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DARK ENERGY SPECTROSCOPIC INSTRUMENT (DESI) SURVEY YEAR 1 RESULTS

DESI 2024 VII: cosmological constraints from the full-shape modeling of clustering measurements



**DARK ENERGY
SPECTROSCOPIC
INSTRUMENT**

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ABSTRACT: We present cosmological results from the measurement of clustering of galaxy, quasar and Lyman- α forest tracers from the first year of observations with the Dark Energy Spectroscopic Instrument (DESI Data Release 1). We adopt the full-shape (FS) modeling of the power spectrum, including the effects of redshift-space distortions, in an analysis which has been thoroughly validated in a series of supporting papers as summarised in [1]. We combine the full-shape information with DESI's DR1 constraints from the baryon acoustic oscillations (BAO) of these tracers. In the flat Λ CDM cosmological model, DESI (FS+BAO), combined with a baryon density prior from Big Bang Nucleosynthesis and a weak prior on the scalar spectral index, determines matter density to $\Omega_m = 0.2962 \pm 0.0095$, and the amplitude of mass fluctuations to $\sigma_8 = 0.842 \pm 0.034$. The addition of the cosmic microwave background (CMB) data tightens these constraints to $\Omega_m = 0.3056 \pm 0.0049$ and $\sigma_8 = 0.8121 \pm 0.0053$, while further addition of the joint clustering and lensing analysis from the Dark Energy Survey Year-3 (DESY3) data further improves these measurements, and leads to a 0.4% determination of the Hubble constant, $H_0 = (68.40 \pm 0.27) \text{ km s}^{-1} \text{ Mpc}^{-1}$. In models with a time-varying dark energy equation of state parametrised by w_0 and w_a , combinations of DESI (FS+BAO) with CMB and type Ia supernovae continue to show the preference, previously found in the DESI DR1 BAO analysis, for $w_0 > -1$ and $w_a < 0$ with similar levels of significance. DESI data, in combination with the CMB, improve the upper limits on the sum of the neutrino masses relative to the case when only the DR1 BAO was available, giving $\sum m_\nu < 0.071 \text{ eV}$ at 95% confidence. We finally constrain deviations from general relativity represented by two modified gravity parameters. DESI (FS+BAO) data alone measure the

parameter that controls the clustering of massive particles, $\mu_0 = 0.11^{+0.45}_{-0.54}$, in agreement with the zero value predicted by general relativity. The combination of DESI with the CMB and the clustering and lensing analysis from DESY3 constrains both modified-gravity parameters, giving $\mu_0 = 0.04 \pm 0.22$ and $\Sigma_0 = 0.044 \pm 0.047$, again in agreement with general relativity.

KEYWORDS: dark energy experiments, redshift surveys, power spectrum, modified gravity

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1 Introduction

The large-scale structure (LSS) of the Universe, as probed by galaxy surveys and the intergalactic medium, has firmly established itself as a reliable probe of cosmology and fundamental physics. The three-dimensional clustering of tracers of the LSS — galaxies, quasars, and Ly α absorption signatures in quasar spectra — can be directly related to cosmological theory. This, in turn, can be used to constrain some of the most familiar quantities in cosmology, including the amount of dark matter and dark energy, the amplitude and spectral index of primordial density perturbations, spatial curvature, and neutrino mass. The progress in such clustering measurements over the last half century has been nothing short of remarkable [2–11], and has resulted in percent-level constraints on some of the aforementioned cosmological parameters.

One prominent feature in the galaxy clustering correlation is the baryon acoustic oscillations (BAO), an oscillatory signature which appears as “wiggles” in the galaxy power spectrum, or a single localised peak in the galaxy correlation function. The scale of the BAO

feature is determined by the sound horizon at the baryon drag epoch, and its observation via tracers at a given redshift z contains information about the ratio of this scale to distance measures ($D_A(z)$ and $c/H(z)$ in the directions perpendicular and parallel to the line-of-sight, respectively), thus containing key cosmological information. More information, however, is available in the “full shape” of the clustering signal, specifically the measured power spectrum $P(k, z)$ over a range of wavenumbers k and tracer redshifts z or, equivalently, the correlation function $\xi(r, z)$ where r is the comoving separation.

Notably, the dependence of the full-shape clustering signal on redshift z informs us about the growth of cosmic structure (e.g. [12–14]), which in turn is very sensitive to the properties of dark energy and modified gravity, and to the total matter content of the universe. The growth-rate constraints allow data to test the underlying theory of gravity at cosmological scales and discriminate between models that share the same expansion history, see e.g. [15–18]. Because DESI measurements are sensitive to both the geometrical quantities and the growth of density perturbations, they are particularly well-suited to supply tests of dark energy and modified gravity. The full-shape clustering signal also contains information about the amplitude and shape of the primordial power spectrum, and hence provides information complementary to that from the cosmic microwave background (CMB) measurements.

