

## THE *CHANDRA*, *HUBBLE SPACE TELESCOPE*, AND VLA VIEW OF THE CIRCUMNUCLEAR EXTENDED EMISSION IN THE NARROW EMISSION LINE GALAXY NGC 2110

DANIEL A. EVANS,<sup>1</sup> JULIA C. LEE,<sup>1</sup> MARIA KAMENETSKA,<sup>2</sup> SARAH C. GALLAGHER,<sup>3</sup> RALPH P. KRAFT,<sup>1</sup>  
MARTIN J. HARDCASTLE,<sup>4</sup> AND KIMBERLY A. WEAVER<sup>5</sup>

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### ABSTRACT

We present results from new *Chandra*, archival *Hubble Space Telescope* (*HST*), and VLA imaging observations of the circumnuclear extended emission in the nearby Type 2 Seyfert galaxy NGC 2110. We find resolved soft-band X-ray emission  $\sim 4''$  ( $\sim 160$  pc) north of the nucleus, which is spatially coincident with [O III] emission, but lies just beyond the northern edge of the radio jet in the source. We find that shock heating of multiphase gas clouds can successfully account for this extended emission, although we cannot rule out alternative models, such as the scattering of nuclear radiation by ionized material, or pure photoionization from the nucleus. In addition, we detect kiloparsec-scale ( $\sim 30''$ ) extended soft-band X-ray emission south of the nucleus. Finally, we compare our results for NGC 2110 with the prototypical type 2 Seyfert galaxy NGC 1068, and suggest that different physical processes could produce extended circumnuclear X-ray emission in Seyfert galaxies.

*Subject headings:* galaxies: active — galaxies: individual (NGC 2110) — galaxies: jets — galaxies: Seyfert

### 1. INTRODUCTION

X-ray observations of the circumnuclear environments of Seyfert galaxies with small-scale radio jets are important diagnostics for understanding the interplay between the ionizing radiation from the nucleus, radio ejecta, and galaxy ISM. In some of the brightest sources, e.g., NGC 1068 (e.g., Young et al. 2001), Mrk 3 (Sako et al. 2000), and the Circinus galaxy (Sambruna et al. 2001), high-resolution *Chandra* and *XMM-Newton* grating spectroscopy reveals their X-ray spectra to be dominated by He-like emission lines most likely to be produced in a plasma photoionized by nuclear radiation. However, alternative models, such as shock heating induced by the interaction of the radio jet and ambient medium may be applicable to other Seyfert galaxies. High-resolution X-ray imaging observations, combined with multiwavelength data, are important probes of circumnuclear environments, and are complementary to X-ray grating data.

NGC 2110, the subject of this paper, is a nearby ( $z = 0.007579$ ,  $D_L = 32.7$  Mpc) Seyfert type 2 galaxy (Shuder 1980). Historically, it is classified as a narrow emission line galaxy (NELG), which is the subclass of Seyfert galaxies with strong 2–10 keV X-ray emission and narrow ( $< 600$  km s<sup>-1</sup>) optical emission lines. The source exhibits complex morphology on a range of angular scales, and shows a rich variety of structure across the electromagnetic spectrum. VLA observations (Ulvestad & Wilson 1989) show that it hosts an S-shaped jet-counterjet system that extends  $\sim 4''$  ( $\sim 600$  pc) in the north-south direction, while H I spectral-line mapping (Gallimore et al. 1999) indicates the presence of a small-scale disk that may partially obscure the southern radio jet. Optically, NGC 2110 is classified as a S0 galaxy, although broad-

band *Hubble Space Telescope* (*HST*) observations (Mulchaey et al. 1994; González Delgado et al. 2002; Ferruit et al. 2004) have revealed a series of complex, spiral-like dust lanes with a distribution consistent with a tightly wrapped single-armed spiral.

Emission line [O III]  $\lambda 5007$  and  $H\alpha + [N II]$  *HST* observations of the inner regions of NGC 2110 (Mulchaey et al. 1994) have revealed a variety of structures. The dominant line emission is concentrated into a strongly curved “jetlike” feature (as characterized by Mulchaey et al. 1994) that extends  $\sim 1''.4$  ( $\sim 230$  pc) north of the nucleus and becomes increasingly broadened with distance. The [O III] emission is noticeably more asymmetric about the nucleus than the  $H\alpha$  emission, and Mulchaey et al. (1994) suggest that this is at least in part due to a north-south gradient in excitation. Detailed spectroscopy of the “jetlike” feature with the *HST* Faint Object Spectrograph (FOS; Ferruit et al. 1999) suggests that photoionizing-shock models (e.g., Dopita & Sutherland 1996) provide a good explanation of the origin of this emission, although photoionization by a central source cannot be ruled out. Recently, Optically Adaptive System for Infrared Spectroscopy (OASIS) integral field spectroscopy of the same region (Ferruit et al. 2004) confirmed that the majority of the narrow [O III] emission is not associated with the nucleus, and may rather be associated with the jetlike feature. This is in contrast to the  $H\alpha$  emission, whose peak intensity is coincident with that of the radio nucleus. On slightly larger scales, Mulchaey et al. (1994) also report the detection of an S-shaped distribution of ionized gas that extends  $\sim 4''$  in either direction of the nucleus. Although this feature, together with the jetlike feature, extends along a similar position angle to the radio emission, neither appears to be spatially coincident with the radio morphology. Instead, the strongest optical line emission occurs where the radio emission is weak, and the optical jetlike feature is more curved than the radio jet. This suggests that, although possibly influenced by the radio-emitting feature, the optical line emission is not directly associated with it.

At X-ray energies, *ROSAT* and Broad Band X-Ray Telescope (BBXRT) observations of NGC 2110 (Weaver et al. 1995) showed the presence of spatially resolved X-ray emission that extends  $\sim 4''$  north of the unresolved nucleus, possibly coincident with

<sup>1</sup> Harvard-Smithsonian Center for Astrophysics, 60 Garden Street, Cambridge, MA 02138.

<sup>2</sup> MIT Kavli Institute for Astrophysics and Space Research, 77 Massachusetts Avenue, NE 80, Cambridge, MA 02139.

<sup>3</sup> UCLA, Division of Astronomy and Astrophysics, 430 Portola Plaza, Mail Code 154705, Los Angeles, CA 90095-1547.

<sup>4</sup> School of Physics, Astronomy and Mathematics, University of Hertfordshire, College Lane, Hatfield AL10 9AB, UK.

