X-Ray Nuclei in Radio Galaxies: Exploring the Roles of Hot and Cold Gas Accretion

D.A. Evans,¹ M.J. Hardcastle² and J.H. Croston²

We present results from *Chandra* and *XMM-Newton* spectroscopic Abstract. observations of the nuclei of z < 0.5 radio galaxies and quasars from the 3CRR catalog, and examine in detail the dichotomy in the properties of low- and highexcitation radio galaxies. The X-ray spectra of low-excitation sources (those with weak or absent optical emission lines) are dominated by unabsorbed emission from a parsec-scale jet, with no contribution from accretion-related emission. These sources show no evidence for an obscuring torus, and are likely to accrete in a radiatively inefficient manner. High-excitation sources (those with prominent optical emission lines), on the other hand, show a significant contribution from a radiatively efficient accretion disk, which is heavily absorbed in the X-ray when they are oriented close to edge-on with respect to the observer. However, the low-excitation/high-excitation division does not correspond to the FRI/FRII division: thus the Fanaroff-Riley dichotomy remains a consequence of the interaction between the jet and the hot-gas environment through which it propagates. Finally, we suggest that accretion of the hot phase of the IGM is sufficient to power all low-excitation radio sources, while high-excitation sources require an additional contribution from cold gas that in turn forms the cold disk and torus. This model explains a number of properties of the radio-loud active galaxy population, and has important implications for AGN feedback mechanisms.

1. Unified Models and Radio-Loud AGN: The Excitation Dichotomy

In standard AGN models, efficient disk accretion of cold matter on to the central supermassive black hole provides the radiation field that photoionizes the optical broad-line region (BLR) and narrow-line region (NLR) and gives rise to X-ray emission via Compton scattering. Without radiatively efficient accretion via the disk, none of these standard features of such an AGN would be observed. Unified models propose that a direct view of the BLR and the optical continuum may be obscured (e.g., in Seyfert 2s) by a dusty 'torus': but in this case the torus re-radiates strongly in the mid-IR band, so that the presence of a luminous AGN can still be inferred.

By analogy with radio-quiet objects, we would expect that face-on radioloud objects (the broad-line radio galaxies and radio-loud quasars) would show both broad and narrow optical lines, while edge-on radio-loud objects (narrowline radio galaxies, NLRG) would show only narrow optical lines, and would

¹Harvard-Smithsonian Center for Astrophysics, 60 Garden Street, Cambridge, MA 02138

 $^{^2 \}rm School of Physics, Astronomy & Mathematics, University of Hertfordshire, College Lane, Hatfield, AL10 9AB, UK$

have a clear mid-infrared signature of the absorbing torus. This radio-loud unified scheme well describes the nuclei of many of the most powerful radio sources (e.g., Barthel 1989; Haas 2004), although it is well known that many radio galaxies do not have the strong optical line-emission that is expected from a conventional AGN (Hine & Longair 1979; Jackson & Rawlings 1997). The objects lacking these narrow lines, the low-excitation radio galaxies (LERGs), in general show no evidence in the mid-IR for an obscuring torus, either at low or high luminosities (Whysong & Antonucci 2004; Ogle et al. 2006). Most low radio-power (FRI) sources galaxies are LERGs, whereas most high radio-power (FRII) sources are high-excitation radio galaxies (HERGs – i.e., NLRGs, BLRGs, and quasars). However, **the low-excitation/high-excitation division does not correspond to the FRI/FRII division**: there is a small number of highexcitation FRI sources (including the nearest radio-loud AGN, Centaurus A), as well as a significant population of low-excitation, radio-powerful FRII sources.

