{0},| \rangle_{w_{1}}, \ldots,| \rangle_{w_{m}}\right)=0, \quad i=1, \ldots, N, \tag{4.37}
\end{equation*}
$$

for any $\alpha \in \Delta^{+}$.
To establish the equality (4.36) we use the invariance of $\tau_{\nu(t)}$ under the $\Gamma$-equivariant function $E_{(0, \alpha)}^{*} \otimes t^{-1} \in\left(\mathfrak{n}_{\mathbb{C}}^{*} \otimes \mathbb{C}_{0, \Gamma \boldsymbol{x}, \infty}(t)\right)^{\Gamma,-1}$. Namely, we have

$$
0=\tau_{\nu(t)}\left(\frac{E_{(0, \alpha)}^{*}}{t} \cdot\left(\langle | \otimes \boldsymbol{v} \otimes v_{0} \otimes| \rangle_{w_{1}} \otimes \ldots \otimes| \rangle_{w_{m}}\right)\right)
$$

and here

$$
\begin{align*}
\frac{E_{(0, \alpha)}^{*}}{t} \cdot\left(\langle | \otimes \boldsymbol{v} \otimes v_{0} \otimes| \rangle_{w_{1}}\right. & \otimes \ldots \otimes\left\rangle_{w_{m}}\right) \\
& =\langle | a_{(0, \alpha)}^{*}[0] \otimes \boldsymbol{v} \otimes v_{0} \otimes| \rangle_{w_{1}} \otimes \ldots \otimes| \rangle_{w_{m}} \\
+\sum_{i=1}^{N} & \langle | \otimes \iota_{t-z_{i}} \frac{\left(E_{(0, \alpha)}^{*}\right)^{(i)}}{t} \boldsymbol{v} \otimes v_{0} \otimes| \rangle_{w_{1}} \otimes \ldots \otimes| \rangle_{w_{m}} \\
& +\langle | \otimes \boldsymbol{v} \otimes a_{(0, \alpha)}^{*}[0] v_{0} \otimes| \rangle_{w_{1}} \otimes \ldots \otimes| \rangle_{w_{m}} \\
& \quad+\sum_{i=1}^{m}\langle | \otimes \boldsymbol{v} \otimes v_{0} \otimes| \rangle_{w_{1}} \otimes \ldots \otimes \iota \iota_{t-w_{i}} \frac{E_{(0, \alpha)}^{*}}{t}| \rangle_{w_{i}} \otimes \ldots \otimes| \rangle_{w_{m}} \tag{4.38}
\end{align*}
$$

The first line is zero by definition of $\langle |$. The second and fourth lines are zero on $\mathbb{Z}$-grading grounds (recall that $E_{\alpha}^{*} \otimes t^{0}=a_{\alpha}^{*}[1]$ ). This leaves only the third line, and thus we have (4.36) as required. The proof of (4.37) is similar (and is as in [VY16a, FFR94, ATY91]).

Recall the definition of $G_{\alpha}$ and $G_{(0, \alpha)}$ from $\S 4.5$. Let $F_{(0, \alpha)}:=\varphi\left(E_{(0, \alpha)}\right) \in \mathfrak{n}_{-}^{\sigma}$, for $\alpha \in \Delta_{0}^{+}$, be the negative root vectors of $\mathfrak{g}^{\sigma}$.

## Lemma 4.11.

(1) Let $v_{\lambda} \in M_{\lambda}$ be a highest weight vector. For any roots $\alpha^{(1)}, \ldots, \alpha^{(k)} \in \Delta^{+}$we have

$$
v_{\lambda}\left(G_{\alpha^{(1)}} G_{\alpha^{(2)}} \ldots G_{\alpha^{(k)}} \cdot\right)=(-1)^{k}\left(F_{\alpha^{(1)}} F_{\alpha^{(2)}} \ldots F_{\alpha^{(k)}} v_{\lambda}\right)(\cdot) .
$$

(2) Let $v_{\lambda_{0}} \in M_{\lambda_{0}}^{\sigma}$ be a highest weight vector. For any $\mathfrak{g}^{\sigma}$-weights $\alpha^{(1)}, \ldots, \alpha^{(k)} \in \Delta_{0}^{+}$we have

$$
v_{\lambda_{0}}\left(G_{\left(0, \alpha^{(1)}\right)} G_{\left(0, \alpha^{(2)}\right)} \ldots G_{\left(0, \alpha^{(k)}\right)} \cdot\right)=(-1)^{k}\left(F_{\left(0, \alpha^{(1)}\right)} F_{\left(0, \alpha^{(2)}\right)} \ldots F_{\left(0, \alpha^{(k)}\right)} v_{\lambda}\right)(\cdot)
$$

Proof. Consider part (1). It follows from (4.20) that $v_{\lambda_{i}}\left(E_{\alpha} \cdot\right)=-v_{\lambda_{i}}\left(G_{\alpha} \cdot\right)$. Because the two actions of $U(\mathfrak{n})$ commute, we therefore have

$$
\begin{aligned}
v_{\lambda}\left(G_{\alpha^{(1)}} G_{\alpha^{(2)}} \ldots G_{\alpha^{(k)}} \cdot\right)=- & v_{\lambda}\left(E_{\alpha^{(1)}} G_{\alpha^{(2)}} \ldots G_{\alpha^{(k)}} \cdot\right) \\
& =-v_{\lambda}\left(G_{\alpha^{(2)}} \ldots G_{\alpha^{(k)}} E_{\alpha^{(1)}} \cdot\right)=\cdots=(-1)^{k} v_{\lambda}\left(E_{\alpha^{(k)}} \ldots E_{\alpha^{(2)}} E_{\alpha^{(1)}} \cdot\right)
\end{aligned}
$$

for any roots $\alpha^{(1)}, \ldots, \alpha^{(k)} \in \Delta^{+}$. Part (1) follows by definition of the contragredient dual. For part (2) the argument is the same.

Lemmas 4.9, 4.10 and 4.11 together imply that the weight function $\psi$ of (4.30) is given recursively as in (3.3). Indeed, repeated use of the relations in Lemma 4.9 allow us to express

$$
\tau_{\nu(t)}\left(\langle |, v_{1}, \ldots, v_{N}, v_{0}, G_{c(1)}[-1]| \rangle_{w_{1}}, \ldots, G_{c(m)}[-1]| \rangle_{w_{m}}\right)
$$

as a linear combination of terms of the form

$$
\tau_{\nu(t)}\left(\langle |, v_{1}^{\prime}, \ldots, v_{N}^{\prime}, v_{0}^{\prime},| \rangle_{w_{1}}, \ldots,| \rangle_{w_{m}}\right)
$$

for certain $v_{0}^{\prime}, v_{1}^{\prime}, \ldots, v_{N}^{\prime}$, with each $v_{i}^{\prime}, 1 \leq i \leq N$, of the form $G_{\alpha^{(1)}} \ldots G_{\alpha^{(k)}}| \rangle$ and $v_{0}^{\prime}$ of the form $G_{\left(0, \alpha^{(1)}\right)} \ldots G_{\left(0, \alpha^{(k)}\right)}| \rangle$. Then Lemma 4.10 tells us that $\tau_{\nu(t)}\left(\langle |, v_{1}^{\prime}, \ldots, v_{N}^{\prime},| \rangle_{w_{1}}, \ldots,| \rangle_{w_{m}}, v_{0}\right)=$ $v_{\lambda_{0}}\left(v_{0}^{\prime}\right) \prod_{i=1}^{N} v_{\lambda_{i}}\left(v_{i}^{\prime}\right)$. So we can use Lemma 4.11 to exchange $G$ 's for $F$ 's. It follows that the vectors $\psi$ in (3.3) and (4.30) coincide, because the structure of the recursive definition of $\psi$ in (3.3) is chosen so as to match the relations of Lemma 4.9.

It remains to consider (ii), that is to show that $\psi$ is an eigenstate of the Gaudin algebra, and to compute its eigenvalue for the quadratic Hamiltonians. The argument is exactly as in [VY16a, FFT10] following [FFR94]. For completeness, let us recall it. Let $u \in \mathbb{C}^{\times}$be an additional nonzero point, with $\Gamma$-orbit disjoint from those of the $x_{i}, 1 \leq i \leq p$. To the point $u$ we assign a copy of the $H_{u} \oplus \mathfrak{h} \otimes \mathbb{C}((t-u))$-module $\mathbb{W}_{0}=\mathrm{M} \otimes \pi_{0}$, cf. (4.14). Then we have a unique $\left(H_{\infty, \boldsymbol{x}, u, 0}^{\Gamma} \oplus \mathfrak{h}_{\infty, \boldsymbol{x}, u, 0}^{\Gamma}\right)$ invariant linear functional

$$
\tau_{\nu(t)}: W_{-\iota_{t}-1 \nu(t)}^{\Gamma, V} \otimes \bigotimes_{i=1}^{N} W_{\iota t-z_{i}} \nu(t) \otimes W_{\iota t \nu(t)}^{\Gamma} \otimes \bigotimes_{j=1}^{m} W_{\iota t-w_{j} \nu} \nu(t) \otimes \mathbb{W}_{0} \longrightarrow \mathbb{C}
$$

normalised such that

$$
\tau_{\nu(t)}\left(\langle | \otimes| \rangle_{z_{1}} \otimes \ldots \otimes| \rangle_{z_{p}} \otimes| \rangle_{0} \otimes| \rangle_{w_{1}} \otimes \ldots \otimes| \rangle_{w_{p}} \otimes| \rangle_{u}\right)=1
$$

(Compare (4.24) and (4.31).)
Let $Z \in \mathfrak{z}(\widehat{\mathfrak{g}}) \subset \mathbb{V}_{0}^{k}$ be a singular vector in the vacuum Verma module. Recall the Feigin-Frenkel homomorphism of vertex algebras, $\rho: \mathbb{V}_{0}^{-h^{\vee}} \rightarrow \mathbb{W}_{0}$ from (4.15). Now, for any vector

$$
v \in \mathbb{C}_{\chi} \otimes \bigotimes_{i=1}^{N} M_{\lambda_{i}}^{*} \otimes M_{\lambda_{0}}^{*, \sigma}
$$

we may evaluate the quantity

$$
\begin{equation*}
\tau_{\nu(t)}\left(v, G_{c(1)}[-1]| \rangle, \ldots, G_{c(m)}[-1]| \rangle, \rho(Z)\right) \tag{4.39}
\end{equation*}
$$

in two ways. First, the functional $\tau_{\nu(t)}$ is $\mathfrak{g}_{\infty, \boldsymbol{x}, u, 0}^{\Gamma}$-invariant. If the cyclotomic Bethe equations (3.8) hold then for each $j, 1 \leq j \leq m$, the vector $G_{c(j)}[-1]| \rangle \in W_{\mu_{j}}$ is singular for the action of $\hat{\mathfrak{g}}_{\left(w_{j}\right)}$, as in Lemma 4.8. Therefore

$$
\begin{align*}
\tau_{\nu(t)}\left(v, G_{c(1)}[-1]| \rangle, \ldots, G_{c(m)}[-1]| \rangle, \rho(Z)\right) & =\tau_{\nu(t)}\left(Z(u) \cdot v, G_{c(1)}[-1]| \rangle, \ldots, G_{c(m)}[-1]| \rangle,| \rangle\right) \\
& =\psi(Z(u) \cdot v)=(\varphi(Z(u)) \cdot \psi)(v) . \tag{4.40}
\end{align*}
$$

On the other hand, using the invariance of the functional $\tau_{\nu(t)}$ under the rational function

$$
\frac{1}{(n-1)!}\left(\frac{\partial}{\partial u}\right)^{n-1} \sum_{r \in \mathbb{Z}_{T}} \frac{\sigma^{r} b_{s}}{\omega^{-r} t-u} \in \mathfrak{h}_{\infty, x, u, 0}^{\Gamma},
$$

and the fact that $\left[b_{s}[n], G_{c(i)}[-1]\right]=0$, we have that for all $w \in \mathbb{W}_{0}$,

$$
\begin{align*}
\tau_{\nu(t)}\left(v, G_{c(1)}[ \right. & \left.\left.-1]\left\rangle, \ldots, G_{c(m)}[-1]\right|\right\rangle, b_{s}[-n] w\right) \\
& =\tau_{\nu(t)}\left(v, G_{c(1)}[-1]| \rangle, \ldots, G_{c(m)}[-1]| \rangle, w\right)\left(\frac{1}{(n-1)!}\left(\frac{\partial}{\partial u}\right)^{n-1}\left\langle\alpha_{s}^{\vee}, \nu(u)\right\rangle\right) . \tag{4.41}
\end{align*}
$$

Here, cf. (4.25),

$$
\left\langle\alpha_{s}^{\vee}, \nu(u)\right\rangle=\left\langle\alpha_{s}^{\vee}, \chi\right\rangle+\sum_{r \in \mathbb{Z}_{T}}\left(\sum_{i=1}^{N} \frac{\left\langle\alpha_{s}^{\vee}, L_{\sigma}^{r} \lambda_{i}\right\rangle}{u-\omega^{r} z_{i}}-\sum_{j=1}^{m} \frac{\left\langle\alpha_{s}^{\vee}, L_{\sigma}^{r} \alpha_{c(j)}\right\rangle}{u-\omega^{r} w_{j}}\right)+\frac{\left\langle\alpha_{s}^{\vee}, T \lambda_{0}+\Lambda_{0}\right\rangle}{u} .
$$

It is known that $\rho(\mathcal{Z}(\widehat{\mathfrak{g}})) \subset \pi_{0}$ [FFR94]. That means $\rho(Z)$ is a linear combination of terms of the form $b_{s_{1}}\left[-n_{1}\right] \ldots b_{s_{M}}\left[-n_{M}\right]| \rangle, s_{1}, \ldots, s_{M} \in I, n_{1}, \ldots, n_{M} \in \mathbb{Z}_{\geq 1}$. Therefore, by repeated use of formula (4.41) one has that $\tau_{\nu(t)}\left(v, G_{c(1)}[-1]| \rangle, \ldots, G_{c(m)}[-1]| \rangle, \rho(Z)\right)$ is proportional to $\tau_{\nu(t)}\left(v, G_{c(1)}[-1]| \rangle, \ldots, G_{c(m)}[-1]| \rangle,| \rangle\right)$. Together with (4.40), this establishes that $\psi$ is an eigenvector of $\varphi(Z(u))$, as required.