<sup>5</sup> NASA Goddard Space Flight Center, Code 662, Greenbelt, MD 20771.

the optical emission-line S-shaped feature. They argued that at least a fraction of the soft excess observed in the X-ray spectrum of NGC 2110 below 1 keV (e.g., Mushotzky 1982) originates in this extended region. However, the physical origin of the extended emission was unclear, and several models have been proposed to explain it. In one model, the impact of the radio jet into the ISM may induce shock heating, producing the observed soft X-ray emission, together with the optical emission lines. Ferruit et al. (1999) provide an in-depth discussion of photoionizing-shock models, and demonstrated that if in situ shock heating were to power the X-ray emission, shock velocities in excess of  $500 \text{ km s}^{-1}$  would be required (cf. the measured [O III] line width of  $200 \text{ km s}^{-1}$  in this region (e.g., Wilson et al. 1985). However, the shock heating could occur in a region much closer to the nucleus, where the measured [O III] velocities are higher, and subsequently expand outwards. Another interpretation, discussed by Weaver et al. (1995) is that the extended X-ray emission is electron-scattered nuclear radiation. Alternatively, the extended emission may result from pure photoionization of ambient gas by the nucleus. A final interpretation, given the common detection with *Chandra* of resolved kiloparsec-scale X-ray jets in low-power radio galaxies (e.g., Worrall et al. 2001), is that the resolved X-ray emission is from the VLA-detected radio ejecta. Dong et al. (2004) showed a series of zeroth-order *Chandra* HETGS images of the nucleus of NGC 2110, which hinted at the presence of extended soft-band X-ray emission north of the nucleus, and we use the superb spatial resolution and good low-energy sensitivity of the *Chandra* ACIS instrument to provide further constraints on the nature of this emission.

In this paper, we present a multiwavelength imaging study of the circumnuclear extended emission in NGC 2110. We make a detailed comparison between the *Chandra* X-ray and archival radio and optical observations of the extended emission north of the nucleus, and discuss models that might explain its origin. In addition, we present the first evidence for large-scale soft-band X-ray emission south of the nucleus. Finally, we compare our results for NGC 2110 with the prototypical type 2 Seyfert galaxy NGC 1068, and argue that different physical processes could produce the extended circumnuclear X-ray emission in Seyfert galaxies. A detailed analysis of the *Chandra* HETGS spectra of NGC 2110 will be presented in a future paper (D. A. Evans et al. 2007, in preparation), and the energetics arguments presented in this paper are in part a prelude to our analysis of the high-energy resolution *Chandra* gratings data.

This paper is organized as follows. Section 2 contains a description of the data and a summary of their analysis. In § 3 we present the *Chandra* imaging observation and compare it to observations with *HST* and the VLA. In § 4 we discuss several models that may account for the small-scale northern extended emission, while in § 5 we discuss the large-scale southern extended X-ray emission. In § 6 we compare NGC 2110 with the prototypical Seyfert 2 galaxy NGC 1068. We end with our conclusions in § 7. All results presented in this paper use a cosmology in which  $\Omega_{m,0} = 0.3$ ,  $\Omega_{\Lambda,0} = 0.7$ , and  $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$ . At the redshift of NGC 2110<sup>6</sup> ( $z = 0.007579$ ),  $1''$  corresponds to 160 pc. This value is slightly different ( $-131 \text{ km s}^{-1}$ ) than the velocity (with respect to the cosmic microwave background [CMB]) derived from stellar Mg I *b* absorption lines (Nelson & Whittle 1995). However, none of our conclusions would be changed if we adopted this alternate value, and we note that previous publications on NGC 2110 have used a variety of redshift estimates.

All spectral fits include absorption through our Galaxy using  $N_{\text{H,Gal}} = 1.76 \times 10^{21} \text{ cm}^{-2}$  (Dickey & Lockman 1990). Errors quoted in this paper are 90% confidence for one parameter of interest (i.e.,  $\chi^2_{\text{min}} + 2.7$ ), unless otherwise stated. When distinguishing between different model fits to the data, we consider *F*-statistic results, and we adopt a threshold of 95% for a significant improvement in the fitting statistic.

## 2. OBSERVATIONS AND ANALYSIS

NGC 2110 was observed on 2000 April 22 (ObsID 883) with the S3 chip of the *Chandra*/ACIS charge-coupled device (CCD) camera. The observation was made in the VFaint mode, using a 128 row subarray (giving a  $1' \times 8'$  field of view), in order to reduce the CCD frame time to 0.44 s and lessen the effect of pileup. The source was positioned near the standard aim point on the S3 chip. The data were reprocessed using CIAO v3.2.2 with the CALDB v3.10 calibration database to create a new level 2 events file filtered for the *ASCA* grades 0, 2, 3, 4, 6, and with the  $0''.5$  pixel randomization removed. To check for intervals of high particle background, light curves were extracted for the ACIS-S3 chip, excluding the core. No periods of high background were found, and so the entire 44.3 ks observation was used in our analysis.

The *Chandra* observation of the nucleus of NGC 2110 is affected by pileup to the extent that the accurate modeling of the nuclear spectrum is significantly hampered. We demonstrate this using two independent methods. First, the counts extracted from a source-centered circle of radius  $1''.23$  (2.5 pixels) in the 0.5–7 keV energy band gave 0.38 counts per frame, for which a pileup fraction of  $\sim 15\%$ – $20\%$  is predicted using the PIMMS software. As a second independent test, we estimated the spatial extent of the pileup by producing images consisting of grade 7 events, which are largely produced by photon pileup. Inspection of the images shows the pileup to be concentrated within the central two pixels.

In addition to the X-ray data, we utilize multiwavelength archival data from observations of NGC 2110 with the VLA and *HST*. NGC 2110 was observed with the VLA in its A configuration at 5.0 GHz for  $\sim 70$  minutes on 1994 February 27. We extracted the data from the archive and calibrated them in the standard manner with AIPS, using 3C 48 as the primary flux calibrator and J0541–056 as the phase calibrator. The quality of the final image of NGC 2110 was improved using one iteration of phase self-calibration, and the noise level ( $40 \mu\text{Jy beam}^{-1}$ ) approaches the expected thermal limit. We verified by examining B-configuration archival data (taken on 1985 May 27) that the A-configuration data spatially samples all the detectable structure. The reduced *HST* data were provided to us by J. Mulchaey, and Mulchaey et al. (1994) should be consulted for a detailed description of the data-reduction methods used.

## 3. RESULTS

### 3.1. Soft-Band Extended Emission

#### 3.1.1. Small-Scale Emission

Figure 1 shows a Gaussian-smoothed ( $\sigma = 0''.5$ ) tri-band *Chandra* ACIS-S image of NGC 2110. Here, red corresponds to 0.5–1 keV photons, green to 1–1.5 keV photons, and blue to 4–5 keV photons. We have deliberately emphasized the 0.5–1 keV photons, in order to contrast them with the 4–5 keV photons. Figure 1 shows that the extended X-ray emission to the north of the nucleus, first reported by Weaver et al. (1995) with *ROSAT*, is dominated by

<sup>6</sup> Taken from the SIMBAD database.