2. The Origin of X-ray Emission

2.1. Overview

The physical origin of nuclear X-ray emission in radio-loud AGN has been a topic of considerable debate. In particular, it has been unclear as to whether the emission primarily originates in an quasi-isotropic accretion flow, or instead is associated with an intrinsically beamed parsec-scale radio jet. The detection with *ROSAT* of unabsorbed, power-law X-ray emission in the B2 (Canosa et al. 1999) and 3CRR (Hardcastle & Worrall 1999) samples, together with observed correlations between both the X-ray and VLA radio core fluxes and luminosities led those authors to suggest a nuclear jet-related origin for at least the soft Xray emission. However, the torus required to obscure the optical and UV disk emission should also obscure any X-ray emission associated with the accretion disk. Therefore, AGN inclined at low to intermediate angles with respect to the observer should show a component of heavily absorbed, accretion-related, nuclear X-ray emission, similar to that observed in Seyfert 2 galaxies. This was borne out by early studies of individual objects (e.g., Ueno et al. 1994) as well as more detailed studies of large samples with hard X-ray instruments such as ASCA and BeppoSAX (e.g., Sambruna et al. 1999; Grandi et al. 2006).

Chandra and XMM-Newton have revolutionized the study of radio-galaxy nuclei. Chandra is particularly suited to this task, owing to its high angular resolution and corresponding ability to spatially separate AGN emission from that of the surrounding hot-gas environment. In turn, the excellent sensitivity of XMM-Newton allows us to search for, with greater sensitivity, components of heavily absorbed (and likely accretion-related) X-ray emission. As the orientation-dependent effects of relativistic beaming and the putative obscuring torus are expected to play a large part in determining the observed properties of a radio-galaxy nucleus, it is important to select sources based on their lowfrequency (and hence isotropic) emission characteristics, such as in the 3C and 3CRR samples. In these proceedings, we concentrate on Chandra and XMM-Newton observations of radio galaxies drawn from the 3CRR catalog in the redshift range z < 0.5, as discussed by Donato et al. (2004), Evans et al. (2006), Balmaverde et al. (2006), and Hardcastle et al. (2006). Of the 86 z < 0.5 3CRR radio sources, 40 have been observed with *Chandra* or *XMM-Newton*.

2.2. X-ray Spectra and Correlations

Spectral analysis with *Chandra* and *XMM-Newton* of the z < 0.5 sources shows that *every one* possesses an unabsorbed component of nuclear X-ray emission, as first seen with *ROSAT*. However, narrow-line (i.e., high-excitation) radio galaxies (almost all of which have FRII morphologies, with a handful showing FRI radio features) show an *additional* heavily absorbed nuclear component, with a column density often in excess of 10^{23} cm⁻², accompanied by narrow Fe K α line emission. The properties of these narrow-line FRIIs are thus consistent with the expectation from unified models. However, *no* low-excitation radio galaxy in the 3CRR sample shows any evidence for this type of heavily absorbed nuclear component (Hardcastle et al. 2006).

Figure 1a shows the 1-keV luminosity of the *unabsorbed* power-law component against the 5-GHz luminosity of the radio core. This illustrates the correlation between the soft X-ray emission and the radio core, as first discussed by Hardcastle & Worrall (1999). The correlation implies that the X-ray emission is affected by relativistic beaming in the same manner as the radio, and so suggests a physical relationship between the two, most plausibly at the base of the jet. There is no systematic difference in the behavior of FRIs and FRIIs, or NLRGs and LERGs. The broad-line objects in the sample lie above the trendline established by the NLRG and LERG, and Evans et al. (2006) attribute this to the presence of unabsorbed accretion-related emission in their spectra, consistent with the near face-on orientation of these objects in unified AGN models.

3. The Ubiquity of the Torus and Constraints on the Accretion Mode

We have established that low-excitation radio galaxies possess no evidence for heavily absorbed nuclear X-ray emission. What does this imply for the ubiquity of the torus and the nature of the accretion flow? We assumed that in each LERGs, there exists a 'hidden' component of accretion-related emission that is obscured by a column 10^{23} cm⁻², in addition to jet-related component of X-ray emission that dominates the spectrum. We then determined the 90%-confidence upper limit to the 2–10 keV luminosity of this hidden component. The accretionrelated luminosity of the HERGs is given by the unabsorbed luminosity of the heavily obscured emission in NLRGs, and by the offset (Fig. 1a) in the case of broad-line objects.