In particular, consider the quadratic singular vector $S \in \mathfrak{J}(\widehat{\mathfrak{g}})$ from (2.23). It is known (see e.g., [Fre07, §8.1.4]) that

$$
\rho(S)=\frac{1}{2} \sum_{s, t \in I}\left\langle\omega_{s}, \omega_{t}\right\rangle b_{s}[-1] b_{t}[-1]| \rangle-\sum_{\alpha \in \Delta^{+}} \frac{\langle\alpha, \alpha\rangle}{4} b_{\alpha}[-2]| \rangle .
$$

where $\omega_{s}, s \in I$, are the fundamental weights. Hence

$$
\begin{align*}
\varphi(S(u)) \cdot \psi & =\left(\frac{1}{2} \sum_{s, t \in I}\left\langle\omega_{s}, \omega_{t}\right\rangle\left\langle\alpha_{s}^{\vee}, \nu(u)\right\rangle\left\langle\alpha_{t}^{\vee}, \nu(u)\right\rangle-\sum_{\alpha \in \Delta^{+}} \frac{\langle\alpha, \alpha\rangle}{4}\left\langle\alpha^{\vee}, \nu^{\prime}(u)\right\rangle\right) \psi \\
& =\left(\frac{1}{2}\langle\nu(u), \nu(u)\rangle-\left\langle\nu^{\prime}(u), \varrho\right\rangle\right) \psi \tag{4.42}
\end{align*}
$$

where $\varrho=\frac{1}{2} \sum_{\alpha \in \Delta^{+}} \alpha$. Taking the residue in $u-z_{i}$ (which comes only from the term $\left.\langle\nu(u), \nu(u)\rangle\right)$ we obtain the eigenvalue of the quadratic Hamiltonians $\mathcal{H}_{i, 0}$ given in Theorem 3.2. (In this calculation
one needs the following steps

$$
\begin{aligned}
\sum_{s=1}^{T-1} \frac{\left\langle\lambda_{k}, L_{\sigma}^{s} \lambda_{k}\right\rangle}{z_{k}\left(1-\omega^{s}\right)} & =\frac{1}{2} \sum_{s=1}^{T-1}\left(\frac{\left\langle\lambda_{k}, L_{\sigma}^{s} \lambda_{k}\right\rangle}{z_{k}\left(1-\omega^{s}\right)}+\frac{\left\langle L_{\sigma}^{-s} \lambda_{k}, \lambda_{k}\right\rangle \omega^{-s}}{z_{k}\left(\omega^{-s}-1\right)}\right) \\
& =\frac{1}{2} \sum_{s=1}^{T-1}\left(\frac{\left\langle\lambda_{k}, L_{\sigma}^{s} \lambda_{k}\right\rangle}{z_{k}\left(1-\omega^{s}\right)}+\frac{\left\langle L_{\sigma}^{s} \lambda_{k}, \lambda_{k}\right\rangle \omega^{s}}{z_{k}\left(\omega^{s}-1\right)}\right)=\frac{1}{2} \sum_{s=1}^{T-1} \frac{\left\langle\lambda_{k}, L_{\sigma}^{s} \lambda_{k}\right\rangle}{z_{k}}
\end{aligned}
$$

where the change of variable $s \mapsto-s$ was performed in the second sum.) This completes the proof of Theorem 3.2.
4.8. Proof of Theorem 3.4. We keep the notations of the previous section. Let us now suppose that $\chi=0$. Then we have from (4.25) that

$$
\begin{aligned}
\iota_{t^{-1}} \nu(t) & =0+\frac{T}{t}\left(\sum_{i=1}^{N} \Pi_{0} \lambda_{i}-\sum_{j=1}^{m} \Pi_{0} \alpha_{c(j)}+\lambda_{0}+\Lambda_{0} / T\right)+\mathcal{O}\left(\frac{1}{t^{2}}\right) \\
& =0+\frac{T}{t}\left(\lambda_{\infty}+\Lambda_{0} / T\right)+\mathcal{O}\left(\frac{1}{t^{2}}\right)
\end{aligned}
$$

where $\lambda_{\infty}$ is as in (3.4).
Now we modify the construction of $\S 2.4$ as follows. Let $\mathcal{M}_{\infty}$ be a module over $\left(\mathfrak{g}^{\mathrm{op}} \otimes \mathbb{C}\left[\left[t^{-1}\right]\right]\right)^{\Gamma}$ (rather than $\left(\mathfrak{g}^{\text {คР }} \otimes t^{-1} \mathbb{C}\left[\left[t^{-1}\right]\right]\right)^{\Gamma}$ as in $\left.\S 2.4\right)$. We make it into a module $\mathcal{M}_{\infty}^{k / T}$ over $\left(\mathfrak{g}^{\text {คค }} \otimes \mathbb{C}\left[\left[t^{-1}\right]\right]\right)^{\Gamma} \oplus$ $\mathbb{C} K_{\infty}$ by declaring that $K_{\infty}$ acts by multiplication by $k / T \in \mathbb{C}$. We have the induced right module $U\left(\hat{\mathfrak{g}}_{\infty}^{\Gamma}\right)$ module of level $k / T$,

$$
\begin{equation*}
\mathbb{M}_{\infty}^{k / T}:=\mathcal{M}_{\infty}^{k / T} \otimes_{\left.U\left(\left(\mathfrak{g} \otimes \mathbb{C}\left[t^{-1}\right]\right]\right)^{\Gamma} \oplus \mathbb{C} K_{\infty}\right)} U\left(\hat{\mathfrak{g}}_{\infty}^{\Gamma}\right) \tag{4.43}
\end{equation*}
$$

The tensor product

$$
\begin{equation*}
\mathbb{M}:=\mathbb{M}_{\infty}^{k / T} \otimes \bigotimes_{i=1}^{N} \mathbb{M}_{z_{i}}^{k} \otimes \mathbb{M}_{0}^{k / T} \tag{4.44}
\end{equation*}
$$

is again a module over $\widehat{\mathfrak{g}}_{\infty, N, 0}$ on which $K$ acts as $k$. Pulling back by the embedding (2.8), we have that $\mathbb{M}$ becomes a module over $\mathfrak{g}_{\infty, \boldsymbol{z}, 0}^{\Gamma}$ and we can form the space of coinvariants $\mathbb{M} / \mathfrak{g}_{\infty, \boldsymbol{z}, 0}^{\Gamma}$. Let us write

$$
\begin{equation*}
\widehat{\mathfrak{g}}_{\infty, N, 0}^{+}:=\left(\mathfrak{g}^{\mathbf{o p}} \otimes \mathbb{C}\left[\left[t^{-1}\right]\right]\right)^{\Gamma} \oplus \bigoplus_{i=1}^{N} \mathfrak{g} \otimes \mathbb{C}\left[\left[t-z_{i}\right]\right] \oplus(\mathfrak{g} \otimes \mathbb{C}[[t]])^{\Gamma} \tag{4.45}
\end{equation*}
$$

We have the natural inclusion $\mathfrak{g}^{\sigma} \hookrightarrow \mathfrak{g}$ and hence the embedding

$$
\begin{equation*}
\mathfrak{g}^{\sigma} \longleftrightarrow \widehat{\mathfrak{g}}_{\infty, N, 0}^{+} ; \quad X \longmapsto\left(-X[0]_{\infty} ; X[0]_{z_{1}}, \ldots, X[0]_{z_{N}} ; X[0]_{0}\right) . \tag{4.46}
\end{equation*}
$$

Pulling back by this embedding, the tensor product

$$
\begin{equation*}
\mathcal{M}:=\mathcal{M}_{\infty} \otimes \bigotimes_{i=1}^{N} \mathcal{M}_{z_{i}} \otimes \mathcal{M}_{0} \tag{4.47}
\end{equation*}
$$

is a module over $\mathfrak{g}^{\sigma}$ and we have the space of coinvariants $\mathcal{M} / \mathfrak{g}^{\sigma}$.
Proposition 4.12. There is a canonical isomorphism of vector spaces

$$
\mathbb{M} / \mathfrak{g}_{\infty, z, 0}^{\Gamma} \cong_{\mathbb{C}} \mathcal{M} / \mathfrak{g}^{\sigma}
$$

Proof. We have

$$
\mathbb{M}=U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}\right) \otimes_{U\left(\mathfrak{g}_{\infty, N, 0}^{+}\right)} \mathcal{M}
$$

Here we regard $\mathcal{M}_{\infty}$ as a left module over $U\left(\mathfrak{g}^{\sigma, o p}\right)$.
There are natural embeddings of Lie algebras $\mathfrak{g}^{\sigma} \hookrightarrow \widehat{\mathfrak{g}}_{\infty, N, 0}^{+}$and $\mathfrak{g}^{\sigma} \hookrightarrow \mathfrak{g}_{\infty, z, 0}^{\Gamma}$; the first is given in (4.46); for the second, $X \in \mathfrak{g}^{\sigma}$ embeds as the constant function $X(t)=X$ in $\mathfrak{g}_{\infty, z, 0}^{\Gamma}$. In turn, both $\widehat{\mathfrak{g}}_{\infty, N, 0}^{+}$and $\mathfrak{g}_{\infty, z, 0}^{\Gamma}$ embed into $\widehat{\mathfrak{g}}_{\infty, N, 0}$ - see (2.8) - and the following diagram of embeddings commutes


We may identify $\mathfrak{g}^{\sigma}, \widehat{\mathfrak{g}}_{\infty, N, 0}^{+}$and $\mathfrak{g}_{\infty, z, 0}^{\Gamma}$ with their images in $\widehat{\mathfrak{g}}_{\infty, N, 0}$. Then

$$
\widehat{\mathfrak{g}}_{\infty, N, 0}=\widehat{\mathfrak{g}}_{\infty, N, 0}^{+}+\mathfrak{g}_{\infty, \boldsymbol{z}, 0}^{\Gamma} \quad \text { and } \quad \hat{\mathfrak{g}}_{\infty, N, 0}^{+} \cap \mathfrak{g}_{\infty, \boldsymbol{z}, 0}^{\Gamma}=\mathfrak{g}^{\sigma} .
$$

Therefore [Di74, Proposition 2.2.9] there is an isomorphism

$$
U\left(\widehat{\mathfrak{g}}_{\infty, N, 0}\right) \cong U\left(\mathfrak{g}_{\infty, z, 0}^{\Gamma}\right) \otimes_{U\left(\mathfrak{g}^{\sigma}\right)} U\left(\hat{\mathfrak{g}}_{\infty, N, 0}^{+}\right)
$$

of vector spaces and in fact of left $U\left(\mathfrak{g}_{\infty, z, 0}^{\Gamma}\right)$-modules. Hence we have an isomorphism of left $U\left(\mathfrak{g}_{\infty, z, 0}^{\Gamma}\right)$-modules,

$$
\mathbb{M} \cong U\left(\mathfrak{g}_{\infty, z, 0}^{\Gamma}\right) \otimes_{U\left(\mathfrak{g}^{\sigma}\right)} U\left(\hat{\mathfrak{g}}_{\infty, N, 0}^{+}\right) \otimes_{U\left(\mathfrak{g}_{\infty, N, 0}^{+}\right)} \mathcal{M}=U\left(\mathfrak{g}_{\infty, z, 0}^{\Gamma}\right) \otimes_{U\left(\mathfrak{g}^{\sigma}\right)} \mathcal{M}
$$

The result follows by the following elementary lemma.
Lemma 4.13. Suppose $\mathfrak{b} \hookrightarrow \mathfrak{a}$ is an embedding of complex Lie algebras. Let $V$ be a $\mathfrak{b}$-module and $M$ an $\mathfrak{a}$-module. Then $\left(\left(U(\mathfrak{a}) \otimes_{U(\mathfrak{b})} V\right) \otimes M\right) / \mathfrak{a} \cong_{\mathbb{C}}(V \otimes M) / \mathfrak{b}$.