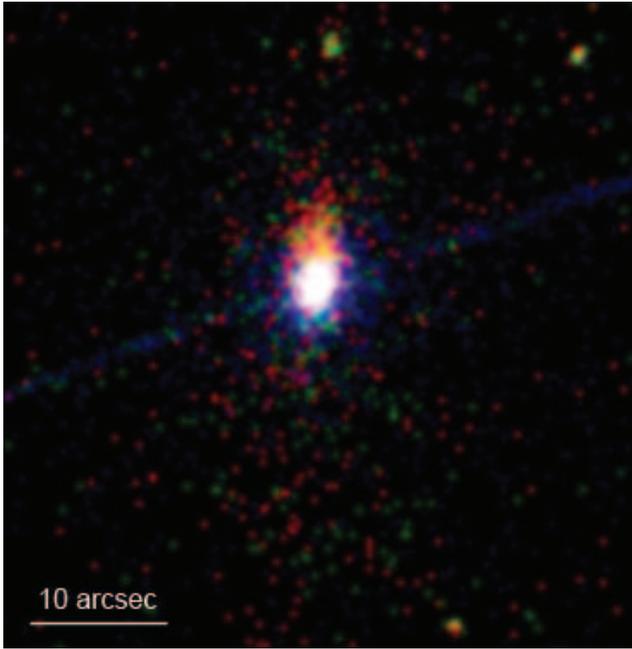


FIG. 1.— Tri-color (red: 0.5–1 keV, green: 1–1.5 keV, blue: 4–5 keV) ACIS-S image of NGC 2110, showing soft extended emission toward the north of the nucleus. The image has been binned to  $0''.25 \text{ pixel}^{-1}$ , and subsequently smoothed with a Gaussian of  $\sigma = 0''.5$ .

soft X-ray photons. This is in contrast to the nucleus, which shows both soft and hard ( $>4$  keV) X-ray emission. There is a hint of 1–1.5 keV X-ray emission just south of the nucleus, which may be emission similar to the northern feature, but obscured by the small-scale gas disk discussed by Gallimore et al. (1999).

Figure 2 shows a 0.5–1.5 keV *Chandra* image of NGC 2110, binned to  $0''.25 \text{ pixel}^{-1}$  and subsequently smoothed with a Gaussian of  $\sigma = 0''.5$ , with contours from *HST* [O III] (Mulchaey et al. 1994) and 5 GHz VLA radio observations overlaid. The relative astrometry of the X-ray, radio, and [O III] images have been adjusted to make the unresolved cores spatially coincident. Fig-

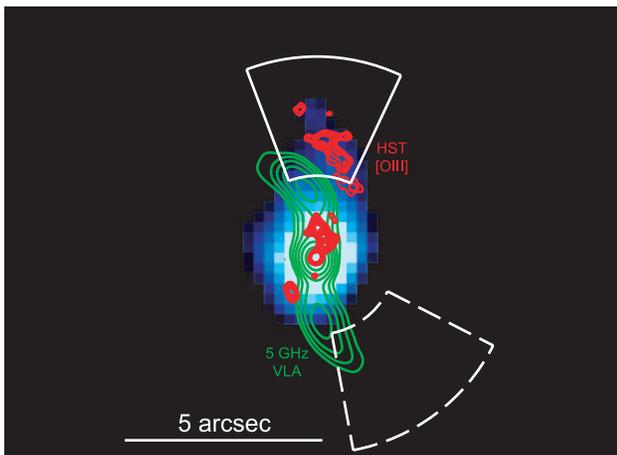


FIG. 2.— 0.5–1.5 keV *Chandra* image of NGC 2110 (blue) with VLA (green) and *HST* [O III] (red) contours overlaid. The *Chandra* image has been binned to  $0''.25 \text{ pixel}^{-1}$ , and subsequently smoothed with a Gaussian of  $\sigma = 0''.5$ . The relative astrometry of the X-ray, radio, and [O III] images have been adjusted to make the unresolved cores spatially coincident. The white pie section marks the region used to extract the spectrum of the extended X-ray emission, and the dashed white section marks the background sampled.

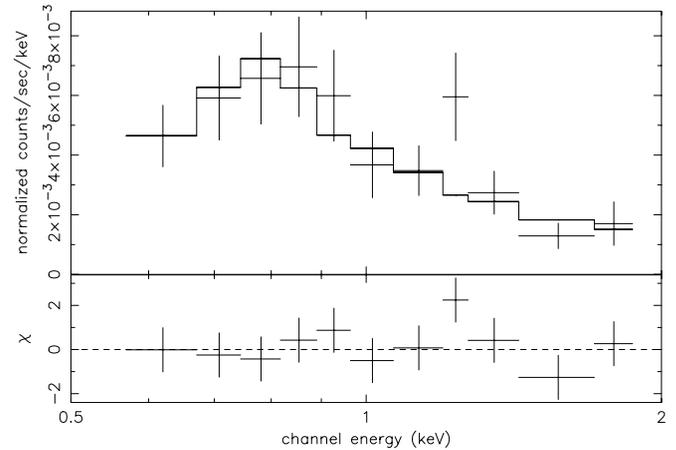


FIG. 3.— 0.5–2 keV X-ray counts spectrum and contributions to  $\chi^2$  of the extension north of the nucleus. The model fit shown is the sum of two thermal components of temperature 0.34 and 5.5 keV, with abundance fixed at solar.

ure 2 shows that the X-ray and radio emission extend along a similar position angle north of the nucleus, but are not spatially coincident: the X-ray emission lies beyond the northern edge of the radio jet, before the jet begins to diverge away from it. However, the X-ray emission is coincident with the S-shaped [O III] emission. Our results may support a common origin for the soft X-ray and [O III] emission, which we discuss more fully in § 4.

We extracted the X-ray spectrum of the extended X-ray emission, using a source-centered annulus of inner radius  $2''$  and outer radius  $5''$  between the position angles  $-25^\circ$  and  $20^\circ$ . In order to reduce any possible contamination from the wings of the unresolved nucleus, we adopted a matched background extraction region of identical radii between the position angles  $190^\circ$  and  $235^\circ$ . The source and background extraction regions are marked in Figure 2. Imaging showed possible contamination from the nucleus at hard energies ( $>2$  keV), and so we restricted our spectral analysis to 0.5–2 keV. The data were grouped to 20 counts per bin. The X-ray spectrum is shown in Figure 3.

