Uncovering extraplanar gas in UGCA 250 with the Ultra-deep MHONGOOSE Survey

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Accepted 2025 March 5. Received 2025 March 2; in original form 2024 November 18

ABSTRACT

We use the neutral atomic hydrogen (H I) observations of the edge-on galaxy UGCA 250, taken as part of the MeerKAT H I Observations of Nearby Galactic Objects-Observing Southern Emitters (MHONGOOSE) survey to investigate the amount, morphology, and kinematics of extraplanar gas. The combination of high column density sensitivity and high spatial resolution of the survey over a large field of view is ideal for studying the underlying physics governing the extraplanar gas. These data reveal nine additional detections within the field of view along with UGCA 250, with eight of them being within ~200 km s⁻¹ of the galaxy's systemic velocity. The galaxy seems to have a tail-like feature extending away from it in the southern direction up to ~41 kpc (in projection). We also detect a cloud at anomalous velocities, but we did not find any optical counterpart. We construct a detailed tilted ring model for this edge-on galaxy to gain a deeper understanding of the vertical structure of its neutral hydrogen. The model that best matches the data features a thick disc with a scale height of ~3±1 kpc and an H I mass of about 15 per cent of the total H I mass. This extraplanar gas is detected for the first time in UGCA 250. Our analysis favours a mixed origin for the extraplanar gas in UGCA 250, likely arising from a combination of internal stellar feedback and external tidal interactions.

Key words: galaxies: haloes – galaxies: individual – galaxies: kinematics and dynamics – galaxies: structure.

1 INTRODUCTION

It is becoming increasingly clear that gas accretion is necessary to sustain the observed levels of star formation rates in disc galaxies. Gas accretion on to galaxies can occur via different mechanisms, such as wet mergers with gas-rich satellites (Rubin et al. 2012) or gradual cooling of gas accreted on to the galactic halo from the cosmic web (van de Voort et al. 2011; Wetzel & Nagai 2015; Cornuault et al. 2018; Ho, Martin & Turner 2019; Iza et al. 2022), possibly triggered by hydrodynamical interactions with the galactic fountain (e.g. Fraternali et al. 2002; Melioli et al. 2008, 2009; Marinacci et al. 2010; Fraternali 2017). The halo of a galaxy will form its connection with the intergalactic medium and one must focus on the disc–halo interface region to understand the gas exchange processes between the disc and the circumgalactic medium (CGM).

Deep H I observations of several nearby spiral galaxies have shown the presence of cold extraplanar gas (EPG) in a gaseous layer that surrounds the star forming disc up to a height of some kpc (e.g. Sancisi et al. 2008; Hess et al. 2009; Lucero et al. 2015). This layer typically has a mass of 10–30 per cent of the mass of the HI in the disc (e.g. Oosterloo, Fraternali & Sancisi 2007; Marasco et al. 2019). Observations of nearby spiral galaxies with different inclination angles provide information on the structure and kinematics of neutral EPG in disc galaxies (see Boomsma et al. 2008 for face-on galaxies; Fraternali et al. 2001; Hess et al. 2009; Marasco et al. 2019 for intermediate inclination galaxies; and Oosterloo et al. 2007; Rand & Benjamin 2008; Zschaechner et al. 2011 for edge-on galaxies).

For example, the Westerbork Synthesis Radio Telescope observations of the edge-on galaxy NGC 891 revealed a slowly rotating (also known as lagging) HI halo, which extends vertically out to 14 kpc from the thin HI disc and, in one quadrant up to 22 kpc, with filamentary structures far from the disc (Swaters, Sancisi & van der Hulst 1997; Oosterloo et al. 2007). In the face-on galaxy NGC 6946, vertical motions are observed across the star forming disc and are in clear connection with EPG (Boomsma et al. 2008). In some galaxies, non-circular motions have also been detected. The EPG of NGC 2403, UGCA 105, NGC 4559 have a coherent radial infall motion towards the centre of the galaxy (Fraternali et al. 2002;

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α, δ (J2000)	11 ^h 53 ^m 24.1 ^s , -28°33 [′] 11.4 ^{″′}				
$V_{\rm sys} (\rm km s^{-1})$	1705 ^a				
PA (°)	256 ^a				
<i>i</i> (°)	82 ^a				
d (Mpc)	20.4 ^b				
$\log(M_*/M_{\odot})$	9.80 ^c				
SFR (M _☉ /yr)	0.54 ^c				
D ₂₅ (kpc)	21.6^{c}				
$D_{\rm HI}({\rm kpc})$	69.0 ^a				

^{*a*}From this work (at H I column density $\sim 1.2 \times 10^{19}$ cm⁻²). ^{*b*}Distance is taken from Kourkchi et al. (2022).

^cFrom Leroy et al. (2019); α, δ (J2000): Coordinates; V_{sys} : Heliocentric

systemic velocity; PA: Position angle; *i*: Inclination angle; *d* : Distance; $\log(M_*/M_{\odot})$: Stellar mass; SFR: Star formation rate; D_{25} : Optical diameter; D_{H1} : HI diameter.

Barbieri et al. 2005; Schmidt et al. 2014; Vargas et al. 2017). Trends of EPG, such as the presence and magnitude of any decrease in rotation speed with height are useful in tracing their origin.

Various mechanisms have been proposed to explain the origin of EPG. One is an internal mechanism in which a large part of this halo gas is brought up from the disc by galactic fountains (e.g. Shapiro & Field 1976; Bregman 1980). In this model, supernova explosions expel the hot gas to large distances from the disc into the halo. In the halo the gas cools down and condenses before returning to the disc. The other is an external mechanism in which a galaxy accretes gas from the CGM (Binney 2005; Kaufmann et al. 2006). Additionally, a 'mixed' scenario has been proposed where the EPG is formed by the interplay between the galactic fountain and the hot CGM. The interactions between these two components produce a mixture with low cooling time-scales and slow rotational speed, which explains the observed EPG kinematics and provides a potential source of pristine gas accretion on to the galaxy (Fraternali & Binney 2008; Marasco, Fraternali & Binney 2012; Li et al. 2023). A powerful method to determine the origin of EPG is to study their kinematics. In particular, edge on galaxies provide a great opportunity to measure the vertical distribution of the interstellar medium.