In particular (taking $M$ to be the trivial one-dimensional module), $\left(U(\mathfrak{a}) \otimes_{U(\mathfrak{b})} V\right) / \mathfrak{a} \cong_{\mathbb{C}} V / \mathfrak{b}$.
Proof. Consider the linear map $V \otimes M \cong_{\mathbb{C}}\left(U(\mathfrak{b}) \otimes_{U(\mathfrak{b})} V\right) \otimes M \hookrightarrow\left(U(\mathfrak{a}) \otimes_{U(\mathfrak{b})} V\right) \otimes M \rightarrow$ $\left(\left(U(\mathfrak{a}) \otimes_{U(\mathfrak{b})} V\right) \otimes M\right) / \mathfrak{a}$ sending $v \otimes m \mapsto[(1 \otimes v) \otimes m]$. This map is surjective. We must show that it has kernel $\mathfrak{b} .(V \otimes M)$. And indeed, the kernel is the intersection of $\left(U(\mathfrak{b}) \otimes_{U(\mathfrak{b})} V\right) \otimes M$ with $\mathfrak{a} .\left(\left(U(\mathfrak{a}) \otimes_{U(\mathfrak{b})} V\right) \otimes M\right)$, which is $\mathfrak{b} .\left(\left(U(\mathfrak{b}) \otimes_{U(\mathfrak{b})} V\right) \otimes M\right) \cong_{\mathbb{C}} \mathfrak{b} .(V \otimes M)$.

We may now choose modules - cf. (4.28) -

$$
\begin{align*}
\mathcal{M}_{z_{i}} & =M_{\lambda_{i}}^{*} \cong \widetilde{W}_{\iota_{t-z_{i}} \nu(t)}, \quad i=1, \ldots, N, \\
\mathcal{M}_{0} & =M_{\lambda_{0}}^{*, \sigma} \cong \widetilde{W}_{t_{t} \nu(t)}^{\Gamma}, \\
\mathcal{M}_{\infty} & =M_{\lambda_{\infty}}^{\sigma, \varphi} \cong \widetilde{W}_{-\iota_{t-1} \nu(t)}^{\Gamma, V}, \tag{4.48}
\end{align*}
$$

where the isomorphisms are as in Lemma 4.2, Lemma 4.3 and, now, Lemma 4.4. The embedding of $\mathfrak{g}$-modules $M_{\lambda_{i}}^{*} \cong \widetilde{W}_{\iota t-z_{i} \nu(t)} \hookrightarrow W_{\iota t-z_{i} \nu} \nu(t)$ extends to an embedding of modules over $\mathfrak{g} \otimes \mathbb{C}\left[\left[t-z_{i}\right]\right] \oplus \mathbb{C} K$ of level $-h^{\vee}$. (Positive modes $A[n]_{z_{i}}, n>0, A \in \mathfrak{g}$, act as zero on $\widetilde{W}_{\iota t-z_{i}} \nu(t)$ as in Lemma 4.2.) $\mathbb{M}_{z_{i}}^{-h^{\vee}}$ is by definition the $\widehat{\mathfrak{g}}_{x_{i}}$-module of level $-h^{\vee}$ induced from $\mathcal{M}_{z_{i}}=M_{\lambda_{i}}^{*}$. So we get a canonical homomorphism $\mathbb{M}_{z_{i}}^{-h^{\vee}} \rightarrow W_{\iota t-z_{i}} \nu(t)$ of $\widehat{\mathfrak{g}}_{z_{i}}$-modules of level $-h^{\vee}$. The marked points 0 and $\infty$ work
similarly. In total we get a canonical homomorphism of $\widehat{\mathfrak{g}}_{\infty, p, 0}$-modules

$$
\mathbb{M}_{\infty}^{-h^{\vee} / T} \otimes \bigotimes_{i=1}^{N} \mathbb{M}_{z_{i}}^{-h^{\vee}} \otimes \mathbb{M}_{0}^{-h^{\vee} / T} \rightarrow W_{-\iota_{t}-1 \nu(t)}^{\Gamma, \vee} \otimes \bigotimes_{i=1}^{N} W_{\iota t-z_{i} \nu(t)} \otimes W_{\iota t \nu(t)}^{\Gamma}
$$

Hence, from the $\mathfrak{g}_{\infty, z, 0}^{\Gamma}$-invariant functional $\psi_{\nu(t)}$ of (4.27), we obtain a $\mathfrak{g}_{\infty, z, 0}^{\Gamma}$-invariant functional $\mathbb{M}_{\infty}^{-h^{\vee} / T} \otimes \bigotimes_{i=1}^{N} \mathbb{M}_{z_{i}}^{-h^{\vee}} \otimes \mathbb{M}_{0}^{-h^{\vee} / T} \rightarrow \mathbb{C}$. By Proposition 4.12 this is the same thing as a $\mathfrak{g}^{\sigma}$-invariant functional

$$
\begin{equation*}
M_{\lambda_{\infty}}^{\sigma, \varphi} \otimes \bigotimes_{i=1}^{N} M_{\lambda_{i}}^{*} \otimes M_{\lambda_{0}}^{*, \sigma} \rightarrow \mathbb{C} \tag{4.49}
\end{equation*}
$$

Recall that we are embedding $\mathfrak{g}^{\sigma} \hookrightarrow \mathfrak{g}^{\sigma, o p} \oplus \bigoplus_{i=1}^{N} \mathfrak{g} \oplus \mathfrak{g}^{\sigma}$ by $X \mapsto(-X, X, \ldots, X, X)$ and that the antipode map $S: U\left(\mathfrak{g}^{\sigma}\right) \rightarrow U\left(\mathfrak{g}^{\sigma}\right)$ is defined by $S(X)=-X$ for $X \in \mathfrak{g}^{\sigma}$. So the space of $\mathfrak{g}^{\sigma}$-invariant functionals $M_{\lambda_{\infty}}^{\sigma, \varphi} \otimes \bigotimes_{i=1}^{N} M_{\lambda_{i}}^{*} \otimes M_{\lambda_{0}}^{*, \sigma} \rightarrow \mathbb{C}$ is the space

$$
\operatorname{Hom}_{\mathfrak{g}^{\sigma}}\left(\left(M_{\lambda_{\infty}}^{\sigma}\right)^{\varphi \circ S} \otimes \bigotimes_{i=1}^{N} M_{\lambda_{i}}^{*} \otimes M_{\lambda_{0}}^{*, \sigma}, \mathbb{C}\right)
$$

Proposition 4.14. There is an isomorphism of vector spaces, and in fact of left $\left(U(\mathfrak{g})^{\otimes N} \otimes\right.$ $\left.U\left(\mathfrak{g}^{\sigma}\right)\right)^{\mathfrak{g}^{\sigma}}$-modules,

$$
\begin{equation*}
\left(\bigotimes_{i=1}^{N} M_{\lambda_{i}} \otimes M_{\lambda_{0}}^{\sigma}\right)_{\lambda_{\infty}}^{\mathfrak{n}^{\sigma}} \cong \operatorname{Hom}_{\mathfrak{g}^{\sigma}}\left(\left(M_{\lambda_{\infty}}^{\sigma}\right)^{\varphi \circ S} \otimes \bigotimes_{i=1}^{N} M_{\lambda_{i}}^{*} \otimes M_{\lambda_{0}}^{*, \sigma}, \mathbb{C}\right) \tag{4.50}
\end{equation*}
$$

Proof. We have

$$
\begin{align*}
\left(\bigotimes_{i=1}^{N} M_{\lambda_{i}} \otimes M_{\lambda_{0}}^{\sigma}\right)_{\lambda_{\infty}}^{\mathfrak{n}^{\sigma}} & =\operatorname{Hom}_{\mathfrak{h}^{\sigma} \oplus \mathfrak{n}^{\sigma}}\left(\mathbb{C} v_{\lambda_{\infty}}, \bigotimes_{i=1}^{N} M_{\lambda_{i}} \otimes M_{\lambda_{0}}^{\sigma}\right) \\
& \cong \operatorname{Hom}_{\mathfrak{g}^{\sigma}}\left(U\left(\mathfrak{g}^{\sigma}\right) \otimes_{U\left(\mathfrak{h}^{\sigma} \oplus \mathfrak{n}^{\sigma}\right)} \mathbb{C} v_{\lambda_{\infty}}, \bigotimes_{i=1}^{N} M_{\lambda_{i}} \otimes M_{\lambda_{0}}^{\sigma}\right) \\
& =\operatorname{Hom}_{\mathfrak{g}^{\sigma}}\left(M_{\lambda_{\infty}}^{\sigma}, \bigotimes_{i=1}^{N} M_{\lambda_{i}} \otimes M_{\lambda_{0}}^{\sigma}\right) \tag{4.51}
\end{align*}
$$

This is the space of those maps $\phi: M_{\lambda_{\infty}}^{\sigma} \rightarrow \bigotimes_{i=1}^{N} M_{\lambda_{i}} \otimes M_{\lambda_{0}}^{\sigma}$ such that

$$
\begin{equation*}
\phi(x \cdot v)=\left(\Delta^{N+1} x\right) \cdot \phi(v), \quad v \in M_{\lambda_{\infty}}^{\sigma}, x \in \mathfrak{g}^{\sigma} . \tag{4.52}
\end{equation*}
$$

It carries a natural left action of $\left(U(\mathfrak{g})^{\otimes N} \otimes U\left(\mathfrak{g}^{\sigma}\right)\right)^{\mathfrak{g}^{\sigma}}$ : given one such a map $\phi$, the map $v \mapsto X \phi(v)$ is another, for any $X \in\left(U(\mathfrak{g})^{\otimes N} \otimes U\left(\mathfrak{g}^{\sigma}\right)\right)^{\mathfrak{g}^{\sigma}}$.

Then

$$
\begin{equation*}
\operatorname{Hom}_{\mathfrak{g}^{\sigma}}\left(M_{\lambda_{\infty}}^{\sigma}, \bigotimes_{i=1}^{N} M_{\lambda_{i}} \otimes M_{\lambda_{0}}^{\sigma}\right) \cong \operatorname{Hom}_{\mathfrak{g}^{\sigma}}\left(M_{\lambda_{\infty}} \otimes\left(\bigotimes_{i=1}^{N} M_{\lambda_{i}} \otimes M_{\lambda_{0}}^{\sigma}\right)^{\vee, S}, \mathbb{C}\right) \tag{4.53}
\end{equation*}
$$

where for a finitely-weighted left $U\left(\mathfrak{g}^{\sigma}\right)$-module $M$ we denote by $M^{\vee, S}$ its restricted dual made into a left $U\left(\mathfrak{g}^{\sigma}\right)$-module by twisting by the antipode map $S$. Indeed, given a map $\phi$ such that (4.52)
holds we define a map $\Phi: M_{\lambda_{\infty}} \otimes\left(\otimes_{i=1}^{N} M_{\lambda_{i}} \otimes M_{\lambda_{0}}^{\sigma}\right)^{\vee, S} \rightarrow \mathbb{C}$ by $\Phi(v \otimes \mu)=\mu(\phi(v)), v \in M_{\lambda_{\infty}}, \mu \in$ $\left(\otimes_{i=1}^{N} M_{\lambda_{i}} \otimes M_{\lambda_{0}}^{\sigma}\right)^{\vee, S}$. It obeys $\Phi(x .(v \otimes \mu))=0$ where $x .(v \otimes \mu):=(x . v) \otimes \mu+v \otimes\left(\mu \circ\left(-\Delta^{N+1} x.\right)\right)$, for $x \in \mathfrak{g}^{\sigma}$.

Finally, we may twist the action of $\mathfrak{g}^{\sigma}$ on $M_{\lambda_{\infty}} \otimes\left(\otimes_{i=1}^{N} M_{\lambda_{i}} \otimes M_{\lambda_{0}}^{\sigma}\right)^{\vee, S}$ by the automorphism $S \circ \varphi=\varphi \circ S$. Since $S^{2}=\operatorname{id}$ here, $\left(M_{\lambda}^{\vee, S}\right)^{S \circ \varphi}=M_{\lambda}^{\vee, \varphi}=: M_{\lambda}^{*}$. Thus

$$
\begin{equation*}
\operatorname{Hom}_{\mathfrak{g}^{\sigma}}\left(M_{\lambda_{\infty}}^{\sigma} \otimes\left(\bigotimes_{i=1}^{N} M_{\lambda_{i}} \otimes M_{\lambda_{0}}^{\sigma}\right)^{\vee, S}, \mathbb{C}\right) \cong \operatorname{Hom}_{\mathfrak{g}^{\sigma}}\left(\left(M_{\lambda_{\infty}}^{\sigma}\right)^{\varphi \circ S} \otimes \bigotimes_{i=1}^{N} M_{\lambda_{i}}^{*} \otimes M_{\lambda_{0}}^{*, \sigma}, \mathbb{C}\right) \tag{4.54}
\end{equation*}
$$

as required.