We attempted to fit a series of models to the X-ray spectrum of the extended emission. Guided by the prevalence of jet-induced shock heating in radio galaxies (e.g., Kraft et al. 2003), we initially fitted the X-ray spectrum of the extension with a single-temperature, collisionally ionized thermal plasma (APEC) model, with abundance fixed at solar, and modified by Galactic absorption. However, the fit was poor ( $\chi^2 = 35.4$  for 9 dof), due to the model greatly overproducing Fe L lines between 0.7 and 1.2 keV for the best fitting  $kT = 0.62$  keV. The fit was significantly improved ( $\chi^2 = 6.6$  for 8 dof) when the abundance was left to be a free parameter, although its value tended to an implausibly low value of  $(1.01^{+3.23}_{-1.01}) \times 10^{-2}$  solar. An acceptable fit ( $\chi^2 = 8.3$  for 7 dof) was achieved with the combination of two thermal models of abundances fixed at solar; one of temperature  $kT_1 = 0.34^{+0.11}_{-0.07}$  keV, and the other significantly hotter, with a temperature  $kT_2 \sim 5.5$  keV and a normalization 3 times greater than the cool component. Figure 3 shows the counts spectrum and  $\chi^2$  residuals for this model.

We fitted the X-ray spectrum with an alternative model consisting of an unabsorbed power law and multitemperature blackbody emission from an accretion disk. This model is appropriate for the scattering of nuclear radiation by a population of electrons. The best-fitting nuclear spectrum is found from the *Chandra* HETGS data (D. A. Evans et al. 2007, in preparation).

TABLE 1  
PARAMETERS OF MODEL FITS TO THE 0.5–2 keV X-RAY EMISSION  
NORTH OF THE NUCLEUS

Model (1)	Description (2)	Parameters (3)	$\chi^2/\text{dof}$ (4)
N1.....	TH	$kT = 0.62 \pm 0.05$ $Z = 1$ (f) $\text{norm} = (7.55 \pm 1.37) \times 10^{-6}$	35.4/9
N2.....	TH+TH	$kT_1 = 0.34^{+0.11}_{-0.07}$ keV $Z_1 = 1$ (f) $\text{norm}_1 = (7.05^{+2.50}_{-2.86}) \times 10^{-6}$ $kT_2 = 5.5^{+5.9}_{-2.9}$ keV $Z_2 = 1$ (f) $\text{norm}_2 = (1.98^{+2.09}_{-0.65}) \times 10^{-5}$	8.31/7
N3.....	DISKBB+PL	$kT = 0.19$ keV (f) $\text{norm}_{\text{diskbb}} = 2.39 \pm 0.32$ $\Gamma = 1.6$ (f) $\text{norm}_{\text{PL}} = 1.72 \times 10^{-6}$ (t)	14.8/10

NOTE.— Col. (2): TH=collisionally ionized plasma (APEC) model, DISKBB=disk blackbody, PL=power law. Col. (3): (f) indicates parameter was frozen, (t) indicates parameter was tied to blackbody normalization. Normalization quoted is in units of  $10^{-14}(1+z)^2 \int n_e n_p dV / 4\pi D_L^2$  (thermal),  $[(R_{\text{in}}/\text{km})/(D_L/10 \text{ kpc})]^2 \cos(i)$  (disk blackbody), and photons  $\text{cm}^{-2} \text{s}^{-1} \text{keV}^{-1}$  at 1 keV (power law).

We emphasize that the intrinsic absorption of the nucleus ( $N_{\text{H}} \sim 4 \times 10^{22} \text{ cm}^{-2}$ ), which we associate with the central engine of NGC 2110, is not included in our model fit to the extended region. With the power-law photon index, disk temperature, and relative normalizations fixed at their HETGS values, we found an adequate fit to the extended-emission spectrum ( $\chi^2 = 14.8$  for 10 dof), although there were noticeable negative residuals

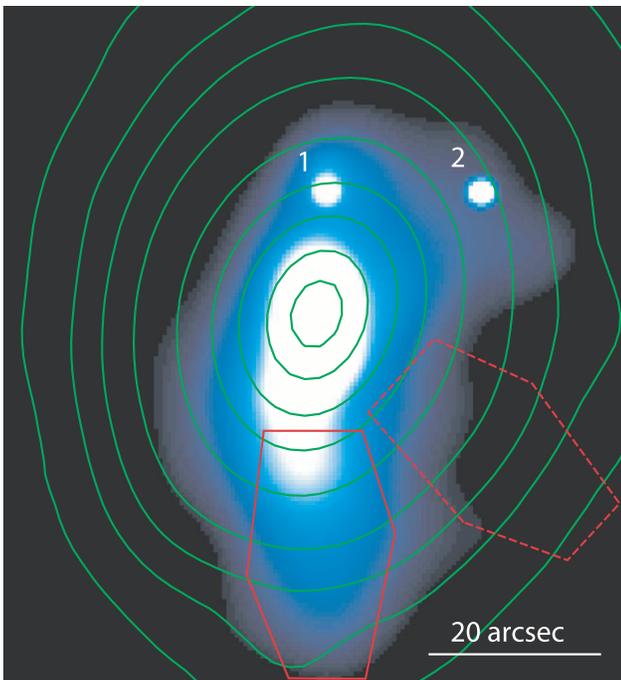


FIG. 4.— Adaptively smoothed 0.5–1 keV *Chandra* ACIS-S image of NGC 2110, with optical DSS contours overlaid. The X-ray morphology does not trace the stellar population, with a noticeable X-ray extension south of the nucleus. Two point sources to the north of the nucleus, marked 1 and 2, are also observed. The large solid red box south of the nucleus marks the region used to extract the spectrum of the extended X-ray emission, and the dashed box marks the background extraction region.

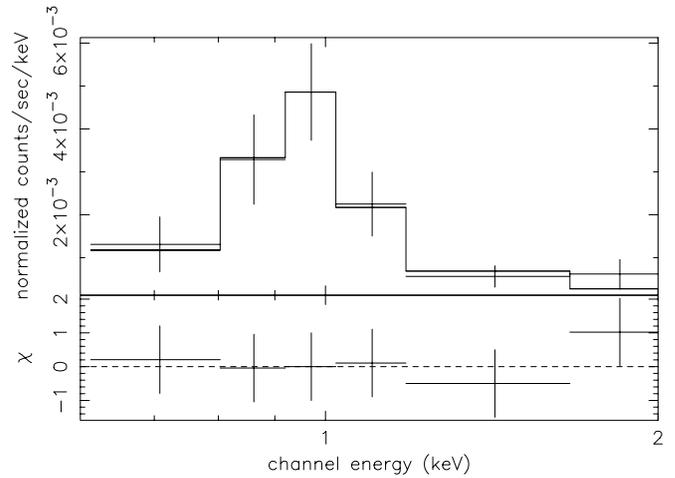


FIG. 5.— 0.5–2 keV X-ray counts spectrum and contributions to  $\chi^2$  of the large-scale extension south of the nucleus. The model fit shown is the sum of a power law of photon index 1.6 and thermal emission of temperature 0.82 keV and abundance fixed at solar.

below  $\sim 0.8$  keV. Table 1 summarizes the results of our spectral fitting.