Here, we use the 'full depth' HI 21 cm observations of the nearly edge-on galaxy UGCA 250 (J1153-28), taken as part of the MeerKAT HI Observations of Nearby Galactic Objects-Observing Southern Emitters (MHONGOOSE) survey (de Blok et al. 2024) to study the detailed spatial and kinematical distribution of the gas, with a focus on the faint outer edges. UGCA 250 is a nearby galaxy at a distance of 20.4 Mpc (Kourkchi et al. 2022). The galaxy has an optical diameter (d_{25}) of 21.6 kpc, stellar mass $\log(M_*/M_{\odot}) =$ 9.80, and star formation rate of $0.54 \,\mathrm{M_{\odot} \, yr^{-1}}$, which places it on the star formation main sequence for spiral galaxies. Table 1 lists the basic parameters of UGCA 250. The MHONGOOSE survey is an ultra-deep HI survey that uses the MeerKAT radio interferometer to observe the HI in 30 nearby galaxies. The combination of high column density sensitivity and high spatial resolution over a large field of view is ideal for studying the underlying physics governing the EPG as well as the disc-halo connection.

This paper is organized as follows. In Section 2, we describe the observations and data reduction techniques. In Section 3, we presents the results, including the detection of companion galaxies, H I morphology of UGCA 250, and tilted ring models. In Section 4, we discuss the various models used to understand the origin of the EPG. Finally, in Section 5, we summarize our key findings.

2 OBSERVATIONS AND DATA REDUCTION

The MeerKAT observations of UGCA 250 were taken as part of the MHONGOOSE survey in 10 tracks each lasting 5.5 h, resulting in a total on source time of 50 h. Each track consists of 10 min of observing time on one of the primary calibrators, J0408–6545 or J1939–6342, and the secondary calibrator, J1154–3505, was observed every 50 min for 2 min. We used the narrow-band mode of MeerKAT which has 32 768 channels of 3.265 kHz each, giving a total bandwidth of 107 MHz. This corresponds to a rest frame velocity resolution of 0.7 km s⁻¹ at 1420 MHz. All data were subsequently binned to 1.4 km s⁻¹ for the reduction process.

We use the Containerised Automated Radio Astronomy Calibration (CARACAL)¹ pipeline (Józsa et al. 2020) to carry out the standard calibration and data reduction processes. CARACAL is based on STIMELA, a radio interferometry scripting framework based on container technologies and PYTHON (Makhathini 2018), which uses many publicly available radio interferometry software packages such as CASA, CUBICAL, and WSCLEAN, etc. The pipeline flags the calibrators for radio frequency interference, derives and applies the cross-calibration solutions, and splits off the target. Further, we use the self-calibration of the continuum to improve the quality of the calibration. After applying the self-calibration, we subtract the continuum from the measurement set.

The HI cubes were then imaged at various resolutions using WSCLEAN as part of the CARACAL. We make six HI cubes that range in resolution from 6 to 90 arcsec using different combinations of robust weighting (*r*) and Gaussian tapering (*t*), leading to varying HI column density sensitivity ranging from $N_{\rm HI} \sim 6.0 \times 10^{19} \,\rm cm^{-2}$ to $5.7 \times 10^{17} \,\rm cm^{-2}$ at 3σ over $16 \,\rm km \, s^{-1}$. For most of the analysis, we use r15t00 configuration, where r = 1.5 and t = 0. This yields a resolution of 34.3 arcsec $\times 25.6 \,\rm arcsec \, (\sim 3 \,\rm kpc$ at the distance of UGCA 250), and achieves a 3σ column density sensitivity of approximately $2.7 \times 10^{18} \,\rm cm^{-2}$ over $16 \,\rm km \, s^{-1}$, which offers an optimal balance between column density sensitivity and spatial resolution. A more detailed description of the data reduction and data quality assessment is presented in Blok et al. (2024).

We use SoFiA-2 (Westmeier et al. 2021) to search for sources and to generate the moment maps of the detections, which are standard maps produced for all the MHONGOOSE observations. We smooth the cube spatially and spectrally with the 'smooth and clip' algorithm of SoFiA-2. We apply a threshold of 4σ using two spatial Gaussian smoothing kernels: no smoothing, one times the beam; in addition to three spectral boxcar smoothing kernels: no smoothing, 12 km s^{-1} and 35 km s^{-1} . We then create moment maps of detections from the masked data cubes and we correct the moment-0 maps for the primary beam attenuation of MeerKAT (see Blok et al. 2024; Healy et al. 2024, for more details).

In this paper, we use deep optical observations made using OMEGACAM on the VLT Survey Telescope (VST) for UGCA 250 from dedicated proprietary observations (ID: 110.25AJ, PI: Maccagni; ID: 112.266Y, PI: Marasco) of the MHONGOOSE galaxies (hereafter MHONGOOSE–VST). These observations reach a limiting surface brightness of $\mu_g = 28.5 \text{ mag/arcsec}^2$, $\mu_r = 27.3 \text{ mag/arcsec}^2$ in the *g* and *r* bands respectively, for a point source at 5σ over a circular area with FWHM ~ 1 arcsec. To create the RGB image shown in Fig. 1, we combine the *g* and *r* bands from MHONGOOSE–VST with publicly available *i*-band observations from the VST Early-type Galaxy Survey (VEGAS; Spavone et al. 2017; Iodice

¹https://github.com/caracal-pipeline/caracal



Figure 1. Composite optical image UGCA 250. The RGB colour is provided by a combination of g, r, and i-band VST data. At brighter magnitudes, the galaxy exhibits a thin, elongated structure typical of an edge-on galaxy. At fainter magnitudes, it shows a vertical extension and appears asymmetric, with an extension towards the southwest.

et al. 2021). Although VEGAS primarily targets elliptical galaxies, UGCA 250 was observed because it lies approximately 30 arcmin away from the elliptical galaxy NGC 3923 (Kourkchi & Tully 2017). At brighter magnitudes, the optical morphology resembles an edgeon galaxy, characterized by a thin, elongated structure. However, at fainter magnitudes, it exhibits a vertical extension and the galaxy appears asymmetric with a with a noticeable discontinuity towards the southwestern side. This discontinuity is more likely a feature connected to the main disc rather than a separate dwarf galaxy. As will be discussed in Section 3.1, HI observations do not reveal emission at a distinct velocity at that location, a signature expected if a separate dwarf were present. Since the MHONGOOSE-VST observations do not fully cover the surrounding dwarf galaxies, we instead use the VEGAS g-band image in Fig. 2, which provides a more complete view of the environment around UGCA 250, including all nearby dwarfs.