Thus, the functional of (4.49) defines an element of $\left(\otimes_{i=1}^{N} M_{\lambda_{i}} \otimes M_{\lambda_{0}}^{\sigma}\right)_{\lambda_{\infty}}^{\mathfrak{n}^{\sigma}}$. At the same time, by (4.51), this functional is uniquely determined by its restriction to the subspace

$$
\mathbb{C} v_{\lambda_{\infty}} \otimes \bigotimes_{i=1}^{N} M_{\lambda_{i}}^{*} \otimes M_{\lambda_{0}}^{*, \sigma} \cong \mathbb{C}\langle | \otimes \bigotimes_{i=1}^{N} \widetilde{W}_{\iota t-z_{i}} \nu(t) \otimes \widetilde{W}_{\iota t \nu(t)}^{\Gamma} .
$$

By definition, on this subspace it agrees with the functional of (4.29). In other words, it defines the same vector $\psi \in \bigotimes_{i=1}^{N} M_{\lambda_{i}} \otimes M_{\lambda_{0}}^{\sigma}$ as in (4.30). We saw in the previous section that $\psi$ vector of (4.30) is the weight function defined in (3.3). So we have shown that the weight function $\psi$ is singular when $\chi=0$, completing the proof of Theorem 3.4.

## Appendix A. Proof of Proposition 2.5

Let us write

$$
\partial_{x}^{(n)}:=\frac{1}{n!}\left(\frac{\partial}{\partial x}\right)^{n}
$$

Consider the $\Gamma$-equivariant rational function

$$
\begin{equation*}
f(t)=\sum_{k=0}^{T-1} \frac{\sigma^{k} A}{\left(\omega^{-k} t-u\right)^{p+1}} \in \mathfrak{g}_{\infty, z, u, 0}^{\Gamma} \tag{A.1}
\end{equation*}
$$

The expansion of (A.1) at $u$ is

$$
\iota_{t-u} f(t)=A[-p-1]_{u}+\sum_{k=1}^{T-1} \sum_{n=0}^{\infty}\left[\begin{array}{c}
p+n  \tag{A.2}\\
n
\end{array}\right] \frac{(-1)^{p+1} \omega^{k(p+1)}}{\left(\omega^{k}-1\right)^{n+p+1} u^{n+p+1}}\left(\sigma^{k} A\right)[n]_{u} \in \mathfrak{g} \otimes \mathbb{C}((t-u)) .
$$

Its expansions at $z_{i}, i=1, \ldots, N$, and at 0 are given by

$$
\begin{gather*}
\iota_{t-z_{i}} f(t)=-\partial_{u}^{(p)} \sum_{k=0}^{T-1} \sum_{n=0}^{\infty} \frac{\omega^{-k n}}{\left(u-\omega^{-k} z_{i}\right)^{n+1}}\left(\sigma^{k} A\right)[n]_{z_{i}} \in \mathfrak{g} \otimes \mathbb{C}\left[\left[t-z_{i}\right]\right],  \tag{A.3}\\
\iota_{t} f(t)=-T \partial_{u}^{(p)} \sum_{n=0}^{\infty} \frac{1}{u^{n+1}}\left(\Pi_{n} A\right)[n]_{0} \in \mathfrak{g} \otimes \mathbb{C}[[t]], \tag{A.4}
\end{gather*}
$$

where, recall, $A[n]_{z_{i}}=A \otimes\left(t-z_{i}\right)^{n} \in \mathfrak{g} \otimes \mathbb{C}\left(\left(t-z_{i}\right)\right)$. It is regular at $\infty$ and its expansion there is given by

$$
\begin{equation*}
-\iota_{t^{-1}} f(t)=-\partial_{u}^{(p)} \sum_{k=0}^{T} \sum_{n=0}^{\infty}\left(\omega^{k} u\right)^{n} \omega^{k} \sigma^{k} A[-n-1]_{\infty}=-T \partial_{u}^{(p)} \sum_{n=0}^{\infty} u^{n}\left(\Pi_{-n-1} A\right)[-n-1]_{\infty}, \tag{A.5}
\end{equation*}
$$

where recall that for the modes at infinity we write $A[n]_{\infty}=A \otimes t^{n}$ in our conventions.
Thus, since $\left[f(t) .\left(x \otimes v_{0}\right)\right]=0$, we have

$$
\begin{align*}
& {\left[S(u) x \otimes v_{0}\right]:=\left[x \otimes \frac{1}{2} I^{a}[-1] I_{a}[-1] v_{0}\right]=\left[I^{a}(u) x \otimes \frac{1}{2} I_{a}[-1] v_{0}\right]} \\
& \quad+\sum_{p=1}^{T-1} \frac{\omega^{p}}{\left(\omega^{p}-1\right) u}\left[x \otimes \frac{1}{2}\left(\sigma^{p} I^{a}\right)[0] I_{a}[-1] v_{0}\right]+\sum_{p=1}^{T-1} \frac{\omega^{p}}{\left(\omega^{p}-1\right)^{2} u^{2}}\left[x \otimes \frac{1}{2}\left(\sigma^{p} I^{a}\right)[1] I_{a}[-1] v_{0}\right] \tag{A.6}
\end{align*}
$$

where, putting the expansions at $0, z_{1}, \ldots, z_{N}, \infty$ together, we shall write, for any $A \in \mathfrak{g}$,

$$
\begin{equation*}
A(u):=\left(\sum_{i=1}^{N} \sum_{n=0}^{\infty} \sum_{k=0}^{T-1} \frac{\omega^{-k n}\left(\sigma^{k} A\right)[n]_{z_{i}}}{\left(u-\omega^{-k} z_{i}\right)^{n+1}}+T \sum_{n=0}^{\infty} \frac{\left(\Pi_{n} A\right)[n]_{0}}{u^{n+1}}+T \sum_{n=0}^{\infty} u^{n}\left(\Pi_{-n-1} A\right)[-n-1]_{\infty}\right) . \tag{A.7}
\end{equation*}
$$

The equality in (A.6) is an example of what was called "cyclotomic swapping" in [VY16a]. Then, defining the element $F \in \mathfrak{g}^{\sigma}$ and number $K$ as in (2.22), we see after one further such "swapping" that

$$
\begin{equation*}
S(u)=\frac{1}{2} I_{a}(u) I^{a}(u)+\frac{1}{u} F(u)+\frac{1}{u^{2}} K . \tag{A.8}
\end{equation*}
$$

Consider first the terms in $\frac{1}{2} I_{a}(u) I^{a}(u)$ acting at sites $i, j \in\{1, \ldots, N\}, i \neq j$. Let us write

$$
\begin{equation*}
z_{i k}:=\omega^{-k} z_{i} \quad \text { and } \quad \partial_{i k}^{(n)}=\frac{1}{n!}\left(\frac{\partial}{\partial z_{i k}}\right)^{n} \tag{A.9}
\end{equation*}
$$

Then $\frac{1}{\left(u-z_{i k}\right)^{n+1}}=\partial_{i k}^{(n)} \frac{1}{u-z_{i k}}$ and we get the terms

$$
\begin{gather*}
\frac{1}{2} \sum_{n, m=0}^{\infty} \sum_{\substack{i, j=1 \\
j \neq i}}^{N} \sum_{k, l=0}^{T-1} \partial_{i k}^{(n)} \partial_{j l}^{(m)} \frac{1}{u-z_{i k}} \frac{1}{u-z_{j l}} \omega^{-k n-l m}\left(\sigma^{k} I_{a}\right)[n]_{z_{i}}\left(\sigma^{l} I^{a}\right)[m]_{z_{j}} \\
=\frac{1}{n, m=0} \sum_{\substack{i, j=1 \\
j \neq i}}^{N} \sum_{k, l=0}^{T-1} \partial_{i k}^{(n)} \partial_{j l}^{(m)}\left(\frac{1}{u-z_{i k}} \frac{1}{z_{i k}-z_{j l}}+\frac{1}{u-z_{j l}} \frac{1}{z_{j l}-z_{i k}}\right) \omega^{-k n-l m}\left(\sigma^{k} I_{a}\right)[n]_{z_{i}}\left(\sigma^{l} I^{a}\right)[m]_{z_{j}} \\
=\frac{1}{2} \sum_{n, m=0}^{\infty} \sum_{\substack{i, j=1 \\
j \neq i}}^{N} \sum_{k, l=0}^{T-1} \partial_{i k}^{(n)} \partial_{j l}^{(m)} \frac{1}{u-z_{i k}} \frac{1}{z_{i k}-z_{j l}} \omega^{-k n-l m}\left\{\left(\sigma^{k} I_{a}\right)[n]_{z_{i}},\left(\sigma^{l} I^{a}\right)[m]_{z_{j}}\right\} \\
\left.=\sum_{n, m=0}^{\infty} \sum_{\substack{i, j=1 \\
j \neq i}}^{N} \sum_{k, l=0}^{T-1} \partial_{i k}^{(n)} \partial_{j l}^{(m)} \frac{1}{u-z_{i k}} \frac{1}{z_{i k}-z_{j l}} \omega^{-k n-l m}\left(\sigma^{k} I_{a}\right)[n]_{z_{i}}\left(\sigma^{l} I^{a}\right)[m]_{z_{j}} \quad \text { (A.10) }\right) \tag{A.10}
\end{gather*}
$$

where $\{a, b\}:=a b+b a$ denotes the anti-commutator. Now here

$$
\begin{align*}
\partial_{i k}^{(n)} \partial_{j l}^{(m)} \frac{1}{u-z_{i k}} \frac{1}{z_{i k}-z_{j l}}= & \sum_{p=0}^{n}\left(\partial_{i k}^{(p)} \frac{1}{u-z_{i k}}\right)\left(\partial_{i k}^{(n-p)} \partial_{j l}^{(m)} \frac{1}{z_{i k}-z_{j l}}\right) \\
& =\sum_{p=0}^{n} \frac{1}{\left(u-z_{i k}\right)^{p+1}}(-1)^{n-p}\left[\begin{array}{c}
n+m-p \\
m
\end{array}\right] \frac{1}{\left(z_{i k}-z_{j l}\right)^{n-p+m+1}} \tag{A.11}
\end{align*}
$$

Thus we find the following terms in $\frac{1}{2} I_{a}(u) I^{a}(u)$ :

$$
\begin{align*}
& \sum_{n, m=0}^{\infty} \sum_{\substack{i, j=1 \\
j \neq i}}^{N} \sum_{k, l=0}^{T-1} \sum_{p=0}^{n} \frac{\omega^{-k n-l m}}{\left(u-z_{i k}\right)^{p+1}}(-1)^{n-p}\left[\begin{array}{c}
n+m-p \\
m
\end{array}\right] \frac{1}{\left(z_{i k}-z_{j l}\right)^{n-p+m+1}}\left(\sigma^{k} I_{a}\right)[n]_{z_{i}}\left(\sigma^{l} I^{a}\right)[m]_{z_{j}} \\
= & \sum_{i=1}^{N} \sum_{k=0}^{T-1} \sum_{p=0}^{\infty} \frac{1}{\left(u-\omega^{-k} z_{i}\right)^{p+1}}\left(\sum_{\substack{j=1 \\
j \neq i}}^{N} \sum_{l=0}^{T-1} \sum_{r, m=0}^{\infty} \frac{\omega^{-k(r+p)-l m}(-1)^{r}\left[\begin{array}{c}
r+m \\
m
\end{array}\right]}{\left(\omega^{-k} z_{i}-\omega^{-l} z_{j}\right)^{r+m+1}}\left(\sigma^{k} I_{a}\right)[r+p]_{z_{i}}\left(\sigma^{l} I^{a}\right)[m]_{z_{j}}\right) \\
= & \sum_{i=1}^{N} \sum_{k=0}^{T-1} \sum_{p=0}^{\infty} \frac{\omega^{-k p+k} \sigma^{k}}{\left(u-\omega^{-k} z_{i}\right)^{p+1}}\left(\sum_{\substack{j=1 \\
j \neq i}}^{N} \sum_{l=0}^{T-1} \sum_{r, m=0}^{\infty} \frac{\omega^{(k-l) m}(-1)^{r}\left[\begin{array}{r}
r+m \\
m
\end{array}\right]}{\left(z_{i}-\omega^{k-l} z_{j}\right)^{r+m+1}} I_{a}[r+p]_{z_{i}}\left(\sigma^{l-k} I^{a}\right)[m]_{z_{j}}\right) \\
= & \sum_{i=1}^{N} \sum_{k=0}^{T-1} \sum_{p=0}^{\infty} \frac{\omega^{-k p+k} \sigma^{k}}{\left(u-\omega^{-k} z_{i}\right)^{p+1}}\left(\sum_{\substack{j=1 \\
j \neq i}}^{N} \sum_{l=0}^{T-1} \sum_{r, m=0}^{\infty} \frac{\omega^{-l m}(-1)^{r}\left[\begin{array}{r}
r+m \\
m
\end{array}\right]}{\left(z_{i}-\omega^{-l} z_{j}\right)^{r+m+1}} I_{a}[r+p]_{z_{i}}\left(\sigma^{l} I^{a}\right)[m]_{z_{j}}\right) \tag{A.12}
\end{align*}
$$

In part by a similar calculation we find the terms in $\frac{1}{2} I_{a}(u) I^{a}(u)$ acting solely at one site $i \in$ $\{1, \ldots, N\}$, namely

$$
\begin{align*}
& \sum_{i=1}^{N} \sum_{k=0}^{T-1} \sum_{p=0}^{\infty} \frac{\omega^{-k p+k} \sigma^{k}}{\left(u-\omega^{-k} z_{i}\right)^{p+1}} \times \\
& \quad\left(\sum_{l=1}^{T-1} \sum_{r, m=0}^{\infty} \frac{\omega^{-l m}(-1)^{r}\left[\begin{array}{c}
r+m \\
m
\end{array}\right]}{\left(\left(1-\omega^{-l}\right) z_{i}\right)^{r+m+1}} \frac{1}{2}\left\{I_{a}[r+p]_{z_{i}},\left(\sigma^{l} I^{a}\right)[m]_{z_{i}}\right\}+\sum_{n=0}^{p-1} \frac{1}{2} I_{a}[n]_{z_{i}} I^{a}[p-n-1]_{z_{i}}\right), \tag{A.13}
\end{align*}
$$

where the second term is an "on-diagonal" term in $\frac{1}{2} I_{a}(u) I^{a}(u)$.