Finally, given the observations of photoionized line emission in high-energy-resolution X-ray observations of certain nearby Seyfert galaxies (e.g., Sako et al. 2000; Sambruna et al. 2001; Ogle et al. 2003), it is interesting to consider photoionized gas as the emission mechanism for the extended X-ray emission in NGC 2110. However, we note that the detection of photoionized line emission is hampered here due to the relatively poor quality of the *Chandra* ACIS spectrum of the extended emission ( $\sim 200$  counts), together with its CCD energy resolution. In short, we cannot distinguish spectrally between these plausible models for the extended X-ray emission. However, the 0.5–2.0 keV X-ray luminosity of  $\sim 3 \times 10^{39} \text{ ergs s}^{-1}$  is relatively insensitive to the spectral modeling, and this constraint allows us to investigate simple arguments that demonstrate the energetic viability of photoionization (see § 4.4).

### 3.1.2. Large-Scale Emission

On larger ( $\sim 30''$ ) scales, there is a noticeable asymmetry between the soft X-ray emission located north and south of the nucleus. Figure 4 shows an adaptively smoothed 0.5–1 keV image of NGC 2110, with contours from an optical Digitized Sky Survey (DSS) image overlaid. The X-ray morphology does not trace the stellar population. We extracted the 0.5–2 keV spectrum of the large-scale southern extension from the source and background regions marked in red in Figure 4. The spectrum, grouped to 20 counts per bin, is shown in Figure 5. A model fit to a single-component thermal APEC model of temperature  $0.96^{+0.21}_{-0.17}$  keV and abundance fixed at solar provided a good fit to the spectrum ( $\chi^2 = 1.4$  for 4 dof). The parameters of the best-fitting model are shown in Table 2.

TABLE 2  
PARAMETERS OF MODEL FITS TO THE 0.5–2 keV X-RAY EMISSION  
SOUTH OF THE NUCLEUS

Model	Description	Parameters	$\chi^2/\text{dof}$
S1.....	TH	$kT = 0.96^{+0.21}_{-0.17}$ keV $Z = 1$ (f) $\text{norm}(4.23^{+1.25}_{-1.21}) \times 10^{-6}$	1.4/4

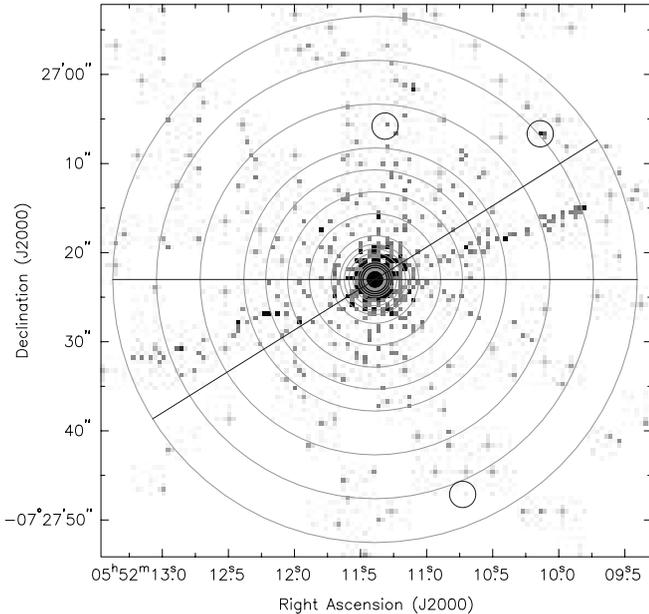


FIG. 6.—4–4.5 keV *Chandra* image of NGC 2110, showing the regions used to extract the radial profile. The radial profile, sampled from the concentric annuli shown in gray, extends from  $1''$  to  $30''$ , with background taken from a large source-centered annulus of inner radius  $37''$  and outer radius  $50''$ . The profile excludes pie slices between the position angles  $90^\circ$ – $122^\circ$  and  $270^\circ$ – $302^\circ$ , in order to exclude emission from the frame transfer streak, and circles that mask emission from nonnuclear point sources.

### 3.2. Associated Point Sources?

Figure 4 also shows two point sources to the north of the nucleus of NGC 2110. The sources, marked (1) and (2), are most prominent at energies  $\lesssim 1.5$  keV. We inspected an archival *HST* continuum image and found no optical counterparts to these sources. There are 61 and 74 background-subtracted 0.5–7 keV counts in each, respectively. Their 0.5–10 keV X-ray spectra are well described by unabsorbed power laws of photon index  $1.6 \pm 0.4$  and  $1.2 \pm 0.3$ , respectively, with unabsorbed fluxes  $(1.4 \pm 0.6) \times 10^{-14}$  and  $(2.2 \pm 0.5) \times 10^{-14}$  ergs  $s^{-1} cm^{-2}$ . We performed queries with the NED and SIMBAD databases, and found no known sources in their vicinity. We used the soft-band  $\log N$ – $\log S$  relationship for the *Chandra* Deep Fields (Tozzi et al. 2001; Bauer et al. 2004) to estimate the number of background X-ray sources expected within a circle of radius  $50''$  centered on the nucleus, and found this number to be  $\sim 0.5$ – $1.5$ , which is approximately consistent with the two sources we detected. While we cannot determine the distance to these sources, if they were at the redshift of NGC 2110, their 0.5–10 keV unabsorbed luminosities would be  $(1.6 \pm 0.5) \times 10^{39}$  and  $(2.7 \pm 0.7) \times 10^{39}$  ergs  $s^{-1}$ , respectively, which would make them ultraluminous X-ray sources (ULXs; e.g., Swartz et al. 2004).