3 RESULTS

The source finding yields the HI emission associated with a total of 10 galaxies (including UGCA 250), within the field of view of $\sim 1.5^{\circ}$ which corresponds to a projected radius of ~535 kpc at the distance of the galaxy. Within this field of view also lies the elliptical galaxy NGC 3923, positioned 30 arcmin (or approximately 200 kpc in projection) from UGCA 250 and sharing a similar redshift ($V_{sys} \sim 1739 \text{ km s}^{-1}$). Despite its proximity, no H I emission was detected in NGC 3923. The 5-sigma H I mass upper limit for this non-detection is $10^{6.4}$ M_{\odot}, after correcting for the primary beam attenuation, assuming a velocity width of $250 \,\mathrm{km \, s^{-1}}$ and the same distance as that of UGC 250. The HI properties of the detected galaxies are given in Table 2. All the detected galaxies, except 6dFJ1151533-281047 are within 200 km s⁻¹ of the systemic velocity of UGCA 250. The left panel of Fig. 2 shows the HI contours of UGCA 250 and its satellite galaxies overlaid on the VEGAS g-band optical image (The moment 0 and moment 1 maps of individual satellite galaxies are shown in Figs A3

and A4). The satellite galaxy IDs that are labelled are based on Table 2.

The H I contours are displayed at different resolutions: 26.7 arcsec \times 18.8 arcsec (r10t00) in purple, 34.3 arcsec \times 25.6 arcsec (r15t00) in cyan, 65.3 arcsec \times 63.8 arcsec (r05t60) in magenta, and 93.8 arcsec \times 91.7 arcsec (r10t90) in blue. The lower resolution contours (magenta and blue) reveal both intra-group gas and gas around specific satellites such as ESO440–G023 (ID 1), 2MFGC09308 (ID 2), LEDA737259 (ID 3), and MKT115222-280126 (ID 8), and 6dFJ1151533-281047 (ID 10).

We calculate the integrated flux of each detection using the SoFiA-2 output. Further, we calculate their HI masses using the standard formula, $M_{\rm HI} = 2.356 \times 10^5 D^2 S_{\rm int} M_{\odot}$, where D is the distance in Mpc, S_{int} is the integrated flux density in Jy km s⁻¹ (Roberts 1962). Six of the nine satellite galaxies (marked with * in Table 2) were detected in the 5.5-h single-track observations (Blok et al. 2024). Full-depth observations revealed three additional galaxies: 2MFGC 09308, MKT115222-280126, and MKT115551-282149, with H I masses $\log(M_{\rm H I}/\rm M_{\odot}) \lesssim 6.60$, consistent with the expected sensitivity limits of single-track and full-depth observations. MKT115551-282149, with $\log(M_{\rm H\,I}/M_{\odot}) \lesssim 6.10$, was detected at a distance of 20 Mpc, demonstrating the remarkable sensitivity of the full-depth MHONGOOSE observations. Similarly, a low-mass, gas-rich, low-surface brightness galaxy in the Dorado Group was detected at 17.7 Mpc (Maccagni et al. 2024). These results highlight the enhanced capabilities of full-depth observations in detecting lowmass galaxies at larger distances.

3.1 HI morphology of UGCA 250

The right panel of Fig. 2 displays a zoomed-in view of UGCA 250 at a resolution of 26.7 arcsec \times 18.8 arcsec (r10t00), which clearly shows the extent of gas at different column densities. Additionally, we show low-resolution contours in cyan, magenta, and blue, corresponding to 34.3 arcsec \times 25.6 arcsec, 65.3 arcsec \times 63.8 arcsec, and 93.8 arcsec \times 91.7 arcsec respectively, in the outer regions. The red contour corresponds to a HI column density of 1.0 \times 10^{20} cm⁻². Even at this column density, a noticeable position angle warp in the West direction is apparent. As we transition to lower HI column densities, the disc does not expand significantly in the radial direction, but it does grow substantially in the vertical direction. The warp becomes more prominent at these lower column densities, and the HIdistribution of UGCA 250 appears asymmetric, and slightly lopsided, with the disc column densities dropping more rapidly towards the northeast and southeast compared to the northwest and southwest. The thickening observed in the HI disc is also apparent in the optical stellar disc (see Fig. 1).

The galaxy has a diameter of ~69 kpc at the projected H I column density sensitivity of ~ 3.3×10^{18} cm⁻². The prominent feature of this gas disc is that there is a tail like feature (henceforth filament) extending away from it in the southern direction up to ~34 kpc (in projection). The blob below this is at a projected distance of ~41.5 kpc from the centre of galaxy, and has a H I mass, log($M_{\rm HI}$) ~ 6.34. The HImass of filament along with the blob is, log($M_{\rm HI}$) ~ 7.01.

The moment maps derived at a slightly coarser resolution of $34.3 \operatorname{arcsec} \times 25.6 \operatorname{arcsec} (r15t00, r = 1.5 \operatorname{and} t = 0)$ are shown in Fig. 3. The left, middle, and right panels show the moment 0, moment 1, and moment 2 maps, respectively. The middle panel shows the intensity weighted velocity field, which is rather complex. The inner disc is regularly rotating and the iso-velocity contours seem to turn over at lower H I column densities. There is a cloud-like



Figure 2. The H1integrated intensity contours are overlaid on the VST optical *g*-band image. (a) Moment 0 map of UGCA 250 and its nine satellite galaxies. The contours are shown from data at various resolutions: 26.7 arcsec × 18.8 arcsec (rlot00) in purple with levels of 1.2×10^{19} cm⁻² × 2ⁿ, n = 0,1,2,3..; 34.3 arcsec × 25.6 arcsec (rlot00) in cyan at 3.0×10^{18} cm⁻²; 65.3 arcsec × 63.8 arcsec (rlot00) in magenta at 1.0×10^{18} cm⁻²; and 93.8 arcsec × 91.7 arcsec (rlot90) in blue at 6.2×10^{17} cm⁻². (b) Moment-0 map of the galaxy UGCA 250 with the same H I contours, where the thick red contour corresponds to 1.0×10^{20} cm⁻². Note that the H I column density values have not been corrected for the primary beam attenuation. The ellipses in the bottom left corner represent the beams at each resolution. The zoomed-in moment 0 and moment 1 maps of individual satellite galaxies are shown in Figs A3 and A4 in the Appendix. (c) The red, orange, and yellow show the H I integrated intensity from three different resolution data cubes: $14.4 \operatorname{arcsec} \times 9.8 \operatorname{arcsec} \times 26.7 \operatorname{arcsec} \times 18.8 \operatorname{arcsec}$, and 65.3 arcsec × 63.8 arcsec. These H I maps are overlaid on a composite optical RGB image of UGCA 250 in the background. The creation of this image followed the technique outlined in English (2017).