The cross terms in $\frac{1}{2} I_{a}(u) I^{a}(u)$ acting at sites $i \in\{1, \ldots, N\}$ and 0 are given by (here we use an obvious trick to make the calculation resemble the one above)

$$
\begin{align*}
& \left.T \sum_{n, m=0}^{\infty} \sum_{i=1}^{N} \sum_{k=0}^{T-1} \partial_{i k}^{(n)} \partial_{z}^{(m)} \frac{1}{u-z_{i k}} \frac{1}{u-z}\right|_{z=0} \omega^{-k n}\left(\sigma^{k} I_{a}\right)[n]_{z_{i}}\left(\Pi_{m} I^{a}\right)[m]_{0} \\
& =\left.T \sum_{n, m=0}^{\infty} \sum_{i=1}^{N} \sum_{k=0}^{T-1} \partial_{i k}^{(n)} \partial_{z}^{(m)}\left(\frac{1}{u-z_{i k}} \frac{1}{z_{i k}-z}+\frac{1}{u-z} \frac{1}{z-z_{i k}}\right)\right|_{z=0} \omega^{-k n}\left(\sigma^{k} I_{a}\right)[n]_{z_{i}}\left(\Pi_{m} I^{a}\right)[m]_{0} \\
& =\left.T \sum_{n, m=0}^{\infty} \sum_{i=1}^{N} \sum_{k=0}^{T-1} \sum_{p=0}^{n} \frac{1}{\left(u-z_{i k}\right)^{p+1}} \frac{(-1)^{n-p}\left[\begin{array}{l}
n+m-p \\
m
\end{array}\right.}{\left(z_{i k}-z\right)^{n-p+m+1}}\right|_{z=0} \omega^{-k n}\left(\sigma^{k} I_{a}\right)[n]_{z_{i}}\left(\Pi_{m} I^{a}\right)[m]_{0} \\
& \quad+T \sum_{n, m=0}^{\infty} \sum_{i=1}^{N} \sum_{k=0}^{T-1} \sum_{p=0}^{m} \frac{1}{(u-z)^{p+1}} \frac{(-1)^{m-p}\left[\left.\begin{array}{c}
n+m-p \\
n \\
\left(z-z_{i k}\right)^{n-p+m+1} \\
=
\end{array}\right|_{z=0} ^{\infty} \omega^{-k n}\left(\sigma^{k} I_{a}\right)[n]_{z_{i}}\left(\Pi_{m} I^{a}\right)[m]_{0}\right.}{\sum_{n, m=0}^{T-1} \sum_{i=1}^{n} \frac{1}{\left(u-\omega^{-k} z_{i}\right)^{p+1}} \frac{(-1)^{n-p}[n+m-p}{\left.z_{i}^{n-p+m+1}\right]} \omega^{k(-p+m+1)}\left(\sigma^{k} I_{a}\right)[n]_{z_{i}}\left(\Pi_{m} I^{a}\right)[m]_{0}} \\
& \quad+T \sum_{n, m=0}^{\infty} \sum_{i=1}^{N} \sum_{k=0}^{T-1} \sum_{p=0}^{m} \frac{1}{u^{p+1}} \frac{(-1)^{n+1}\left[\begin{array}{l}
n+m-p \\
n
\end{array}\right.}{z_{i}^{n-p+m+1} \omega^{k(-p+m+1)}\left(\sigma^{k} I_{a}\right)[n]_{z_{i}}\left(\Pi_{m} I^{a}\right)[m]_{0}}
\end{align*}
$$

Next, the cross terms in $\frac{1}{2} I_{a}(u) I^{a}(u)$ acting at sites $i \in\{1, \ldots, N\}$ and $\infty$ are

$$
\begin{align*}
& T \sum_{k=0}^{T-1} \sum_{n, m=0}^{\infty} \sum_{i=1}^{N} \frac{\omega^{-k n}}{\left(u-\omega^{-k} z_{i}\right)^{n+1}} u^{m}\left(\Pi_{-m-1} I_{a}\right)[-m-1]_{\infty}\left(\sigma^{k} I^{a}\right)[n]_{z_{i}} \\
& \quad=T \sum_{k=0}^{T-1} \sum_{n, m=0}^{\infty} \sum_{i=1}^{N}\left(\sum_{p=0}^{m-n-1}\left[\begin{array}{c}
n+p \\
n
\end{array}\right] u^{m-n-1-p}\left(\omega^{-k} z_{i}\right)^{p}+\sum_{p=0}^{n}\left[\begin{array}{c}
m \\
p
\end{array}\right] \frac{\left(\omega^{-k} z_{i}\right)^{m-p}}{\left(u-\omega^{-k} z_{i}\right)^{n-p+1}}\right) \times \\
& \omega^{-k n}\left(\Pi_{-m-1} I_{a}\right)[-m-1]_{\infty}\left(\sigma^{k} I^{a}\right)[n]_{z_{i}} \tag{A.15}
\end{align*}
$$

where we used the identity

$$
\frac{u^{m}}{(u-z)^{n+1}}=\sum_{p=0}^{m-n-1}\left[\begin{array}{c}
n+p \\
n
\end{array}\right] u^{m-n-1-p} z^{p}+\sum_{p=0}^{n}\left[\begin{array}{l}
m \\
p
\end{array}\right] \frac{z^{m-p}}{(u-z)^{n-p+1}} .
$$

This can be seen by writing the left hand side as

$$
\frac{u^{m}}{(u-z)^{n+1}}=\frac{1}{n!} \frac{\partial^{n}}{\partial z^{n}} \frac{u^{m}}{u-z}=\frac{1}{n!} \frac{\partial^{n}}{\partial z^{n}} \sum_{k=0}^{m-1} u^{k} z^{m-k-1}+\frac{1}{n!} \frac{\partial^{n}}{\partial z^{n}} \frac{z^{m}}{u-z} .
$$

After taking the $n$ derivatives with respect to $z$, each of the two terms on the left hand side then evaluate to the two terms on the left hand side of the above identity.

The pole term in $\left(u-\omega^{-k} z_{i}\right)$ in (A.15) is (here $\left.r=n-p, n=r+p\right)$

$$
\begin{align*}
& T \sum_{i=1}^{N} \sum_{k=0}^{T-1} \sum_{r=0}^{\infty} \frac{1}{\left(u-\omega^{-k} z_{i}\right)^{r+1}} \sum_{p, m=0}^{\infty} \omega^{-k r-k m} z_{i}^{m-p}\left[\begin{array}{c}
m \\
p
\end{array}\right]\left(\Pi_{-m-1} I_{a}\right)[-m-1]_{\infty}\left(\sigma^{k} I^{a}\right)[r+p]_{z_{i}} \\
& \quad=T \sum_{i=1}^{N} \sum_{k=0}^{T-1} \sum_{r=0}^{\infty} \frac{\omega^{-k r+k} \sigma^{k}}{\left(u-\omega^{-k} z_{i}\right)^{r+1}}\left(\sum_{p, m=0}^{\infty} z_{i}^{m-p}\left[\begin{array}{c}
m \\
p
\end{array}\right]\left(\Pi_{-m-1} I_{a}\right)[-m-1]_{\infty} I^{a}[r+p]_{z_{i}}\right) \\
& =T \sum_{i=1}^{N} \sum_{k=0}^{T-1} \sum_{r=0}^{\infty} \frac{\omega^{-k r+k} \sigma^{k}}{\left(u-\omega^{-k} z_{i}\right)^{r+1}}\left(\sum_{n, p=0}^{\infty} z_{i}^{n}\left[\begin{array}{c}
n+p \\
p
\end{array}\right]\left(\Pi_{-n-p-1} I_{a}\right)[-n-p-1]_{\infty} I^{a}[r+p]_{z_{i}}\right) \tag{A.16}
\end{align*}
$$

where we use the fact that $\sigma^{k} \Pi_{-m-1}=\omega^{-k m-k} \Pi_{-m-1}$.
The powers of $u$ in (A.15) are (here $r=m-n-1-p, p=m-n-1-r$ )

$$
\begin{align*}
& T \sum_{i=1}^{N} \sum_{k=0}^{T-1} \sum_{m, n=0}^{\infty} \sum_{r=0}^{m-n-1} u^{r} \omega^{-k m+k+k r} z_{i}^{m-n-1-r}\left[\begin{array}{c}
m-1-r \\
n
\end{array}\right]\left(\Pi_{-m-1} I_{a}\right)[-m-1]_{\infty}\left(\sigma^{k} I^{a}\right)[n]_{z_{i}} \\
= & T \sum_{r=0}^{\infty} u^{r} \sum_{i=1}^{N} \sum_{k=0}^{T-1} \sum_{n=0}^{\infty} \sum_{m=r+n+1}^{\infty} \omega^{-k m+k+k r} z_{i}^{m-n-1-r}\left[\begin{array}{c}
m-1-r \\
n
\end{array}\right]\left(\Pi_{-m-1} I_{a}\right)[-m-1]_{\infty}\left(\sigma^{k} I^{a}\right)[n]_{z_{i}} \tag{A.17}
\end{align*}
$$

The cross terms in $\frac{1}{2} I_{a}(u) I^{a}(u)$ acting at sites 0 and $\infty$ can be conveniently written as

$$
\begin{align*}
T^{2} \sum_{r=0}^{\infty} \frac{1}{u^{r+1}} \sum_{n=r}^{\infty}\left(\Pi_{n} I_{a}\right)[n]_{0}\left(\Pi_{r-n+1} I^{a}\right) & {[-n+r-1]_{\infty} } \\
+ & T^{2} \sum_{r=0}^{\infty} u^{r} \sum_{n=0}^{\infty}\left(\Pi_{n} I_{a}\right)[n]_{0}\left(\Pi_{-r-n-2} I^{a}\right)[-n-r-2]_{\infty} \tag{A.18}
\end{align*}
$$

Finally the diagonal terms in $\frac{1}{2} I_{a}(u) I^{a}(u)$ acting at sites 0 and $\infty$ respectively are

$$
\begin{equation*}
\frac{T^{2}}{2} \sum_{r=1}^{\infty} \frac{1}{u^{r+1}} \sum_{n=0}^{r-1}\left(\Pi_{n} I_{a}\right)[n]_{0}\left(\Pi_{r-n-1} I^{a}\right)[r-n-1]_{0} \tag{A.19}
\end{equation*}
$$

and

$$
\begin{equation*}
\frac{T^{2}}{2} \sum_{r=0}^{\infty} u^{r} \sum_{n=0}^{r}\left(\Pi_{-n-1} I_{a}\right)[-n-1]_{\infty}\left(\Pi_{-r+n-1} I^{a}\right)[-r+n-1]_{\infty} \tag{A.20}
\end{equation*}
$$

Next we turn to the term $F(u) / u$ in (A.8). Noting that

$$
\frac{1}{u} \frac{1}{(u-z)^{n+1}}=\sum_{p=0}^{n}(-1)^{p} \frac{1}{z^{p+1}} \frac{1}{(u-z)^{n-p+1}}-\frac{1}{u} \frac{(-1)^{n}}{z^{n+1}}
$$

we find that in $F(u) / u$ the terms acting at the sites $i \in\{1, \ldots, N\}$ are

$$
\begin{align*}
& \sum_{i=1}^{N} \sum_{k=0}^{T-1} \sum_{n=0}^{\infty}\left(\sum_{p=0}^{n}(-1)^{p} \frac{\omega^{-k n}}{\left(\omega^{-k} z_{i}\right)^{p+1}\left(u-\omega^{-k} z_{i}\right)^{n-p+1}}-\frac{1}{u} \frac{\omega^{-k n}(-1)^{n}}{\left(\omega^{-k} z_{i}\right)^{n+1}}\right) \sigma^{k} F[n]_{z_{i}} \\
& \quad=\sum_{i=1}^{N} \sum_{k=0}^{T-1} \sum_{r=0}^{\infty} \sum_{n=r}^{\infty}(-1)^{n-r} \frac{\omega^{-k n+k(n-r+1)}}{\left(u-\omega^{-k} z_{i}\right)^{r+1}} \frac{1}{z_{i}^{n-r+1}} \sigma^{k} F[n]_{z_{i}}-\frac{1}{u} \sum_{i=1}^{N} \sum_{n=0}^{\infty} \frac{(-1)^{n}}{z_{i}^{n+1}} \sum_{k=0}^{T-1} \omega^{k} \sigma^{k} F[n]_{z_{i}} \\
& =\sum_{i=1}^{N} \sum_{k=0}^{T-1} \sum_{r=0}^{\infty} \sum_{n=0}^{\infty}(-1)^{n} \frac{\omega^{-k r+k}}{\left(u-\omega^{-k} z_{i}\right)^{r+1}} \frac{1}{z_{i}^{n+1}} \sigma^{k} F[n+r]_{z_{i}}-\frac{T}{u} \sum_{i=1}^{N} \sum_{n=0}^{\infty} \frac{(-1)^{n}}{z_{i}^{n+1}} \Pi_{-1} F[n] z_{i} . \tag{A.21}
\end{align*}
$$

The terms acting at the sites 0 and $\infty$ in $F(u) / u$ are clearly

$$
T \sum_{n=0}^{\infty} \frac{1}{u^{n+2}}\left(\Pi_{n} F\right)[n]_{0}+T \sum_{n=0}^{\infty} u^{n-1}\left(\Pi_{-n-1} F\right)[-n-1]_{\infty}
$$

Now we collect terms, and use $\sigma^{k} \Pi_{m}=\omega^{k m} \Pi_{m}, \sigma^{k} \mathcal{H}_{i, p}=\mathcal{H}_{i, p}$, and the fact that $\Pi_{k} F=0$ for all $k \in \mathbb{Z} / T \mathbb{Z} \backslash\{0\}$, to obtain the result.