### 3.3. Hard-Band X-Ray Emission

We attempted to confirm that the hard ( $>3$  keV) X-ray emission is unresolved by extracting and modeling the radial profile of NGC 2110 from the ACIS-S3 chip. The results described here are for fits to 4–4.5 keV events, but are consistent with the results for any  $>3$  keV events. The 4–4.5 keV radial profile was centered on the nucleus, and excluded pie slices between the position angles  $90^\circ$ – $122^\circ$  and  $270^\circ$ – $302^\circ$ , in order to exclude emission from the frame transfer streak, and circles that mask any emission from the point sources discussed in § 3.2. The radial profile extended from  $1''$  (so as to exclude events known to be affected

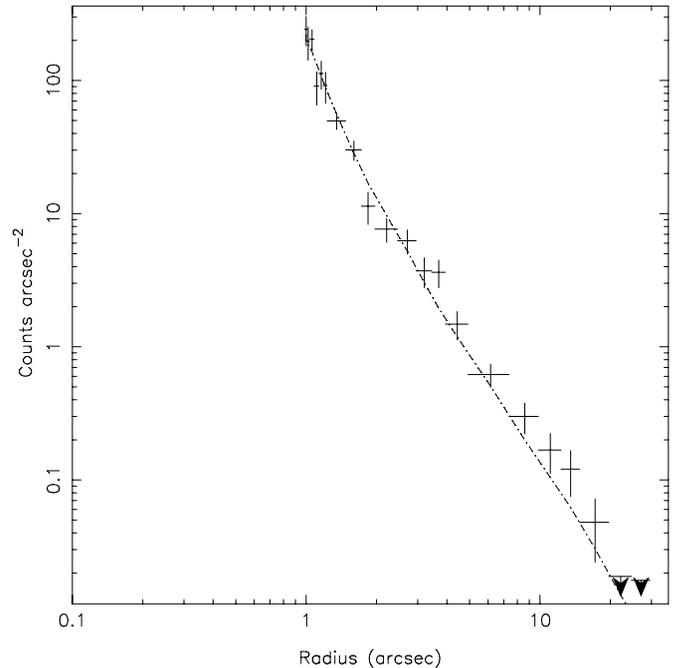


FIG. 7.—Background-subtracted 4–4.5 keV radial surface brightness profile of NGC 2110, shown with a model *Chandra* PSF.

by photon pileup), out to  $30''$ , with background taken from a large source-centered annulus of inner radius  $37''$  and outer radius  $50''$ . Figure 6 shows the regions used to extract the radial profile. A point-spread function (PSF) of  $\Gamma = 1.6$  was modeled using ChaRT and MARX, and was convolved with a Gaussian of rms width  $0''.35$  in order to provide a good match with the radial profile data close to the core. The radial profile is shown in Figure 7.

An acceptable fit to the radial profile ( $\chi^2 = 21.2$  for 20 dof) was achieved with the model of the *Chandra* PSF. However, at distances  $\gtrsim 8''$ , there is a small excess of observed counts with respect to the model PSF. We inspected the 4–4.5 keV image, and found a slight excess of counts toward the northwest of the nucleus. While the origin of the excess emission is unclear, it is not spatially coincident with the soft X-ray emission, which has a north-south orientation. We speculate that, given the readout direction of the ACIS instrument, the apparent excess hard emission is due to the combination of serial and parallel charge transfer inefficiency (M. Nowak 2006, private communication).

## 4. ORIGIN OF NORTHERN EXTENDED EMISSION

### 4.1. X-Ray Emission from the Radio Jet?

Many low-power radio galaxies show X-ray jets that are inferred to have a synchrotron origin, given the good spatial agreement between the X-ray and radio emission. Since the X-ray and radio emission in NGC 2110 are not spatially coincident (Fig. 2), it seems unlikely that the X-ray emission is directly associated with the radio jet.

### 4.2. Shock-heated Gas?

An alternative model is that clouds of gas have been ablated and shock heated by the radio jet. We test this model by comparing the pressures of the radio jet with the X-ray and [O III] emission, in order to determine if the pressure of the jet is sufficient to drive a shock into the gas. We adopt the standard assumption of equipartition between the magnetic field and particles to

estimate the minimum pressure of the radio jet. The 5 GHz flux density of the brightest part of the jet is 40 mJy. We model the jet as an ellipsoid of dimensions  $1''.4 \times 0''.2 \times 0''.2$  (volume  $3.5 \times 10^{60}$  cm<sup>3</sup>). We assume the electron energy distribution extends from a Lorentz factor  $\gamma_{\min} = 2$  up to  $\gamma_{\max} = 10^5$ , with an electron energy index  $p = 2.4$  (the results have only a weak dependence on the choice of  $\gamma$ ; see Hardcastle et al. 2004). Under the simple assumption that the jet is in the plane of the sky, and that relativistic beaming is unimportant, the minimum pressure is  $\sim 10^{-10}$  Pa, and the minimum energy is  $\sim 10^{52}$  ergs.

In order to estimate the pressure of the hot gas, we adopt a two-temperature plasma as the best-fitting spectral model (model N2 in Table 1). We adopt a cylindrical region of radius  $0''.37$  and length  $1''.9$  (volume =  $9.9 \times 10^{61}$  cm<sup>3</sup>), located where the X-ray emission is brightest and adjacent to the radio emission. The APEC normalization is  $10^{-14}(1+z)^2 \int n_e n_p dV / 4\pi D_L^2$ , where  $n_e$  and  $n_p$  are the electron and proton number densities (cm<sup>-3</sup>), and  $D_L$  is the luminosity distance (cm). From this, we find the pressure of the gas ( $P \sim 2.2n_p k_B T$ ) to be  $\sim 3 \times 10^{-9}$  Pa, and its energy  $\sim 3 \times 10^{54}$  ergs. We calculate the sound speed of the gas to be  $\sim 8 \times 10^5$  m s<sup>-1</sup>, and find the characteristic sound crossing time to be  $\sim 2 \times 10^5$  yr. The rate of energy transfer from the jet is therefore  $\sim 5 \times 10^{41}$  ergs s<sup>-1</sup>. Finally, Ferruit et al. (1999) used *HST* FOC observations to estimate the temperature and density of the optical line-emitting gas. They derive an [O III] temperature of  $\sim 17,400$  K and, using [S II] line ratios, find a density  $\sim 1200$  cm<sup>-3</sup>, implying a pressure  $\sim 6 \times 10^{-10}$  Pa.

The pressure of the X-ray gas is approximately 30 times greater than the minimum internal pressure in the radio jet. However, for a shock to be driven into the gaseous atmosphere, the pressure of the radio jet must be comparable to that of the X-ray gas. On the other hand, the radio pressure calculated is formally a *minimum* pressure, and the underlying assumption that the energies in radiating electrons and magnetic fields are close to equipartition may not be applicable in the jet of NGC 2110. As an example, the diffuse radio structures that are apparent in FR I-type radio galaxies require large departures from equipartition in order to maintain pressure balance with the gaseous atmosphere (e.g., Croston et al. 2003), and we speculate that a similar situation may be occurring in NGC 2110. Indeed, we suggest that departures from equipartition might be an important difference between the “frustrated,” diffuse jets and lobes observed in Seyfert galaxies such as NGC 2110 and FR I-type radio galaxies, and the powerful jets observed in FR II-type radio galaxies radio galaxies.