Table 2.	H I properties	of satellite	galaxies.
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ID	Galaxy ID	RA (h m s)	Dec (°''')	$V_{\rm sys}$ (km s ⁻¹)	W_{20} (km s ⁻¹)	W_{50} (km s ⁻¹)	Flux (Jy km s ⁻¹)	H I mass $log(M_{\rm HI}/M_{\odot})$
1	ESO 440-G 023*	11 52 33.1	-29 07 19.1	1866	114.4	92.3	0.31	7.49
2	2MFGC 09308	11 51 39.3	$-28\ 45\ 17.5$	1887	71.6	46.8	0.04	6.60
3	LEDA 737259*	11 54 14.8	$-28\ 48\ 19.8$	1761	42.7	27.5	0.62	7.79
4	ESO 440-G 014*	11 50 03.9	$-28\ 40\ 18.2$	1857	68.9	44.1	0.10	6.98
5	LEDA 746393*	11 53 10.4	$-28\ 03\ 30.4$	1828	95.1	71.6	0.09	6.96
6	UGCA 250*	11 53 24.3	$-28\ 33\ 08.9$	1694	268.7	285.3	72.0	9.85
7	ESO 440-G 030*	11 55 25.9	$-28\ 44\ 10.2$	1801	115.8	91.0	0.78	7.89
8	MKT115222-280126	11 52 22.1	$-28\ 01\ 26.2$	1636	42.7	28.9	0.02	6.30
9	MKT115551-282149	11 55 51.4	-28 21 49.3	1595	68.9	45.5	0.01	6.10
10	6dFJ1151533-281047*	11 51 53.3	$-28\ 10\ 50.2$	1422	78.5	50.9	0.08	6.89
	Cloud at anomalous velocities (C)	11 53 20.5	$-28\ 35\ 47.0$	1605	-	-	0.065	6.81
	Blob below filament (B)	11 53 38.4	$-28\ 39\ 43.0$	1640	_	-	0.02	6.34
	Entire filament (with blob)	11 53 31.3	-28 37 34.0	1617	-	-	0.10	7.01

*Galaxies that were previously detected in 'single track' (5.5 h) shallow data (Blok et al. 2024); ID: Identifier; Optical ID: Optical identification; RA: Right ascension; Dec: Declination; V_{sys} : Systemic velocity in km s⁻¹; W_{20} and W_{50} : H I line width at 20 per cent peak and line width at 50 per cent peak in km s⁻¹; Flux: Integrated flux in Jy km s⁻¹; H I mass: H I mass in solar masses (M_{\odot}).

feature in the velocity field in the south of the galaxy which is at anomalous velocities. The visual inspection of the cube has revealed a cloud of gas coinciding with the main emission of the galaxy, but at a slightly different velocity, i.e at ~1605 km s⁻¹ in comparison to the systemic velocity ~1694 km s⁻¹. This can be seen in the channel maps presented in Fig. 4 in the channels with velocities from 1587 km s⁻¹ to 1621 km s⁻¹.

The right panel shows the second moment map, indicative of velocity dispersion for symmetrical, single Gaussian profiles. We can see two very high velocity dispersion regions with values higher than 40 km s^{-1} . The high values near the cloud in the southern direction are not caused by intrinsically broader HIprofiles, but rather due to the cloud being at a different velocity compared to the galaxy's rotation. We separate this cloud and show its moment maps in Fig. 5. The velocity field of the cloud shows a possible hint of rotation, but the velocity width is less than 15 km s⁻¹. Despite its H I mass, $\log(M_{HI}/$

 M_{\odot}) ~ 6.80, we did not find any optical counterpart for this cloud. Using the cloud size of 57 arcsec × 48 arcsec and assuming a stellar mass-to-light ratio of 1, we estimate a 5 σ upper limit for the stellar mass to be log(M_*/M_{\odot}) ~ 6.1. There is also high velocity dispersion in the northeast of UGCA 250. In the centre, an x-shaped region with high dispersion values is observed (see right panel of Fig. 3), likely due to beam smearing associated with the rise of rotation curve and the inclination of this nearly edge-on galaxy (e.g. Voigtländer et al. 2013).

The features such as asymmetries, filament, cloud, and high velocity dispersion are also noticeable in the position-velocity (PV) diagrams along both the major and minor axes, shown in the Appendix (Figs A1 and A2). Hereafter, we refer to the major axis pv diagram as the lv diagram and the minor axis pv diagram as the bv diagrams, taken at 30 arcsec intervals along the minor axis, highlight features such as the cloud and the



Figure 3. Moment maps of UGCA 250 at a resolution of 34.3 arcsec \times 25.6 arcsec. Left panel: the H I intensity map (moment 0). Middle panel: velocity field (moment 1), The isovelocity contours are spaced at 15 km s⁻¹ and the systemic velocity is indicated by a red line. The colour scheme used in this panel is developed by English et al. (2024). Right panel: observed line width/dispersion map of the galaxy, UGCA 250.

filament extending from the disc. The lv diagrams along the major axis indicate that the galaxy is more extended on the receding half, emphasizing an asymmetry that becomes more pronounced at higher vertical distances above the galactic plane.