## Appendix B. $Y_{W}$-map

B.1. Vertex Lie algebras. In this appendix we use the notion of a vertex Lie algebra and related concepts as defined in [VY16b], to which we refer the reader for the precise definition. In particular, all our vertex Lie algebras are finitely generated as $\mathbb{C}[\mathcal{D}]$-modules, of the form

$$
\mathscr{L}:=\mathbb{C}[\mathcal{D}] \otimes\left(L^{o} \oplus \mathbb{C} \mathrm{c}\right) / \mathcal{D} \mathbb{C}[\mathcal{D}] \otimes \mathbb{C} \mathrm{c} \cong_{\mathbb{C}} \mathbb{C}[\mathcal{D}] \otimes L^{o} \oplus \mathbb{C} \mathrm{c},
$$

where $L^{o}$ is a finite dimensional vector space and $\mathcal{D}$ is the translation operator of the vertex Lie algebra. Moreover, we will always consider the case when $\mathscr{L}$ is $\mathbb{Z}_{\geq 0}$-graded, in the sense that the underlying finite dimensional vector space $L^{o}$ is itself $\mathbb{Z}_{\geq 0}$-graded, $\mathcal{D}$ is an operator of degree $1, \mathrm{c}$ is of degree 0 and $\operatorname{deg}\left(a_{(n)} b\right)=\operatorname{deg} a+\operatorname{deg} b-n-1$ for any $a, b \in \mathscr{L}$ where ${ }_{(n)}: \mathscr{L} \times \mathscr{L} \rightarrow \mathscr{L}$ for $n \in \mathbb{Z}_{\geq 0}$ denote the $n^{\text {th }}$-products of $\mathscr{L}$. We denote by $L(0)$ the degree operator on $\mathscr{L}$, given by $L(0) a=(\operatorname{deg} a) a$ for homogeneous $a \in \mathscr{L}$. We will always assume that $\mathscr{L}$ is equipped with an action $\check{R}: \Gamma \rightarrow$ Aut $\mathscr{L}, \alpha \mapsto \check{R}_{\alpha}$ through automorphisms of $\mathscr{L}$, and consider the corresponding twisted action $R: \Gamma \rightarrow \mathrm{GL}(\mathscr{L}), \alpha \mapsto R_{\alpha}:=\alpha^{L(0)} \check{R}_{\alpha}$.

To any such vertex Lie algebra $\mathscr{L}$ we can associate a genuine Lie algebra, denoted $\mathrm{L}(\mathscr{L})$ and defined as $\operatorname{Lie}_{\mathbb{C}((t))} \mathscr{L}$ using [VY16b, Lemma 2.2] with $\mathcal{A}=\mathbb{C}((t))$, which is isomorphic as a vector space to $L^{o} \otimes \mathbb{C}((t)) \oplus \mathbb{C}$. The element in $\mathrm{L}(\mathscr{L})$ corresponding to $a \otimes t^{n}$ with $a \in L^{o}$ and $n \in \mathbb{Z}$ is denoted by $a(n) \in \mathrm{L}(\mathscr{L})$. We denote by $\mathrm{c}(-1)$ the copy of the element c belonging to $\mathrm{L}(\mathscr{L})$. The definition of $a(n), n \in \mathbb{Z}$, then extends to all $a \in \mathscr{L}$ by repeatedly applying $(\mathcal{D} a)(n)=-n a(n-1)$. In particular then $\mathrm{c}(n)=0$ unless $n=-1$. For a homogeneous vector $a \in L^{o}$ we also make use of the notation $a[n] \in \mathrm{L}(\mathscr{L})$ for the element corresponding to $a \otimes t^{n+\operatorname{deg} a-1}$. There is a vector space direct sum decomposition $\mathrm{L}(\mathscr{L})=\mathrm{L}^{-}(\mathscr{L})+\mathrm{L}^{+}(\mathscr{L})$ where $\mathrm{L}^{ \pm}(\mathscr{L})$ are Lie subalgebras isomorphic as vector spaces to $L^{o} \otimes \mathbb{C}[t t] \oplus \mathbb{C} c$ and $L^{o} \otimes t^{-1} \mathbb{C}\left[t^{-1}\right]$ respectively.

Denote by $\mathbb{V}(\mathscr{L})$ the vacuum Verma module over $\mathrm{L}(\mathscr{L})$, namely the $\mathrm{L}(\mathscr{L})$-module induced from the one-dimensional $\mathrm{L}^{+}(\mathscr{L})$-module $\mathbb{C} v_{0}$ on which c acts as 1 and $L^{o} \otimes \mathbb{C}[[t]]$ acts trivially. It is
naturally endowed with the structure of a vertex algebra. Finally, let $\mathrm{L}(\mathscr{L})^{\Gamma}$ denote the subalgebra of $\Gamma$-invariant elements, where the action of $\Gamma$ on $\mathrm{L}(\mathscr{L})$ is defined by letting $\alpha \in \Gamma$ send any $a(n)$ with $a \in \mathscr{L}$ and $n \in \mathbb{Z}$ to $\alpha^{-n-1}\left(R_{\alpha} a\right)(n)$. Given any $a \in \mathscr{L}$ and $n \in \mathbb{Z}$ we define the corresponding twisted $n^{\text {th }}$-mode

$$
a^{\Gamma}(n):=\sum_{\alpha \in \Gamma} \alpha^{-n-1}\left(R_{\alpha} a\right)(n) \in \mathrm{L}(\mathscr{L})^{\Gamma}
$$

The two main examples of vertex Lie algebras $\mathscr{L}$ we shall consider are given in [VY16b, Examples 2.5 \& 2.6], whose corresponding Lie algebras $\mathrm{L}(\mathscr{L})$ are centrally extended loop algebras and Heisenberg Lie algebras.
B.2. 'Local' and 'global' Lie algebras. Fix a set $\boldsymbol{x}:=\left\{x_{i}\right\}_{i=1}^{p}$ of $p \in \mathbb{Z}_{\geq 0}$ non-zero points in the complex plane. We attach to each point $x_{i} \in \boldsymbol{x}$ a local copy of the Lie algebra $\mathrm{L}(\mathscr{L})$ defined as $\operatorname{Lie}_{\mathbb{C}\left(\left(t-x_{i}\right)\right)} \mathscr{L}$ using [VY16b, Lemma 2.2] with $\mathcal{A}=\mathbb{C}\left(\left(t-x_{i}\right)\right)$, which we denote $\mathrm{L}(\mathscr{L})_{x_{i}}$. We will also use the index $x_{i}$ on the formal modes $a(n)$ or $a[n]$ when we wish to emphasise that these belong to $\mathrm{L}(\mathscr{L})_{x_{i}}$. Simiarly, to the origin we attach a local copy $\mathrm{L}(\mathscr{L})_{0}^{\Gamma}$ of the $\Gamma$-invariant subalgebra $\mathrm{L}(\mathscr{L})^{\Gamma}$. To the point at infinity we attach the Lie algebra $\mathrm{L}\left(\mathscr{L}^{\mathrm{op}}\right)_{\infty}^{\Gamma}$ defined as the subalgebra of $\Gamma$-invariants in $\mathrm{L}\left(\mathscr{L}^{\mathrm{op}}\right)_{\infty}:=\operatorname{Lie}_{\mathbb{C}\left(\left(t^{-1}\right)\right)} \mathscr{L}^{\mathrm{op}}$, where $\mathscr{L}^{\text {op }}$ is the opposite vertex Lie algebra given by the $\mathbb{C}[\mathcal{D}]$-module $\mathscr{L}$ but with opposite $n^{\text {th }}$-products. For any $n \in \mathbb{Z}$ and homogeneous element $a \in \mathscr{L}$, we denote by $a(n)_{\infty}\left(\right.$ resp. $\left.a[n]_{\infty}\right)$ the class of $a \otimes t^{n}\left(\right.$ resp. $\left.a \otimes t^{n+\operatorname{deg} a-1}\right)$ in $\mathrm{L}\left(\mathscr{L}^{\mathrm{op}}\right)_{\infty}$.

Consider the direct sum $\mathrm{L}\left(\mathscr{L}^{\text {op }}\right)_{\infty}^{\Gamma} \oplus \bigoplus_{i=1}^{p} \mathrm{~L}(\mathscr{L})_{x_{i}} \oplus \mathrm{~L}(\mathscr{L})_{0}^{\Gamma}$. We define the ideal

$$
\begin{equation*}
I_{\infty, p, 0}:=\operatorname{span}_{\mathbb{C}}\left\{\mathrm{c}(-1)_{x_{i}}-T \mathrm{c}(-1)_{\infty}\right\}_{i=1}^{p} \cup\left\{\mathrm{c}(-1)_{0}-\mathrm{c}(-1)_{\infty}\right\} \tag{B.1}
\end{equation*}
$$

and the corresponding quotient Lie algebra

$$
\mathrm{L}(\mathscr{L})_{\infty, \boldsymbol{x}, 0}:=\mathrm{L}\left(\mathscr{L}^{\mathrm{op}}\right)_{\infty}^{\Gamma} \oplus \bigoplus_{i=1}^{p} \mathrm{~L}(\mathscr{L})_{x_{i}} \oplus \mathrm{~L}(\mathscr{L})_{0}^{\Gamma} / I_{\infty, p, 0}
$$

Let $\mathbb{C}_{\infty, \Gamma \boldsymbol{x}, 0}(t)$ be the algebra of global rational functions with poles at most at $0, \infty$ and the points in $\Gamma \boldsymbol{x}$. It comes equipped with the derivation $\partial_{t}$ and an action of $\Gamma$ defined through pullback by the multiplication map $t \mapsto \alpha^{-1} t$ for $\alpha \in \Gamma$. Consider the associated Lie algebra $\operatorname{Lie}_{\mathbb{C}_{\infty, \Gamma x, 0}(t)} \mathscr{L}$ defined with the help of [VY16b, Lemma 2.2] for $\mathcal{A}=\mathbb{C}_{\infty, \Gamma \boldsymbol{x}, 0}(t)$. We will denote by $a f:=\rho(a \otimes f)$ the class in $\operatorname{Lie}_{\mathbb{C}_{\infty, \Gamma \boldsymbol{x}, 0}(t)} \mathscr{L}$ of an element $a \otimes f \in \mathscr{L} \otimes \mathbb{C}_{\infty, \Gamma \boldsymbol{x}, 0}(t)$. The action of $\Gamma$ on $\mathscr{L} \otimes \mathbb{C}_{\infty, \Gamma \boldsymbol{x}, 0}(t)$ defined for any $\alpha \in \Gamma$ by

$$
\alpha .(a \otimes f(t)):=\alpha^{-1} R_{\alpha} a \otimes f\left(\alpha^{-1} t\right)
$$


Consider the ideal in $\operatorname{Lie}_{\mathbb{C}_{\infty, \Gamma x, 0}(t)} \mathscr{L}$ defined by

$$
\begin{equation*}
I_{\infty, \Gamma \boldsymbol{x}, 0}:=\operatorname{span}_{\mathbb{C}}\left\{\sum_{\alpha \in \Gamma} \frac{\mathrm{c}}{t-\alpha x}\right\}_{x \in \boldsymbol{x} \cup\{0\}}=\operatorname{span}_{\mathbb{C}}\left\{\sum_{\alpha \in \Gamma} \frac{\mathrm{c}}{t-\alpha x_{i}}\right\}_{i=1}^{p} \cup\left\{\frac{T \mathrm{c}}{t}\right\} \tag{B.2}
\end{equation*}
$$

and denote the corresponding quotient Lie algebra by

$$
\mathrm{L}_{\infty, \Gamma \boldsymbol{x}, 0}(\mathscr{L}):=\operatorname{Lie}_{\mathbb{C}_{\infty, \Gamma \boldsymbol{x}, 0}(t)} \mathscr{L} / I_{\infty, \Gamma \boldsymbol{x}, 0}
$$