Figure 1b shows a plot of the unabsorbed 2–10 keV accretion-related luminosity against 178-MHz radio luminosity, for both the low- and high-excitation radio sources. The upper limits on the accretion-related components in the LERGs, given our assumed absorbing column of 10^{23} cm⁻², lie systematically below the detected HERGs at all radio powers. If no obscuring region is present at all in LERGs, then the luminosity of any accretion-related emission will be substantially lower than that shown.

We now turn to the implications of this result on the nature of the accretion flow in low- and high-excitation sources. One widely discussed model (e.g., Reynolds et al. 1996; Donato et al. 2004) is that there exists a funda-



Figure 1. (a) X-ray luminosity of the unabsorbed X-ray component as a function of 5-GHz radio core luminosity. Open circles are LERG, filled circles NLRG, open stars BLRG, and filled stars quasars. Surrounding circles mean a source is an FRI. Dotted lines show 90% confidence regression lines fitted through the NLRGs and LERGs. BLRGs and quasars lie above the line. (b) X-ray luminosity of the accretion-related component for the combined z < 0.5 sample as a function of 178-MHz total radio luminosity. The regression line is determined by the NLRG.

mentally different accretion *mode* in FRI- and FRII-type sources, such that the accretion-flow luminosities and radiative efficiencies of FRI-type radio galaxies are systematically lower than those of FRII-type radio galaxies. The main difficulty from the point of view of radio-galaxy physics with such a model is that the FRI/FRII dichotomy, and its dependence on host galaxy properties (Ledlow & Owen 1996), can be explained purely in terms of jet power and the interaction with the environment (e.g., Bicknell 1995). However, we can readily modify the accretion model to accommodate a scheme in which **radiatively inefficient accretion flows power LERGs, and efficient accretion via a standard thin disk powers HERGs.**

4. Powering AGN: Hot vs. Cold Gas Accretion

4.1. Bondi Accretion

We have argued that LERGs may be a class of luminous active galaxies that accrete radiatively inefficiently, with almost all the available energy from accretion being channeled into the jets, while high-excitation sources are powered by radiatively efficient accretion. In Hardcastle et al. (2007) we considered whether these different accretion modes may be a result of a different *source* for the accreting gas, building on the recent result of Allen et al. (2006), who showed that some low-luminosity radio galaxies in the centers of clusters could be powered by Bondi accretion from the hot, X-ray emitting medium. We can use X-ray observations of the hot-gas environment of radio sources to estimate the Bondi accretion rate, and in turn constrain the amount of power that the central supermassive black hole can extract from accretion of the hot phase of the IGM. The Bondi rate is given by $\dot{M} = \pi \rho_A G^2 M_{\rm BH}^2 / c_s^3$, where r_A is the Bondi accretion radius and $c_{\rm s}$ is the sound speed in the medium. The available power for AGN activity, $P_{\rm B} = \eta \dot{M} c^2$, where η is an efficiency factor (assumed to be 0.1). To derive black-hole masses for radio galaxies we use the relationship between $M_{\rm BH}$ and K-band absolute bulge magnitude derived by Marconi & Hunt (2003) for nearby sources. To show that the jet can be powered by accretion of the IGM, we require that the (kinetic plus radiative) jet power $Q \leq P_{\rm B}$. In Hardcastle et al. (2007), we presented a detailed description of methods used to estimate the jet power. In short, we used the Willott et al. (1999) relation between jet power and 151-MHz luminosity density: $Q_W =$ $3 \times 10^{38} f^{3/2} L_{151}^{6/7}$ W, where f parametrizes our ignorance of true jet powers. In Figure 2 we plot the observational quantities, radio luminosity and K-

In Figure 2 we plot the observational quantities, radio luminosity and Kband luminosity, together with their conversion to Bondi power and jet power. Figure 2 shows that the nearby FRI radio galaxies almost all lie within a factor of a few of the line of $Q_W = P_B$. Secondly, it shows that the majority of lowexcitation FRII radio galaxies in our sample also lie close to this line. And thirdly, it shows that there is a population of FRII sources, encompassing most of the narrow-line FRII sources (and therefore, presumably, all high-excitation sources), that have jet powers exceeding the available Bondi powers (for our choice of central gas properties), often by more than two orders of magnitude.