The spatial offset between the radio, and [O III] and X-ray emission can be accounted for in this simple model of shock heating. As has been observed in the radio galaxy Centaurus A (e.g., Kraft et al. 2003), radio ejecta can be driven into a gaseous environment, giving rise to a cap of X-ray-emitting shock-heated gas that surrounds the radio emission, but is spatially distinct from it. We suggest that a similar mechanism may be applicable to NGC 2110, where the shock is detached.

An alternative possibility is that there is a considerable ram velocity of the jet into the gas, i.e., that the observed X-ray-emitting gas is already present and has not been shock heated. This obviates the need for a balance between the ambient pressure of the gas and the internal pressure of the jet. However, this possibility has two problems. First, the origin of the X-ray-emitting gas would again be unclear, and second, the similarities in the spatial distribution of radio and X-ray gas would be merely coincidental.

A final interpretation within the shocked-gas scenario is that EUV photons from the bow shock of the radio lobe are sufficient to photoionize the ambient environment (e.g., Dopita &

Sutherland 1996), giving rise to [O III] and X-ray emission, as discussed by, e.g., Ferruit et al. (1999). However, this model is highly unlikely to be able to generate X-ray-emitting gas with temperatures  $\geq 0.1$  keV, since the emissivity of the required strong shock is extremely low.

We conclude that a simple model of shock-heated multiphase gas clouds is an energetically viable mechanism for producing the observed optical emission lines and two-temperature X-ray emission.

#### 4.3. Scattered Nuclear Radiation?

A second possible scenario to explain the extended soft X-ray emission is that nuclear radiation is being scattered by a population of electrons. We evaluate this model by presenting simple energetic arguments based on the comparison of the scattered and direct nuclear fluxes, following Weaver et al. (1995).

We use *Chandra* HETGS observations (D. A. Evans et al. 2007, in preparation) to model the direct nuclear emission as the sum of a power law and multiple-component disk blackbody emission, modified by an absorbing column  $N_H \sim 4 \times 10^{22}$  cm<sup>-2</sup>. We adopt the best-fitting spectral model of the extended X-ray emission to consist of the same spectral components, but different normalizations, and without the intrinsic absorption (model N3 in Table 1).

The ratio of the unabsorbed power-law fluxes of the scattered and “direct” components is  $\sim 10^{-4}$ . We assume the nuclear radiation to be emitted isotropically, and scattered into a cone of opening angle  $45^\circ$  ( $\Omega = 0.48$ ). The optical depth to Compton scattering is then  $\sim 5 \times 10^{-3}$  ( $N_{\text{scat}} \sim 10^{22}$  cm<sup>-2</sup>). Assuming the line-of-sight scattering radius to be  $\sim 240$  pc, we find the number density of electrons to be  $\sim 10$  cm<sup>-3</sup>. This gives a mass of  $\sim 10^7 M_\odot$ , which seems very large.

It is therefore unlikely that the scattering of isotropic nuclear radiation can power the extended X-ray emission. On the other hand, if the nuclear radiation were emitted anisotropically into a narrow cone, it has been argued by Weaver et al. (1995) that electron scattering may remain a viable mechanism, and our findings may be consistent with this interpretation. However, modeling of the mid-infrared SED of NGC 2110 (e.g., Fritz et al. 2006) suggests that the dust in the “torus” is heated essentially isotropically by the optical-UV continuum, and we see no reason for a departure from isotropic emission at X-ray wavelengths. Finally, we note that a test of this model would be to observe polarized optical light in this region.

#### 4.4. Photoionization

In this scenario, pure photoionization of gas by the nuclear radiation is responsible for the extended soft X-ray emission. The detection of photoionized line emission is hampered in our relatively low signal-to-noise ratio, CCD-quality spectrum, although we can still use the superior spatial resolution of *Chandra* to provide new constraints on the energetic viability of this model, using the measured X-ray luminosity of the emission, and the simple energetics arguments first proposed by Weaver et al. (1995). The luminosity in emission lines for a sphere of radius  $r$  is

$$L_{\text{lines}} = \frac{4\pi}{3} n^2 r^3 j(\xi), \quad (1)$$

where  $n$  is the electron number density (cm<sup>-3</sup>),  $j(\xi)$  is the emissivity of the photoionized gas (ergs cm<sup>3</sup> s<sup>-1</sup>), and  $\xi$  is the ionization parameter (ergs cm s<sup>-1</sup>) (see Weaver et al. 1995 for discussions). Assuming a sphere of radius  $1''.5$ , and using the integrated 0.5–2 keV unabsorbed X-ray luminosity of the extended region of

$\sim 3 \times 10^{39}$  ergs  $s^{-1}$ , we find a number density of  $1.35 \text{ cm}^{-3}$ . For our calculations, we use the definition of the ionization parameter

$$\xi = \frac{L_s}{nR^2}, \quad (2)$$

where  $L_s$  is the luminosity of the nucleus and  $R$  is the distance from the nucleus to the photoionized gas. For  $\xi \sim 100$  ergs  $\text{cm s}^{-1}$  (see arguments by Weaver et al. 1995) and a distance  $R = 3.73$ , the luminosity of the nucleus required to photoionize the gas is in excess of  $10^{44}$  ergs  $s^{-1}$ . This is a factor  $\sim 20$  greater than the unabsorbed nuclear luminosity. We conclude that photoionization is a viable mechanism for producing the extended emission (or at least a fraction of it), although as first noted by Weaver et al. (1995) the nuclear radiation may have to be emitted anisotropically if the extended emission were attributed solely to photoionization.

We note that the electron number density we calculate for the photoionized gas is approximately consistent with that which we derived for the collisionally ionized case (§ 4.2). This demonstrates that we are not orders of magnitude off in our derivation of the properties of the photoionized gas, despite our simple assumptions.

Finally, if the extended soft X-ray emission were gas photoionized by the nucleus, then one would expect to observe emission lines from photoionized plasma in high signal-to-noise ratio, high-resolution X-ray spectra. Observations with *Chandra* HETGS in AO-3 and AO-4 (D. A. Evans et al. 2007, in preparation) will help test this further, although even these data are not sufficient in terms of numbers of counts to permit a spectral analysis of the extended region. Preliminary results indicate that the soft-band nuclear X-ray emission is not dominated by photoionized emission lines, unlike that of, e.g., NGC 1068 (Ogle et al. 2003). The emission is instead consistent with direct emission from the nucleus that intersects cold obscuring material, possibly in the form of a torus. However, Guainazzi & Bianchi (2006) claim to detect a radiative recombination continuum (RRC) from O VII in *XMM-Newton* reflection grating spectrometer (RGS) observations of NGC 2110, which may be consistent with photoionization. The energetics arguments in this section simply serve to demonstrate that photoionization is energetically viable.