Noting that the ideal (B.2) is invariant under the action of $\Gamma$, i.e. $\Gamma . I_{\infty, \Gamma \boldsymbol{x}, 0}=I_{\infty, \Gamma \boldsymbol{x}, 0}$, we obtain a well-defined action of $\Gamma$ on the quotient $L_{\infty, \Gamma \boldsymbol{x}, 0}(\mathscr{L})$. We denote the corresponding subalgebra of
$\Gamma$-invariants by

$$
\begin{equation*}
\mathrm{L}_{\infty, \boldsymbol{x}, 0}^{\Gamma}(\mathscr{L}):=\left(\mathrm{L}_{\infty, \Gamma x, 0}(\mathscr{L})\right)^{\Gamma} \tag{B.3}
\end{equation*}
$$

Proposition B.1. There is an embedding of Lie algebras

$$
\begin{equation*}
\iota: \mathrm{L}_{\infty, \boldsymbol{x}, 0}^{\Gamma}(\mathscr{L}) \hookrightarrow \mathrm{L}(\mathscr{L})_{\infty, \boldsymbol{x}, 0} . \tag{B.4}
\end{equation*}
$$

Proof. Taking the Laurent expansion of an element of $\operatorname{Lie}_{\mathbb{C}_{\infty, \Gamma \boldsymbol{x}, 0}(t)} \mathscr{L}$ at $0, \infty$ and the points in $\boldsymbol{x}$ yields an embedding of Lie algebras

$$
\begin{align*}
\iota: \mathrm{Lie}_{\mathbb{C}_{\infty, \Gamma x, 0}(t)} \mathscr{L} & \longleftrightarrow \mathrm{L}\left(\mathscr{L}^{\mathrm{op}}\right)_{\infty} \oplus \bigoplus_{i=1}^{p} \mathrm{~L}(\mathscr{L})_{x_{i}} \oplus \mathrm{~L}(\mathscr{L})_{0}  \tag{B.5}\\
a f & \longmapsto\left(-\rho\left(a \otimes \iota_{t^{-1}} f\right), \rho\left(a \otimes \iota_{t-x_{1}} f\right), \ldots, \rho\left(a \otimes \iota_{t-x_{p}} f\right), \rho\left(a \otimes \iota_{t} f\right)\right) .
\end{align*}
$$

Note that under the embedding (B.5) we have $\iota\left(\sum_{\alpha \in \Gamma} \frac{\mathrm{c}}{t-\alpha x}\right)=\mathrm{c}(-1)_{x}-T \mathrm{c}(-1)_{\infty}$ for every $x \in \boldsymbol{x}$ and $\iota\left(\frac{T \mathrm{c}}{t}\right)=T \mathrm{c}(-1)_{0}-T \mathrm{c}(-1)_{\infty}$ from which we deduce that the ideal $I_{\infty, \Gamma \boldsymbol{x}, 0} \subset \operatorname{Lie}_{\mathbb{C}_{\infty, \Gamma \boldsymbol{x}, 0}(t)} \mathscr{L}$ defined in (B.2) is mapped to the ideal $I_{\infty, p, 0} \subset \mathrm{~L}\left(\mathscr{L}^{\mathrm{op}}\right)_{\infty} \oplus \bigoplus_{i=1}^{p} \mathrm{~L}(\mathscr{L})_{x_{i}} \oplus \mathrm{~L}(\mathscr{L})_{0}$ defined in (B.1). By [VY16b, Lemma 2.9] it follows that (B.5) induces an embedding of quotient Lie algebras

$$
\iota: \mathrm{L}_{\infty, \Gamma x, 0}(\mathscr{L}) \hookrightarrow \mathrm{L}\left(\mathscr{L}^{\mathrm{op}}\right)_{\infty} \oplus \bigoplus_{i=1}^{p} \mathrm{~L}(\mathscr{L})_{x_{i}} \oplus \mathrm{~L}(\mathscr{L})_{0} / I_{\infty, p, 0} .
$$

Finally, by restricting the latter to the subalgebra of $\Gamma$-invariants (B.3) we obtain the desired Lie algebra embedding (B.4).

Let $\mathcal{M}_{x_{i}}$ be a left module over $\mathrm{L}(\mathscr{L})_{x_{i}}$ for each $i=1, \ldots, p$. Following the convention adopted in [VY16b, Section 3], we will always assume that such modules over $\mathrm{L}(\mathscr{L})$ are of level 1, namely that the central element $\mathrm{c}(-1)$ acts as 1 . Let $\mathcal{M}_{0}$ and $\mathcal{M}_{\infty}$ be left modules over $\mathrm{L}(\mathscr{L})_{0}^{\Gamma}$ and $\mathrm{L}\left(\mathscr{L}^{\circ \mathrm{p}}\right)_{\infty}^{\Gamma}$ respectively, or alternatively $\mathcal{M}_{\infty}$ is a right module over $\mathrm{L}(\mathscr{L})_{\infty}^{\Gamma}$. Following also the convention of [VY16b, Section 3], all such modules over $\mathrm{L}(\mathscr{L})^{\Gamma}$ will be assumed to be of level $\frac{1}{T}$. The tensor product $\mathcal{M}_{\infty} \otimes \bigotimes_{i=1}^{p} \mathcal{M}_{x_{i}} \otimes \mathcal{M}_{0}$ is then a left module over the Lie algebra $\mathrm{L}(\mathscr{L})_{\infty, \boldsymbol{x}, 0}$ so that by virtue of Proposition B. 1 it also becomes a left module over the global Lie algebra $L_{\infty, x, 0}^{\Gamma}(\mathscr{L})$. In particular, we may form the space of coinvariants

$$
\mathcal{M}_{\infty} \otimes \bigotimes_{i=1}^{p} \mathcal{M}_{x_{i}} \otimes \mathcal{M}_{0} / \mathrm{L}_{\infty, \boldsymbol{x}, 0}^{\Gamma}(\mathscr{L})
$$

Many of the statements and proofs of [VY16b] concerning cyclotomic coinvariants which did not include the point at infinity can be seen to carry over with minor modifications to the present case.
B.3. $Y_{W}$-map. A left module $M$ over $\mathrm{L}(\mathscr{L})$ is said to be smooth if for all $a \in \mathscr{L}$ and $v \in M$ we have $a(n) v=0$ for all $n \gg 0$. Likewise, we will say that a right module $N$ over $\mathrm{L}(\mathscr{L})$ is co-smooth if for all $a \in \mathscr{L}$ and $\eta \in N$ we have $\eta a(n)=0$ for all $n \ll 0$.

Let $\mathcal{M}_{\infty}$ be any right co-smooth module over $\mathrm{L}(\mathscr{L})^{\Gamma}$. We define the quasi-module map

$$
Y_{W}(\cdot, u): \mathbb{V}(\mathscr{L}) \longrightarrow \operatorname{Hom}\left(\mathcal{M}_{\infty}, \mathcal{M}_{\infty}\left(\left(u^{-1}\right)\right)\right), \quad A \longmapsto \sum_{n \in \mathbb{Z}} A_{(n)}^{W} u^{-n-1}
$$

where $u$ is a formal variable and $A_{(n)}^{W}$ are endomorphisms of $\mathcal{M}_{\infty}$ for each $A \in \mathbb{V}(\mathscr{L})$ and $n \in \mathbb{Z}$, by direct analogy with the definition given in [VY16b] for left smooth modules over $\mathrm{L}(\mathscr{L})^{\Gamma}$ as follows. For any $a \in \mathscr{L}$ we set

$$
Y_{W}\left(a(-1) v_{0}, u\right):=\sum_{\alpha \in \Gamma} \sum_{n \in \mathbb{Z}}\left(R_{\alpha} a\right)(n)(\alpha u)^{-n-1}
$$

Moreover, the map is defined recursively for any other state in $\mathbb{V}(\mathscr{L})$ by letting

$$
\begin{equation*}
Y_{W}(a(-1) B, u):=: Y_{W}\left(a(-1) v_{0}, u\right) Y_{W}(B, u):+\sum_{\alpha \in \Gamma \backslash\{1\}} \sum_{n \geq 0} \frac{1}{((\alpha-1) u)^{n+1}} Y_{W}\left(\left(R_{\alpha} a\right)(n) B, u\right) \tag{B.6}
\end{equation*}
$$

for all $a \in \mathscr{L}$ and $B \in \mathbb{V}(\mathscr{L})$. Here $: A(u) B(u):=A(u)_{+} B(u)+B(u) A(u)_{-}$denotes the usual normal ordering where for $A(u)=\sum_{n \in \mathbb{Z}} A(n) u^{-n-1}$ we define

$$
\begin{align*}
& A(u)_{+}=\sum_{n<0} A(n) u^{-n-1}=\sum_{m \geq 0} A(-m-1) u^{m}  \tag{B.7a}\\
& A(u)_{-}=\sum_{n \geq 0} A(n) u^{-n-1} \tag{B.7b}
\end{align*}
$$

The following is a direct analogue of [VY16b, Proposition 3.6] for the point at infinity, providing an alternative definition the above $Y_{W}$-map on co-smooth modules using co-invariants.

Proposition B.2. Let $\mathcal{M}_{\infty}$ be a right co-smooth module over $\mathrm{L}(\mathscr{L})^{\Gamma}$. We have

$$
\begin{equation*}
\left.\left.-\iota_{u^{-1}}^{[A} \underset{\substack{\uparrow \\ u} \underset{\infty}{m_{\infty}}}{m_{\infty}} \otimes \cdots\right] \underset{\substack{\uparrow}}{m_{\infty}} Y_{W}(A, u) \otimes \cdots\right], \tag{B.8}
\end{equation*}
$$

for all $A \in \mathbb{V}(\mathscr{L})$ and $m_{\infty} \in \mathcal{M}_{\infty}$.

Proof. We use induction on the depth of $A$. When $A=v_{0}$ the result follows from the analogue of [VY16b, Proposition 3.2] including the point at infinity. For the inductive step, we assume that (B.8) holds for states of depth strictly less than that of $A$. Without loss of generality we can take the state $A$ to be of the form $A=a(-1) B$ for some $a \in \mathscr{L}$ and $B \in \mathbb{V}(\mathscr{L})$.

Let $\mathcal{M}_{x_{i}}, i=1, \ldots, p$ be any collection of $p \in \mathbb{Z}_{\geq 0}$ modules attached to the points $x_{i} \in \mathbb{C}$ for $i=1, \ldots, p$ which may include the origin. We write $\boldsymbol{x}:=\left\{x_{i}\right\}_{i=1}^{p}$ and let $\boldsymbol{m} \in \bigotimes_{i=1}^{p} \mathcal{M}_{x_{i}}$. By definition of the space of coinvariants we have

In other words,

Thus the left hand side of (B.8) may be written as

$$
\begin{aligned}
& -\iota_{u^{-1}}\left[\underset{\substack{\hat{\lambda} \\
u}}{B} \otimes \underset{\substack{\hat{\infty}}}{m_{\infty}} Y_{W}\left(a(-1) v_{0}, u\right)_{+} \otimes \underset{\hat{x}}{\boldsymbol{m}}\right]-\iota_{u^{-1}}\left[\underset{\substack{\hat{u} \\
B}}{B} \underset{\substack{\hat{\infty}}}{m_{\infty}} \otimes \sum_{i=1}^{p} \sum_{\alpha \in \Gamma} \sum_{n \geq 0} \frac{\left(R_{\alpha} a\right)(n)_{x_{i}}}{\left(\alpha u-x_{i}\right)^{n+1}} \underset{\hat{x}}{\boldsymbol{x}}\right] .
\end{aligned}
$$

Here we made use of the notation (B.7a). The states at the point $u$ on the right hand side are of lower depth in $\mathbb{V}(\mathscr{L})$, so that we may apply the inductive hypothesis to obtain

The first two terms on the right hand side are already in the desired form. Taking the map $\iota_{u^{-1}}$ explicitly in the remaining term we may rewrite it as

Now consider the following identity

$$
\sum_{m \geq 0} u^{-m-1}\left[g_{m}(t) \cdot \underset{\substack{\uparrow \\ \infty}}{m_{\infty} Y_{W}}(B, u) \otimes \underset{\substack{\uparrow \\ \boldsymbol{x}}}{\boldsymbol{m})}\right]=0, \quad \text { where } \quad g_{m}(t)=\sum_{\alpha \in \Gamma} \alpha^{-1}\left(R_{\alpha} a\right)\left(\alpha^{-1} t\right)^{m} .
$$

Using this we may rewrite (B.9) simply as

$$
\left.\sum_{m \geq 0} u^{-m-1} \underset{\substack{\uparrow \\ m_{\infty}}}{m_{W}} Y_{W}(B, u)\left(\iota_{t^{-1}} g_{m}(t)\right) \otimes \underset{\substack{\uparrow \\ \boldsymbol{x}}}{\boldsymbol{m}}\right] \underset{\substack{\uparrow \\ \infty}}{\left[m_{\infty} Y_{W}(B, u) Y_{W}\left(a(-1) v_{0}, u\right)_{-} \otimes \underset{\boldsymbol{x}}{\boldsymbol{m}}\right] . . . . ~}
$$

Putting the above together and using the recurrence relation (B.6) of the $Y_{W}$-map we obtain
as required.
Recall that the vacuum Verma module $\mathbb{V}(\mathscr{L})$ has the structure of a vertex algebra and so is in particular a vertex Lie algebra. We may thus form the associated Lie algebra $\mathrm{L}(\mathbb{V}(\mathscr{L}))$ as well as its subalgebra of $\Gamma$-invariants $\mathrm{L}(\mathbb{V}(\mathscr{L}))^{\Gamma}$. The following is an immediate generalisation of [VY16b, Proposition 5.8].