3.2 HI kinematics

To investigate the morphology and kinematics of the HI in this galaxy, we construct model data cubes and compare them to the observations. We use the Tilted Ring Fitting Code (TIRIFIC) software (Józsa et al. 2007) for modelling. This software uses a tilted ring approach, representing a galaxy as concentric rings each characterized by specific kinematic and orientation parameters (Rogstad, Lockhart & Wright 1974). These parameters, including central position, systemic velocity, inclination angle with respect to line of sight, position angle of the receding major axis, rotational velocity, velocity dispersion, vertical thickness, and surface brightness distribution can either be fixed or varied with radius. All the modelling described below uses the cubes at resolution of 34.3 arcsec \times 25.6 arcsec (r15t00, r = 1.5 and t = 0). The approaching and receding halves of UGCA 250 are modelled separately since the galaxy clearly shows asymmetries. We first use Fully Automated TIRIFIC (FAT) software (Kamphuis et al. 2015), which is a wrapper around TIRIFIC, to derive the initial guess model, i.e. the initial de-projected HI surface brightness profile and the de-projected rotation curve of the thin disc. Then, we use these parameters as an initial guess for the data cube using TIRIFIC software. An initial estimate for the systemic velocity, centre, position angle, and inclination angle were obtained from the FAT run. These matched well with the rough estimates of kinematic centre and systemic velocity found using the zeroth and first moment maps. The radial variation of position angle and inclination angle were then derived while the galaxy centre and systemic velocity were fixed at the values found above. The position angle (PA) seems to show some radial variation; we find that the value of PA changes only in the outer regions of the receding half and stays constant in the approaching half, as can also be seen in the channel maps of the galaxy in Fig. 4 in the velocities ranging from 1805 to 1838 km s⁻¹. The radial variation of PA was adjusted based on channel maps and the total HI map. We show the PA and inclination angle as a function of radius in the middle and bottom panels of Fig. 6. The radial variation of inclinations suggests the HI disc to be warped along the line of sight, as seen in the moment 0 map. All of these parameters were later refined both by comparing observed and modelled PV diagrams, channel maps, and zeroth-moment maps. The rotation curve was adjusted based on the major-axis PV diagrams tracing emission along the edge with the highest velocities. Fig. 7 shows a PV diagram along the major axis with the derived rotation curve (projected) overlaid on it. The parameters, systemic velocity (1705 km s^{-1}), kinematic centre ($11^{\text{h}} 53^{\text{m}} 24.4^{\text{s}}$, $-28^{\circ} 33' 9.5''$) and the run of the position angle and inclination angle, column density profile, and rotation curve are kept constant throughout all of the models for the entire process. This model, which includes the PA and inclination angle warp component fits the thin disc well, is considered the base model for the galaxy. In the following section, we will include additional components to our model in order to reproduce the H I emission from the EPG.

3.3 Models

In order to understand if the anomalous gas has truly anomalous kinematics or if it is produced by projection effects due to inclination angle or thickness of the disc, etc., we develop several models. We start with the simplest possible model and then add various features (e.g. a warp or flare) one at a time to understand the subtle distinctions between various features. We visually compare the resulting model with data, i.e. we examine the by PV diagrams (i.e. along the minor axis) as well as channel maps of models and data to fine tune parameters in models. We consider several models, including a simple one-component which is a single disc with a vertical sech² profile, a flare where the scale height varies as a function of radius, and two components with a thin disc and a thick corotating disc. Finally, we consider a more sophisticated model of a thin disc and a thick disc with different kinematics and we explore the effect of radial motions (i.e. inflows or outflows). All models include a PA and inclination angle warp component. We have verified that all models that do not include a warp fail at reproducing several of the basic features observed in the data, and will not be considered here.

The selected channel maps shown in Fig. 4 and the bv diagrams shown in Fig. A1 display a number of distinct features sensitive to various model parameters, which were used to distinguish between different models. In the bv diagrams, we can clearly see the anomalous cloud and the filament features of the galaxy. This cloud



Figure 4. Selected channel maps of UGCA 250 at the resolution of 34.3 arcsec \times 25.6 arcsec (r15t00, r = 1.5 and t = 0). The heliocentric velocity in km s⁻¹ is indicated on top-left of each panel. Contours begin at 3σ and increase by factors of 2. The presence of cloud and filament is evident in the channels at velocities ranging from 1587 to 1637 km s⁻¹.

and tidal tail features were not included in the model fitting, as they are distinct features not described by the thin or thick disc models. Some of the prominent features that were considered while modelling were the spacing and orientation of the contours. Because of the high number of degrees of freedom, the parameters of the best-fitting model from TIRIFIC tend to have strong discontinuities, with abrupt changes in fitted values across the parameter space. Furthermore,



Figure 5. Left: the H I intensity map (moment 0) and velocity field (moment 1) of the cloud. The plotted H I contours are 3.3×10^{18} cm⁻² × 2^n , n = 0, 1, 2, 3,... Right: the velocity field of the cloud.



Figure 6. Top: de-projected rotation velocity, middle: position angle, bottom: inclination angle as a function of radius.

TIRIFIC employs χ^2 minimization, which inherently gives higher weight to regions with high signal-to-noise ratios. However, it is worth noting that these regions are not necessarily the most suited to distinguish between models. Therefore, while TIRIFIC's automated fitting capabilities were used when appropriate, manual intervention becomes necessary to guide the fitting process effectively (e.g. Oosterloo et al. 2007; Zschaechner et al. 2012; Kamphuis et al. 2013).

To determine whether the observed thickness of the HI disc can be attributed to a warp in the thin disc, we initially consider a simple model consisting of a single thin disc warped approximately 15° along the line of sight, which is our base model. Figs 8 and 9 show the bv diagrams and channel maps of both the observed data and the various models, respectively. The observed data is presented in the leftmost column, while the second column shows the model of a thin



Figure 7. Position velocity diagram at the resolution of 34.3 arcsec \times 25.6 arcsec (r15t00, r = 1.5 and t = 0) along the major axis of UGCA 250 (PA = 256°, $V_{\text{sys}} = 1705 \text{ km s}^{-1}$). Contours begin at 3σ and increase by factors of 2. The black circles represent the projected rotation velocities.

disc with a scale height ~ 200 pc, comparable to that observed in the Milky Way and other similar edge-on galaxies. The morphology of the channel maps and by diagrams, particularly near the systemic velocity, in this model is markedly different from the observed data, emphasizing that the model fails to accurately reproduce the observed features. In addition, we find that no amount of disc warping or disc thickness alone could reproduce the data adequately. We also tested increasing the velocity dispersion, which tended to thicken the emission from the channel maps in the vertical direction, but this change did not consistently match the morphologies observed in the channel maps.

We then add a substantial flare to the warped disc model (or base model), denoted as 'Flare+W' in the figures, by increasing the scale height with radius, which allows for a better fit to the bv diagrams and channel maps. The scale height increases from 0.2 kpc in the inner regions to 5 kpc in the outer regions. This modified model shows an improvement compared to the previous thin disc+warp model. The emission on the systemic sides of the bv diagrams and the thickening in the channel maps are now more accurately reproduced. The improvements due to the flare are most evident on the systemic velocity sides of the bv diagrams. However, shape, spacing, and orientation of the outer contours near the systemic side still show differences compared to the observed data.