In summary, we have shown that it is possible that all the low-excitation radio galaxies are powered by accretion from the hot phase, consistent with the fact that their nuclear spectra show no evidence for cold material close to the nucleus (i.e., no evidence for the 'torus'). On the other hand, we have seen that narrow-line radio galaxies (and therefore other high-excitation sources) have clear evidence for accretion disks and tori, cannot be powered in this way the large amounts of cold material in the nucleus and the radiative efficiency of accretion are naturally explained if these objects are powered by accretion of cold material via a thin disk in the standard manner.

4.2. Implications

Hardcastle et al. (2007) present a detailed discussion of the implications of hot vs. cold gas accretion for radio sources. We summarize these below:

Feedback: AGN feedback is thought to be an important aspect of galaxy formation models (e.g., Croton et al. 2006; Bower et al. 2006). An important feature of these models is that the AGN should both be able to influence, *and should be influenced by*, the X-ray emitting phase. Direct accretion of the hot phase provides an natural way of ensuring that the AGN activity is regulated by the gas properties at the cluster center. However, this is only possible for a 'hot-mode' radio source. Cold-mode sources do not have this direct connection between the hot phase and the rate of fueling of the AGN: instead, the jet power is controlled solely by the accretion rate of cold gas, and so these sources can potentially inject a significant amount of energy input to the IGM without regulation from their hot-gas environments.



Figure 2. K-band host galaxy luminosity against 151-MHz luminosity for LERGs and HERGs, with conversions into Willott jet power (Q_W) and available Bondi power (P_B) . Open stars are LERGs and filled stars are NLRGs. A circle round a data point indicates an FRI. The central solid line shows equality between the predicted Bondi power and the Willott jet power, and the dashed lines are separated from the solid line by one order of magnitude.

Environments: In the model we have outlined we expect different types of active galaxies to be found in different environments. Cold-mode accretion requires a supply of cold gas: the easiest way for an elliptical galaxy to acquire this is by a merger with a gas-rich system. Samples of high-excitation radio galaxies should thus show evidence for mergers and interactions, as is often observed (e.g., Heckman et al. 1986). Host galaxies of cold-mode systems do not need a rich environment, or to be at the bottom of a deep potential well, so long as galaxy-galaxy mergers can take place. By contrast, hot-mode accretion requires a supply of hot gas and a massive central black hole. Both the black hole mass and the mass of the galaxy-scale X-ray halo (e.g., Mathews & Brighenti 2003) are correlated with the mass of the host galaxy. Thus we expect hot-mode systems – which, observationally, include almost all FRI radio galaxies – to favor massive galaxies, and the most powerful radio sources to tend to be group- or cluster-dominant systems (e.g., Longair & Seldner 1979; Best 2004).

An overview of the properties of LERGs and HERGs is given in Table 4.2.

Acknowledgments. D.A.E. gratefully acknowledges support from NASA through *XMM-Newton* award NNX06AG37G.