## 5. ORIGIN OF SOUTHERN EXTENDED EMISSION

The extended X-ray emission to the south of the nucleus is interestingly not associated with any radio, optical line, or CO emission, making its interpretation difficult. We investigate several models for the extended emission here.

First, we calculate the energy of the extended emission, by modeling the emission region as a cylinder of radius  $6''$  and length  $30''$  (volume =  $4 \times 10^{65} \text{ cm}^3$ ). Using the same arguments as in § 4.2, we find the energy of the gas to be  $10^{55}$  ergs. We speculate that if the extended X-ray emission may be related to the nucleus in some form of outflow.

Alternatively, some Seyfert galaxies show extended X-ray emission on scales of several kpc, coincident with an extended narrow line region (ENLR; e.g., Bianchi et al. 2006). However, owing to the generally poor signal-to-noise ratio observations, the X-ray emission mechanism is uncertain, except for certain nearby Seyfert galaxies in which the extended emission is associated with photoionization. In NGC 2110, no ENLR is detected either to the north or south of the nucleus.

A final possibility is that the X-ray emission is similar to the ‘‘anomalous arms’’ seen in a *Chandra* observation of NGC 4258 (Wilson et al. 2001). For NGC 4258, the X-ray emission is

associated with gas from an inclined ( $i = 64^\circ$ ) disk, postshocked by gas driven down onto the disk by the radio ejecta. For NGC 2110, this model does not explain why only X-ray emission is observed, rather than being accompanied by the  $H\alpha$  emission seen in NGC 4258. Further, the large (projected) difference in the sizes of the radio and X-ray emission are not easily explained unless the galaxy is highly inclined, contrary to observations.

We conclude that the origin of this extended X-ray emission is unclear. However, its prominence to the south of the nucleus is consistent with its being on the portion of the disk nearest the observer, although we cannot definitively determine whether or not it is obscured by the disk.

## 6. COMPARISON WITH NGC 1068

NGC 1068 is a nearby ( $z = 0.003793$ ,  $D_L = 16.3$  Mpc) Sb spiral galaxy, and is often regarded as the prototypical example of a Seyfert 2 nucleus. An imaging observation with *Chandra* (Young et al. 2001) shows that the X-ray morphology of the source consists of a bright, compact ( $\sim 170$  pc) central region, a ( $\sim 550$  pc) extended region that points to the northeast, and a further (several kpc) region that appears to be associated with the spiral arms of the galaxy. The  $\sim 550$  pc northeast X-ray region in NGC 1068 is spatially coincident with both the [O III] emission and radio ejecta, unlike in NGC 2110, in which the X-ray and radio emission are spatially distinct. This may be evidence that in NGC 2110 the interaction of the radio jet and its environment produces a cap of shock-heated gas, unlike in NGC 1068. Young et al. (2001) argue that the northeastern region in NGC 1068 is not associated with starburst activity. In addition, Young et al. (2001) find that a single, collisionally ionized, thermal plasma model provides a poor fit to the spectrum of this region, unless the abundance tends to an implausibly low value. They instead fit the spectrum with the combination of two bremsstrahlung continua, of temperature  $\sim 0.4$  and  $2.8$  keV, plus a series of emission lines, with line energies consistent with results from *XMM-Newton* RGS observations (Kinkhabwala et al. 2002). Detailed modeling of the high-resolution X-ray spectrum of NGC 1068 (Brinkman et al. 2002; Kinkhabwala et al. 2002; Ogle et al. 2003) shows that the emission lines are likely to be associated with a photoionized plasma. This is consistent with UV results from *HST* Space Telescope Imaging Spectrograph (STIS) emission-line mapping observations (Groves et al. 2004). NGC 1068 also shows evidence of a region of optical line-emitting gas at a distance  $\sim 200$  pc from the nucleus, which Kraemer & Crenshaw (2000) interpret as having been *locally* enhanced by the interaction of emission-line knots with the lower density interstellar medium. Although it is plausible that a similar process may produce the optical line emission in NGC 2110, we consider it unlikely that it is able to produce the high-temperature X-ray gas observed, since the radiative coupling of  $\sim$ keV gas is extremely low. In summary, although the relatively poor quality of the CCD spectrum of the extended emission in NGC 2110 does not enable us to detect lines that might be consistent with emission from a photoionized plasma, the differing relative spatial distributions of X-ray, radio, and [O III] emission in NGC 2110 and NGC 1068 may suggest that different physical processes could produce the extended circumnuclear X-ray emission in Seyfert galaxies. This is important to keep in mind when interpreting spatially unresolved X-ray data of distant Seyfert galaxies, e.g., those in deep fields.

## 7. CONCLUSIONS

We have presented results from a *Chandra* observation of the circumnuclear environment of the Seyfert 2 galaxy NGC 2110,

and have combined them with VLA radio and *HST* [O III] data, in order to constrain the physical processes present. To summarize our conclusions:

1. We detect resolved, soft-band X-ray emission  $\sim 4''$  north of the nucleus. The X-ray emission is spatially coincident with [O III] optical line emission, but not with the radio emission, and instead appears to lie north of the radio ejecta.

2. We consider several different models that may account for the extended northern X-ray emission. We find that it is possible that the radio jet has driven a strong shock through a series of multiphase clouds of gas. We demonstrate that nuclear radiation scattered by a population of electrons may account for the extended X-ray emission, if the nucleus emits its radiation anisotropically. Finally, we show that photoionization of ambient gas by the nucleus may be sufficient to be responsible for the extended emission; high signal-to-noise ratio *Chandra* grating observations can directly test this model. Although energetically all of these mechanisms are plausible, we consider shock emission from the radio jet interacting with the ISM to be the favored model because of the spatial properties of the multiwavelength emission.

3. We detect large-scale ( $\sim 30''$ ) extended X-ray emission south of the nucleus in NGC 2110, which we model as a collisionally ionized plasma of temperature  $\sim 0.9$  keV. The origin of this emission is unclear, although its prominence on the southern side is consistent with the overall geometry of the galaxy.

4. A comparison of the spatial distributions of X-ray, radio, and [O III] emission of NGC 2110 and the prototypical Seyfert 2 galaxy NGC 1068 leads us to suggest that different physical processes could produce the extended circumnuclear X-ray environments in Seyfert galaxies.

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