We further explore a two-component model, consisting of a thin disc and a thick disc/halo, denoted as 'Halo+W' in the figures. HIhaloes have been observed in several galaxies, with NGC 891 being among the most notable (Oosterloo et al. 2007). For the thin disc, we assume a scale height of 0.2 kpc with a vertical sech² profile. The deprojected surface density profile of the thick disc is taken to be a scaled version of the thin disc profile. We test different scale heights for the thick disc. The morphology of bv diagrams depends on the scale-height of the thick disc and varying the scale heights between 2 and 4 kpc give broadly similar results, so we use a scale height of ~3 kpc for the thick disc. This is similar for NGC 3198 as presented in Gentile et al. (2013). This model significantly improves the reproduction of the thickening in the minor axis direction evident

However, the channel maps at the receding side velocities still do not match well with this model, which assumes that the halo is corotating with the thin disc, i.e. there is no lag. To address this, we introduced a lag that varies with radial distance from the centre, now labelled as 'Halo+WL' in the figures. This lag causes the observed velocity of the gas to be shifted in the direction of the systemic velocity, with this shift increasing with height. While the by diagrams of the no lag and lag models look very similar, we were mainly able to distinguish these models using the channel maps. The fine-tuned model shows a much better match in the channel maps at terminal side velocities, especially on the receding side (see last row of Fig. 9). The lag in these central regions is about $20 \text{ km s}^{-1} \text{ kpc}^{-1}$ starting at radius, $R \sim 4.6$ kpc, remaining constant out to $R \sim 6.6$ kpc, then decreasing to zero at $R \sim 20$ kpc in the outer regions on the approaching side. Conversely, on the receding side, the lag is about $20 \,\mathrm{km \, s^{-1} \, kpc^{-1}}$ starting at $R \sim 2.6$ kpc, remaining constant out to $R \sim 8.6$ kpc, then decreasing to zero at $R \sim 16.5$ kpc in the outer regions.

Additionally, while models with uniform lag throughout were considered, the fine-tuned model with radial variation proved to provide a better description of the data. Finally, we explored a model incorporating radial motions of the anomalous gas in the outer regions, indicated as 'Halo+WLO' in the figures. This model is similar to the fine-tuned model, but it includes a radial motion outward from the galaxy at 10 km s^{-1} in the outer regions. This adjustment accurately reproduces the direction of the outermost contour in the receding half of the by diagrams, demonstrating its effectiveness in modelling specific features of the galaxy. To determine the orientation of the disc, we used the optical image, which shows distinct dust lanes offset to the North; this offset indicates that the northern rim of the galaxy is the near side. This is crucial, as the interpretation of radial motions, whether characterized as outflows or inflows, relies on accurately identifying the near and far sides (Marasco et al. 2019).

4 DISCUSSION

The key question of this work concerns the origin of the extraplanar and anomalous gas in UGCA 250. Specifically, the EPG has been proposed to originate either from processes within the galactic disc, such as superbubbles generated by supernovae and stellar winds, or from external mechanisms, including interaction with nearby companions or the accretion of primordial gas from the cosmic web (Fraternali & Binney 2006; Sánchez Almeida et al. 2014; Marasco et al. 2019). Understanding the origin of this extraplanar and anomalous gas is crucial for unraveling the processes governing the dynamics and evolution of galaxies.

Our detailed modelling of UGCA 250 suggests that approximately 15 ± 5 per cent of the galaxy's HIcontent is distributed within a diffuse halo with a scale height of 3 ± 1 kpc. This result is consistent with Marasco et al. (2019), which shows that extraplanar H I is nearly ubiquitous in star forming disc galaxies, typically making up about about 5–25 per cent of the total HI mass, with a mean value of 14 per cent, and a typical thickness of a few kpc. A noticeable warp is also evident, not only seen in the zeroth-moment map, but is also captured in our extensive modelling. The fine-tuned model, which incorporates a radially varying lag – approximately 20 km s⁻¹ in the central regions that tapers to zero at $R \sim 20$ kpc – appears to best fit the data. We found no substantial evidence supporting the accretion of primordial gas, as there are no clear kinematical



begin at 3σ (Data); one f a thin disc Halo+WL);

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Figure 8. Minor axis position–velocity (PV) diagrams of UGCA 250 and models. Slice locations are shown in the left panel of Fig. 3. Contours begin at 3σ and increase by factors of 2. All the models include a warp across the line of sight. From the left: the data at 34.3 arcsec × 25.6 arcsec resolution (Data); one component model made of a thin disc (scale height ~ 200 pc) with a warp (Warp); a flare with a warp (Flare+W); two component model made of a thin disc and a corotating halo (scale height ~ 3 kpc) with a warp (Halo+W); a two component model made of a thin disc and a lagging halo and warp (Halo+WL); two component model with a warp, lagging halo, and a radial motion outwards in the outer regions of galaxy (Halo+WLO). The rightmost column displays the moment 2 map, with the line in each panel indicating the slice location corresponding to the by diagram in the respective row.

signatures of inflowing gas in our data. However, given the complex dynamics arising from tidal interactions and the presence of nine satellite galaxies, detecting such signatures, even if they exist, would be challenging. The most definitive way to identify primordial gas accretion is through metallicity measurements, as they can reveal the low metallicities, characteristic of such gas. Unfortunately, these measurements are not available for this study, limiting our ability to conclusively determine the presence or absence of primordial gas. Additionally, there is some evidence of radial motions in the outer regions of the receding side, with an outward direction. Models using outflows only in the central regions, however, could not reproduce the data accurately, particularly the outermost contour in the by diagrams (as seen in Fig. 8). This indicates these outflows might not result from star formation. Instead, they could be linked to radial motions triggered by potential tidal interactions, which is further corroborated by the detection of nine satellite galaxies and intra-group gas.

Despite extensive modelling, there remains a significant amount of H1 that is spatially or kinematically distinct from the disc gas. To isolate this emission, we blank the data where the model values exceed 0.5σ , where σ is the noise in the observed cube. This results in a 'residual' data cube from which we can generate the moment maps of the anomalous gas, i.e. the gas not described by either the thin or thick disc models. These maps are presented in Fig. 10. The Minor axis offset



Major axis offset

Figure 9. Channel map comparison of UGCA 250 observations with various models. The heliocentric velocities are shown on the top left of data frames. All channel maps are aligned such that their centres match, with a white dashed line drawn across the centre for all maps. From the left: the data at 34.3 arcsec \times 25.6 arcsec resolution (Data); one component model made of a thin disc (scale height \sim 200 pc) with a warp (Warp); a flare with a warp (Flare+W); two component model made of a thin disc and a corotating halo (scale height \sim 3 kpc) with a warp (Halo+W); a two component model made of a thin disc and a lagging halo and warp (Halo+WL); two component model with a warp, lagging halo, and a radial motion outwards in the outer regions of galaxy (Halo+WLO).