	Low-excitation (LERG)	High-excitation (HERG)
Definition	No narrow optical line emis-	Prominent optical emission
	sion.	lines, either narrow (NLRG)
		or broad (BLRG), or quasar.
Fanaroff-Riley	Almost all FRIs are LERGs,	Most FRIIs are HERGs, as
classification	as well as a significant pop-	are a handful of FRIs (e.g.,
	ulation of FRIIs.	Cen A).
X-ray spectra	Jet-related unabsorbed	Jet-related unabsorbed
	power law only. Upper	power law $+$ significant
	limits only to "hidden"	accretion contribution
	accretion-related emission.	(heavily absorbed in NL-
		RGs).
Accretion-flow	Highly sub-Eddington.	Reasonable fraction of Ed-
type	Likely radiatively ineffi-	dington. Likely standard
	cient.	accretion disk.
Optical con-	Strong radio/optical/soft X-	Strong radio/optical/soft X-
straints	ray correlations. Optical	ray correlations. Optical
	emission is jet-related.	emission is jet-related.
Fueling mecha-	Bondi accretion of hot ISM.	Additional fuel supply
nism		needed, likely cold gas.
Implications of	Significant feedback be-	Potentially large energy in-
AGN fueling	tween AGN and environ-	put to IGM, decoupled from
	ment.	hot-gas environment.

Table 1. Overview of the properties of low- and high-excitation radio galaxies

References

- Allen, S. W., Dunn, R. J. H., Fabian, A. C., Taylor, G. B., Reynolds, C. S., 2006, MNRAS, 372, 21
- Balmaverde, B., Capetti, A., Grandi, P., 2006, A&A, 451, 35
- Barthel, P. D., 1989, ApJ, 336, 606
- Best, P. N., 2004, MNRAS, 351, 70
- Bicknell, G. V., 1995, ApJS, 101, 29
- Bower, R. G., et al., 2006, MNRAS, 370, 645
- Blundell, K. M., Rawlings, S., 2000, AJ, 119, 1111
- Canosa, C. M., Worrall, D. M., Hardcastle, M. J., & Birkinshaw, M., 1999, MNRAS, 310, 30
- Croton, D., et al., 2006, MNRAS, 365, 111
- Donato, D., Sambruna, R. M., & Gliozzi, M., 2004, ApJ, 617, 915
- Evans, D. A., Worrall, D. M., Hardcastle, M. J., Kraft, R. P., Birkinshaw, M., 2006, ApJ, 642, 96
- Grandi, P., Malaguti, G., Fiocchi, M., 2006, ApJ, 642, 113
- Haas, M., et al., 2004, A&A, 424, 531
- Hardcastle, M. J., & Worrall, D. M., 1999, MNRAS, 309, 969
- Hardcastle, M. J., Evans, D. A., Croston, J. H., 2006, MNRAS, 370, 1893
- Hardcastle, M. J., Evans, D. A., Croston, J. H., 2007, MNRAS, 376, 1849
- Heckman, T. M., et al., 1986, ApJ, 311, 526
- Hine, R. G., Longair, M. S., 1979, MNRAS, 188, 111
- Jackson, N., Rawlings, S., 1997, MNRAS, 286, 241
- Ledlow, M. J., & Owen, F. N., 1996, AJ, 112, 9

- Longair, M. S., Seldner, M., 1979, MNRAS, 189, 433
- Marconi, A., Hunt, L. K., 2003, ApJ, 589, L21
- Mathews, W. G., Brighenti, F., 2003, ARA&A, 41, 191
- Ogle, P., Whysong, D., Antonucci, R., 2006, ApJ, 647, 161
- Reynolds, C. S., di Matteo, T., Fabian, A. C., Hwang, U., & Canizares, C. R., 1996, MNRAS, 283, L111
- Sambruna, R. M., Eracleous, M., Mushotzky, R. F., 1999, ApJ, 526, 60
- Ueno, S., Koyama, K., Nishida, M., Yamauchi, S., Ward, M. J., 1994, ApJ, 431, L1
- Whysong, D., Antonucci, R., 2004, ApJ, 602, 116
- Willott, C. J., Rawlings, S., Blundell, K. M., Lacy, M., 1999, MNRAS, 309, 1017
- Willott, C. J., Rawlings, S., Jarvis, M. J., Blundell, K., 2003, MNRAS, 339, 173