Proposition B.3. Let $M_{\infty}$ be a co-smooth right module over $\mathrm{L}(\mathscr{L})^{\Gamma}$. There is a well-defined co-smooth right $\mathrm{L}(\mathbb{V}(\mathscr{L}))^{\Gamma}$-module structure on $M_{\infty}$ given for all $m_{\infty} \in M_{\infty}$ by

$$
m_{\infty} A^{\Gamma}(n):=m_{\infty} A_{(n)}^{W} .
$$

## References

[AKS] M. Adler, On a trace functional for formal pseudo-differential operators and the symplectic structure for the KdV type equations, Invent. Math. 50 (1979), 219-248.
B. Kostant, The solution to a generalized Toda lattice and representation theory, Adv. Math. 34 (1980), 13-53.
W. Symes, Systems of Toda type, inverse spectral problems and representation theory, Invent. Math. 59 (1980), 195-338.
[ATY91] H. Awata, A. Tsuchiya, Y. Yamada, Integral Formulas for the WZNW Correlation Functions, Nucl. Phys. B365 (1991), 680-696.
[BF94] H. Babujian, R. Flume, Off-shell Bethe ansatz equation for Gaudin magnets and solutions of KnizhnikZamolodchikov equations, Mod. Phys. Lett. A9 (1994), 2029-2040.
[BD96] A. A. Beilinson, V. G. Drinfeld, Quantization of Hitchin's Fibration and Langland's Program, Math. Phys. Studies 19 (1996) 3-7.
[BD98] A. A. Belavin, V. G. Drinfeld, Triangle equations and simple Lie algebras, Classic Reviews in Mathematics and Mathematical Physics. 1. Amsterdam: Harwood Academic Publishers. vii, 91 p. (1998).
[Bla69] R. J. Blattner, Induced and Produced Representations of Lie Algebras, Trans. Amer. Math. Soc. 144 (1969) 457-474.
[Bro10] A. Brochier, A Kohno-Drinfeld theorem for the monodromy of cyclotomic KZ connections, Comm. Math. Phys. 311 (2012) no. 1, 55-96.
[CS] O. Chalykh, A. Silantyev, KP hierarchy for the cyclic quiver, arXiv:1512.08551 [math.QA].
[CM08] A. V. Chervov, A. I. Molev, On Higher Order Sugawara operators, Int. Math. Res. Not. (2009), no. 9, 1612-1635.
[CT06] A. Chervov, D. Talalaev, Quantum spectral curves, quantum integrable systems and the geometric Langlands correspondence, Preprint, 2006. [arXiv hep-th/0604128]
[CY07] N. Crampé, C. A. S. Young, Integrable models from twisted half-loop algebras, J. Phys. A: Math. Theor. 40 (2007) 5491.
[Di74] J. Dixmier, Algèbres Enveloppantes, Gauthier-Villars, Paris, 1974.
[Doy06] B. Doyon, Twisted Modules for Vertex Operator Algebras, in proc. "Moonshine - the First Quarter Century and Beyond, a Workshop on the Moonshine Conjectures and Vertex Algebras" (Edinburgh, 2004), Lond. Math. Soc. Lecture Notes, CUP, ed. by J. Lepowsky, J. McKay and M. Tuite (2008)
[Dra05] J. Draisma, Representation theory on the open bruhat cell, J. Symb. Comp. 39 (2005) 279-303.
[ER96] B. Enriquez, V. Rubtsov, Hitchin systems, higher Gaudin operators and r-matrices, Math. Res. Lett. 3 (1996) 343-357.
[Fed10] R. M. Fedorov. Irregular Wakimoto modules and the Casimir connection, Selecta Math. (N.S.), 16 (2):241266, 2010.
[FF90] B. Feigin, E. Frenkel, Affine Kac-Moody algebras and semi-infinite flag manifolds, Comm. Math. Phys. 128 (1990) 161-189.
[FF92] B. Feigin, E. Frenkel, Affine Kac-Moody algebras at the critical level and Gelfand-Dikii algebras, Int. J. Mod. Phys. A7, Suppl. 1A (1992), 197-215.
[FF07] B. Feigin, E. Frenkel, Quantization of soliton systems and Langlands duality, [arXiv:0705.2486]
[FFR94] B. Feigin, E. Frenkel, N. Reshetikhin, Gaudin model, Bethe ansatz and critical level, Comm. Math. Phys. 166 (1994), no. 1, 27-62.
[FFRb10] B. Feigin, E. Frenkel, L. Rybnikov, Opers with irregular singularity and spectra of the shift of argument subalgebra, Duke Math. J. 155 (2010), no. 2, 337-363.
[FFT10] B. Feigin, E. Frenkel, and V. Toledano Laredo, Gaudin models with irregular singularities, Adv. Math. 223 (2010), no. 3, 873-948
[FMTV00] G. Felder, Y. Markov, V. Tarasov, A. Varchenko, Differential Equations Compatible with KZ equations Math. Phys. Anal. Geom. 3 (2000), no. 2, 139-177.
[Fre05] E. Frenkel, Gaudin model and opers, Prog. in Math. 237 (2005) 1-58.
[Fre07] E. Frenkel, Langlands correspondence for loop groups, Cambridge University Press, 2007.
[FB04] E. Frenkel and D. Ben-Zvi, Vertex algebras and algebraic curves, (Mathematical surveys and monographs. vol 88) American Mathematical Society, 2004 (Second Edition).
[Gau76] M. Gaudin, Diagonalisation d'une classe d'Hamiltoniens de spin, J. Physique 37 (1976), $1087-1098$.
[Gau83] M. Gaudin, La fonction d'onde de Bethe, Collection du Commissariat l'Érgie Atomique, Série Scientifique, Masson, Paris, 1983.
[Hu97] J. Humphreys, Introduction to Lie Algebras and Representation Theory, Springer, 1972.
[Hu08] J. Humphreys, Representations of Semisimple Lie Algebras in the BGG Category $\mathcal{O}$, Grad. Studies in Math. 94, Amer. Math. Soc. , 2008.
[Ka85] M. Kashiwara, The Universal Verma Module and the b-Function, Adv. Studies in Pure Math. 6 (1985), 67-81.
[Kac83] V. Kac, Infinite Dimensional Lie Algebras: An Introduction, Third Edition, Cambridge Univ. Press, 1990.
[Kac98] V. Kac, Vertex Algebras for Beginners, American Mathematical Society, 1998 (Second Edition).
[LV06] S. Lacroix and B. Vicedo, Cyclotomic Gaudin models, Miura opers and flag varieties, arXiv:1607.07397 [math.QA].
[Li06a] H. Li, A new construction of vertex algebras and quasi modules for vertex algebras, 2004 Adv. Math. 202 (2006) Issue 1, 232-286.
[Li06b] H. Li, On certain generalisations of twisted affine Lie algebras and quasi modules for $\Gamma$-vertex algebras, 2006 [arXiv:math/0605155].
[MF78] A. S. Mishchenko, A. T. Fomenko, Euler equation on finite-dimensional Lie groups, Math. USSR-Izv. 12 (1978), 371-389.
[Mo13] A. I. Molev, Feigin-Frenkel center in types B, C and D, Invent. Math., 191 (2013) Issue 1, 1-34.
[MV00] E. Mukhin, A. Varchenko, Remarks on critical points of phase functions and norms of Bethe vectors in Adv. Studies in Pure Math. 27 (2000), Arrangements - Tokyo, 239-246
[MV05] E. Mukhin, A. Varchenko,t Norm of a Bethe vector and the Hessian of the master function, Compositio Math. 141 (2005), no. 4, 1012-1028.
[MV07] E. Mukhin, A. Varchenko, Multiple orthogonal polynomials and a counterexample to Gaudin Bethe Ansatz Conjecture, Trans. Amer. Math. Soc. 359 (2007), no. 11, 5383-5418.
[MTV06] E. Mukhin, V. Tarasov, A. Varchenko, Bethe Eigenvectors of Higher Transfer Matrices, J. Stat. Mech. (2006), no. 8, P08002, 1-44.
[MTV09] E. Mukhin, V. Tarasov, A. Varchenko, Schubert calculus and representations of the general linear group, J. Amer. Math. Soc. 22 (2009), no. 4, 909-940.
[MTV14] E. Mukhin, V. Tarasov, A. Varchenko, Bethe algebra of Gaudin model, Calogero-Moser space, and Cherednik algebra, Int. Math. Res. Not. 5 (2014) 1174-1204.
[RV95] N. Reshetikhin, A. Varchenko, Quasiclassical asymptotics of solutions to the KZ equations, Geometry, topology, and physics, Conf. Proc. Lecture Notes Geom. Topology, IV, Int. Press, Cambridge, MA (1995), 293-322.
[Ryb06] L. .G .Rybnikov, The argument shift method and the Gaudin model, Func. Anal. App. 40 (2006) no.3, 188-199.
[Ryb16] L. G. Rybnikov, Cactus Group and Monodromy of Bethe Vectors, Int. Math. Res. Not. (2016), 00, pp. 1-35.
[Ryb08] L. .G.Rybnikov, Uniqueness of higher Gaudin Hamiltonians, Rep. Math. Phys. 61 (2008) Issue 2, $247-252$.
[Skr06] T. Skrypnyk, Integrable quantum spin chains, non-skew symmetric r-matrices and quasigraded Lie algebras J. Geom. and Phys. 57 (2006), no. 1, 53-67.
[Skr13] T. Skrypnyk, $\mathbb{Z}_{2}$-graded Gaudin models and analytical Bethe ansatz, Nucl. Phys. B 870 (2013), no. 3, 495-529.
[SV91] V. Schechtman, A. Varchenko, Arrangements of hyperplanes and Lie algebra homology, Invent. Math. 106 (1991), no. 1, 139-194.
[Sz02] M. Szczesny, Wakimoto modules for twisted affine Lie algebras, Math. Res. Lett. 9 (2002), 433-448
[Ta04] D. Talalaev, Quantization of the Gaudin System, Funct. Anal. Its Appl. 40 (2006), Issue 1, 73-77.
[Var95] A. Varchenko, Asymptotic solutions the the Knizhnik-Zamolidchikov equation and crystal base Comm. Math. Phys. 171 (1995) no. 1, 99-137.
[VY] A. Varchenko, C. A. S. Young, Cyclotomic discriminantal arrangements and diagram automorphisms of Lie algebras, [arXiv:1603.07125]
[VY16a] B. Vicedo, C. A. S. Young, Cyclotomic Gaudin models: construction and Bethe ansatz, Comm. Math. Phys. 343 (2016) no. 3,971-1024.
[VY16b] B. Vicedo, C. A. S. Young, Vertex Lie algebras and cyclotomic coinvariants Commun. Contemp. Math. DOI: http://dx.doi.org/10.1142/S0219199716500152
[Wa86] M. Wakimoto, Fock representations of affine Lie algebra $A_{(1)}^{1}$, Comm. Math. Phys. 104 (1986), 605-609.

School of Physics, Astronomy and Mathematics, University of Hertfordshire, College Lane, Hatfield AL10 9AB, UK.

E-mail address: benoit.vicedo@gmail.com
E-mail address: c.a.s.young@gmail.com


[^0]:    ${ }^{1}$ Actually, we define these $\mathscr{Z}_{\infty, \boldsymbol{z}, 0}^{n_{\infty}, \boldsymbol{n}, n_{0}}(\mathfrak{g}, \sigma)^{\Gamma}$ first, and then the universal algebra $\mathscr{Z}_{\infty, \boldsymbol{z}, 0}(\mathfrak{g}, \sigma)^{\Gamma}$ is their inverse limit.

[^1]:    ${ }^{2}$ Indeed, suppose $\alpha \in \Delta^{+}$is a root of $\mathfrak{g}$ such that the orbit $\sigma^{\mathbb{Z}} \alpha$ has $t$ elements, where $t \in \mathbb{Z}_{\geq 1}$ divides $T$. Then $\sigma^{t} E_{\alpha}=\omega^{t m} E_{\alpha}$ for some unique $m \in \mathbb{Z} / T \mathbb{Z}$. Let $E_{(m-k T / t, \bar{\alpha})}:=\sum_{j=0}^{t-1} \omega^{-(m-k T / t) j} \sigma^{j} E_{\alpha} \in \overline{\mathfrak{n}}_{(m-k T / t, \bar{\alpha})}$ for $k=$ $0,1, \ldots, t-1$. By picking one root from each $\sigma$-orbit, we construct a basis of the required form.

[^2]:    ${ }^{3}$ That is, for each $v \in W_{\nu_{i}} \alpha \in \Delta^{+}$and $k \in I$ there exists an $n \in \mathbb{Z}$ such that $0=a_{\alpha}[m]_{x_{i}} v=a_{\alpha}^{*}[m]_{x_{i}} v=b_{k}[m]_{x_{i}} v$ for all $m \geq n$.