total H I mass of this anomalous gas is approximately $10^{7.7}$ M_☉, which accounts for about 0.7 per cent of the total mass of the galaxy. Although this is a relatively small fraction and may not significantly impact the overall H I content, it could be tracing underlying ionized gas and is therefore relevant. We observe that the thin disc and most of the thick disc are well modelled. However, three prominent features of the anomalous gas are identified: (i) The primary feature is a taillike filament extending away from the disc in the southern direction up to approximately 34 kpc (in projection). Below this filament is a blob situated at a projected distance of about 41.5 kpc from the centre of the galaxy, with an H I mass of $\log(M_{\rm HI}) \sim 6.34$. The combined H I mass of the filament and this blob is $\log(M_{\rm HI}) \sim 7.01$. It is worth noting that satellite ID 3 (see Fig. 2) is located below

the filament and blob, at a distance of 65 kpc from the blob. (ii) Further, a cloud-like feature in the velocity field south of the galaxy exhibits anomalous velocities. Visual inspection of the cube reveals a cloud of gas coinciding with the main emission of the galaxy but at a different velocity, visible in the channel maps in Fig. 4. No optical counterpart for this cloud has been detected. (iii) There is a region with high velocity dispersion located in the northeast of the galaxy, on the approaching side.

In Fig. 11, the PV diagrams provide further insights into these features. The top left panel shows a PV diagram through the filament region of the galaxy, while the top right panel depicts a PV diagram through the filament and cloud regions. The bottom right panel presents a PV diagram across the high-velocity dispersion region and



Figure 10. Top: H_I gas surface density of the anomalous gas that was not modelled. Bottom: the velocity field of the anomalous gas, the colour shows the velocity in km s⁻¹.

the cloud, and the bottom left panel shows a PV diagram across the high-velocity dispersion region and the filament. Each panel contains an inset indicating the path of the PV slice on the corresponding moment maps. These PV diagrams through these anomalous features hint at a history of close encounters or mergers, leading to tidal disturbances that generated filament-like structures and cloud-like anomalies. The interconnected patterns across the cloud and filament regions support the hypothesis of a prior interaction.

It is important to note that these various features including the filament, cloud, and high velocity dispersion region were not detected with the 5.5-h track of MHONGOOSE observations (Blok et al. 2024), which have a 3σ column density sensitivity of approximately 6.7×10^{19} atoms per cm⁻². As illustrated in Fig. 2, at this column density sensitivity, the vertical distribution of the gas or thick disc is also less apparent. The full 55 h of MeerKAT data, however, reveal these features in significant detail, highlighting the advantages of its superior column density sensitivity. This has led to the detection of nine satellite galaxies, with six identified in the single track observations and three additional galaxies revealed in the full depth data. Furthermore, at lower resolutions and column densities, intragroup gas was also detected along with the satellites. Additionally, the shell elliptical galaxy NGC 3923, which is located about 30 arcmin (or 200 kpc) from UGCA 250, falls within our observation's field of view and has a redshift similar to UGCA 250's. This galaxy is known for its extensive system of stellar shells, the most of any shell galaxy,

which are thought to be remnants of past galactic interactions (e.g. Bílek et al. 2016; Spavone et al. 2017).

The zoomed-in moment 0 and moment 1 maps of the individual satellite galaxies are shown in the Appendix Figs A3 and A4. In these maps, we observe HI emission offset from the stellar counterpart in several cases (e.g. IDs 1, 4, 5, and 7), along with evident asymmetries and disturbed morphologies that indicate interactions within these systems. Further, all detected satellite galaxies lie within the 356 kpc field of view of MeerKAT, and, except for satellite ID 10, all are within $200 \,\mathrm{km \, s^{-1}}$ of the systemic velocity of UGCA 250, Hence, UGCA 250, NGC 3923, and all the detected satellites likely belong to the same group, as all are contained within the group's estimated virial radius of 591 kpc and virial mass of $10^{13.3}$ M_{\odot} (Kourkchi & Tully 2017). This suggests that the larger environment of UGCA 250, including its satellites and NGC 3923, appears to be a dynamically active system. The distinct anomalous gas features suggest tidal interactions or a past event involving something passing close to UGCA 250 as the likely source of this gas. This interpretation is further supported by the optical discontinuity observed in UGCA 250 (see Fig. 1). The discontinuity is likely associated with the main disc rather than a separate dwarf galaxy, as no HI emission at a distinct velocity is detected at that location. Further, the observation of satellite galaxies and detection of intra-group gas indicate potential close encounters or mergers that could create tidal disturbances manifesting as filaments and clouds with peculiar velocities. The outward radial motions in the outer receding regions further point to these interactions shaping the unique EPG kinematics and structure. A more detailed study of the distribution and interactions of the satellite galaxies is beyond the scope of this paper.

While environmental effects likely play a key role in shaping the EPG in UGCA 250, the contribution of stellar feedback must also be considered. UGCA 250 is a regular star-forming galaxy on the main sequence, and to quantify the role of stellar feedback, we follow the framework of Marasco et al. (2019), applying their equation 11 with a mass-loading factor $\beta = 7$ to estimate the expected EPG mass from a galactic fountain flow model. This yields $\log(M_{\rm EPG}/M_{\odot}) \sim$ 8.6, about 0.4 dex lower than our observed value. This discrepancy suggests that, while stellar feedback likely accounts for a significant portion of the thick disc, tidal interactions are responsible for the more extended features, such as the filament and cloud, UGCA 250 shares several characteristics with NGC 2403, including rotational velocity, star formation rate, and EPG properties, where a galactic fountain model interacting with the hot CGM accounts for much of the EPG (Li et al. 2023). Additionally, NGC 2403 is part of a group of five galaxies (Kourkchi et al. 2022), with some of its extraplanar HIattributed to remnant tidal interactions with a satellite (Veronese, de Blok & Walter 2023). We therefore propose a mixed origin for the EPG in UGCA 250, with the thick disc primarily shaped by stellar feedback and the outer structures influenced by tidal interactions. This study highlights the complex interplay between internal galactic processes and external environmental factors shaping the anomalous gaseous structures around UGCA 250.

5 CONCLUSIONS

We use ultra-deep MHONGOOSE survey observations to investigate the EPG in UGCA 250. The survey's combination of high column density sensitivity and high spatial resolution – ranging from 6.0×10^{19} cm⁻² at 6 arcsec to 5.7×10^{17} cm⁻² at 90 arcsec (3σ over 16 km s^{-1}). This offers an ideal framework for examining the physics of EPG and the disc–halo connection. This has enabled the detection of nine satellite galaxies, a distinctive filament extending



Figure 11. Top left: PV diagram through the filament region of the galaxy. The inset in the top right corner shows the moment 1 map with a black dashed line indicating the PV slice path. Top right: PV diagram through the filament and cloud regions. The inset also displays the moment 1 map with the black dashed path. Bottom left: PV diagram across the high velocity dispersion region and filament. The inset shows the moment 2 map with the black dashed path. Bottom right: PV diagram across the high velocity dispersion region and the cloud. The inset again displays the moment 2 map with the same black dashed line. The horizontal colour bar represents H I column density for the PV diagrams, while the vertical colour bar represents either velocity (moment 1) or velocity dispersion (moment 2) in km s⁻¹. In the maps, C denotes the cloud at anomalous velocities, F indicates the filament, B refers to the blob below the filament, and H represents high velocity dispersion.

approximately 41 kpc southward from the galaxy, and a cloud at anomalous velocities. We have not detected any counterpart for this cloud in the VST deep optical image.

We construct a detailed tilted ring model for this edge-on galaxy to gain a deeper understanding of the vertical structure of its neutral hydrogen. Through detailed modelling, we estimate that approximately 15 ± 5 per cent of the galaxy's H I content resides in a halo extending 3 ± 1 kpc above the galactic plane. This halo exhibits a lag that varies with radial distance from the centre, beginning at 20 km s⁻¹ within the central regions and tapering to zero at a radial distance of 20 kpc. Additionally, we found evidence for radial motions in the outward directions in the outer regions of the receding half of the galaxy. The outermost contours in the by diagrams can only be reproduced by incorporating outflows in the outer regions. Since there is little star formation occurring in these outer regions, these outflows are likely caused by tidal interactions rather than stellar feedback. The combination of radial motions in the outer regions, multiple satellite galaxies, and distinct features such as the cloud, filament, and high-velocity dispersion region suggests that tidal interactions are at least in part responsible for the formation of anomalous gas in UGCA 250. Therefore, our analysis suggests a mixed origin for the EPG in UGCA 250, with contributions from both internal processes, such as stellar feedback, and external influences, including tidal interactions with satellite galaxies.

ACKNOWLEDGEMENTS

SK, DJP, NZ, and SHAR are supported by the South African Research Chairs Initiative of the Department of Science and Technology and National Research Foundation (grant number 77825). This work has received funding from the European Research Council (ERC) under the European Union's Horizon 2020 research and innovation programme (grant agreement no. 882793 'MeerGas'). PK and RJD are partially supported by the BMBF project 05A23PC1 for D-MeerKAT. JH acknowledges support from the UK SKA Regional Centre (UKSRC). The UKSRC is a collaboration between the University of Cambridge, University of Edinburgh, Durham University, University of Hertfordshire, University of Manchester, University College London, and the UKRI Science and Technology Facilities Council (STFC) Scientific Computing at RAL. The UKSRC is supported by funding from the UKRI STFC. EA and AB acknowledge support from the Centre National d'Etudes Spatiales (CNES), France. LC acknowledges financial support from the Chilean Agencia Nacional de Investigación y Desarrollo (ANID) through Fondo Nacional de Desarrollo Científico y Tecnologico (FONDECYT) Regular Project 1210992. LVM acknowledges financial support from the grant PID2021-123930OB-C21 funded by MICIU/AEI/ 10.13039/501100011033 and by ERDF/EU and grant CEX2021-001131-S funded by MICIU/AEI.

The MeerKAT telescope is operated by the South African Radio Astronomy Observatory, which is a facility of the National Research Foundation, an agency of the Department of Science and Innovation. We acknowledge the use of the ilifu cloud computing facility – www.ilifu.ac.za, a partnership between the University of Cape Town, the University of the Western Cape, the University of Stellenbosch, Sol Plaatje University, the Cape Peninsula University of Technology, and the South African Radio Astronomy Observatory. The Ilifu facility is supported by contributions from the Inter-University Institute for Data Intensive Astronomy (IDIA a partnership between the University of Cape Town, the University of Pretoria, the University of the Western Cape and the South African Radio Astronomy Observatory), the Computational Biology division at UCT and the Data Intensive Research Initiative of South Africa (DIRISA).

DATA AVAILABILITY

The moment maps of MHONGOOSE observations of UGCA 250 are available for download at the MHONGOOSE website (https://mhongoose.astron.nl).

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APPENDIX A: PV DIAGRAMS OF UGCA250

In this appendix, we present the bv and lv diagrams of UGCA 250 (see Figs A1 and A2). The bv diagrams, extracted along the minor axis with slices at 30 arcsec intervals on both the approaching and receding halves, clearly display features such as the cloud and the filament extending from the disc. For the lv diagrams along the major axis, an inset is included that shows the slice across the moment 1 map used to generate the lv diagram. In the lv diagrams along the major axis, the galaxy appears more extended toward the northern, receding half, further underscoring an asymmetry in its structure. Further, we show the moment 0 and moment 1 maps of individual satellite galaxies in Figs A3 and A4.



Figure A1. The bv diagrams of UGCA 250 at the resolution of 34.3 arcsec \times 25.6 arcsec (r15t00, r = 1.5 and t = 0). The diagrams correspond to slices taken at offsets of 30, 60, 90, 120, 150, 180, 210, and 240 arcsec from the centre, separately for the approaching and receding halves of the galaxy. Contours begin at 3σ and increase by factors of 2. The outermost contour of each map is the 3σ contour at the resolution of 65.3 arcsec \times 63.8 arcsec (r05t60, r = 0.5 and t = 60 arcsec).



Figure A2. The top left panel shows a by diagram of UGCA 250, while the other panels display ly diagrams. All panels are at the resolution of 34.3 arcsec \times 25.6 arcsec (r15t00, r = 1.5 and t = 0). Contours begin at 3σ and increase by factors of 2. Each panel contains an inset indicating the path of the PV slice on the corresponding moment maps.



Figure A3. The HI intensity map (moment 0) and velocity field (moment 1) of satellite galaxies. The HI contours are at levels of 1.2×10^{19} cm⁻² $\times 2^n$, n = 0, 1, 2, 3.



Figure A4. The H I intensity map (moment 0) and velocity field (moment 1) of satellite galaxies. The H I contours are at levels of 1.2×10^{19} cm⁻² $\times 2^n$, n = 0, 1, 2, 3.

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