The above-mentioned long history of measurements of galaxy clustering has, over the past decade or so, been reinvigorated with the data from the Baryon Oscillation Spectroscopic Survey (BOSS) [19] which has been part of the third phase of the Sloan Digital Sky Survey (SDSS-III; [20]), and its extension eBOSS [21]. Full-shape analyses of galaxy and quasar clustering in BOSS have been carried out by the BOSS [22–26] and eBOSS [27–33] collaborations, as well as independent teams who typically studied BOSS and/or eBOSS data [34–46]. The tools developed in these analyses have enabled reliable extraction of cosmological information from clustering. Cosmological results are also expected from the Hobby-Eberly Telescope Dark Energy Experiment (HETDEX) Survey [47].

The Dark Energy Spectroscopic Instrument (DESI) is the first Stage-IV galaxy survey in operation [48–51]. It is conducting a spectroscopic five-year survey over 14,200 square degrees that will collect spectra of about 40 million galaxies and quasars [52–57]. DESI targets five separate tracers: low-redshift galaxies from the Bright Galaxy Survey (BGS) [58], luminous red galaxies (LRG) [59], emission line galaxies (ELG) [60], quasars (QSO) [61], and the Ly α forest features in quasar spectra [62]. DESI’s deep redshift coverage, $0 < z < 4$, will enable it to map out the expansion history and the growth of cosmic structure to high precision. The principal scientific goals of DESI are to obtain tight constraints on dark energy, neutrino mass, and primordial non-Gaussianity. This is complemented by a tremendous amount of other science that is being carried out using data from the DESI instrument.

This paper is part of a series discussing key cosmology results from the first year of observations from DESI, which is based on DESI Data Release 1 (DR1; [63]). This is the second paper that focuses on key cosmological parameter measurements from DESI DR1; in the first paper [64], we presented cosmological measurements from the information in baryon acoustic oscillations in DESI DR1 data, based on the analysis of galaxy and quasar clustering [65], and that in the Ly α forest data [66]. In this paper, we significantly extend those results by complementing the BAO information with the “full-shape” analysis which

models the overall clustering signal of DESI tracers across time and space, and report the resulting cosmological constraints from the combined BAO + full-shape analysis.

The data behind the analysis, and the plans for their release, are presented in [63], while the galaxy/quasar samples are discussed in detail in [67]. The large-scale structure catalogs are fully described in [67, 68]. The DESI DR1 galaxy full-shape analysis, its detailed pipeline choices, the study of systematics, and the cosmological constraints on the Λ CDM model from DESI DR1 galaxy full-shape alone and its combination with BAO are all presented in [1]. A further detailed analysis of modified gravity models is presented in [69]. Moreover, a number of technical details, as well as in-depth discussions and justification arguments for our analysis choices are provided in a series of supporting papers. Specifically, [70–74] provide details and validation of the perturbation theory codes that we use to analyse the (pre-reconstruction) galaxy power spectrum, while [1, 75] show the level of agreement between the codes in a series of controlled settings with simulated and synthetic data vectors. These papers also discuss the role of priors and the projection effects that can arise when presenting high-dimensional posteriors marginalised to show constraints in lower-dimensional parameter spaces of interest. The covariance matrices are described and validated in [76–78]. Our systematic error budget relies on studies which are presented in [74, 79, 80] and summarised in [1]. Throughout the analysis we have made use of a series of mock catalogs described in detail in [81]. Note that the constraints on primordial NG will be presented separately in [82].

2 Data and methodology

In this section we describe the essential inputs to the cosmological analysis — data and methodology. In section 2.1 we describe the data, full-shape measurement methodology, and the blinding procedure that we applied to the measurements. In section 2.2, we describe the external data that we optionally combine with DESI in the analysis. Finally, section 2.3 describes theoretical modeling, as well as the likelihood analysis and other details of our cosmological inference pipeline.

2.1 DESI full-shape measurements

2.1.1 DESI DR1 data

The DESI data that we use are described in [67]. They are derived from the redshifts and positions of over 4.7 million unique galaxies and QSOs over a $\sim 7,500$ square degree footprint¹ covering the redshift range $0.1 < z < 2.1$. These discrete tracers are broken into four target classes: 300,017 galaxies from the magnitude-limited bright galaxy survey (BGS); 2,138,600 luminous red galaxies (LRG); 1,415,707 emission line galaxies (ELG)² and 856,652 quasars (QSO) (see table 1 of [1]). These tracers are split into six redshift bins: one bin with the BGS ($0.1 < z < 0.4$), three bins with the LRGs ($0.4 < z < 0.6$, $0.6 < z < 0.8$, and $0.8 < z < 1.1$), one bin with the ELGs ($1.1 < z < 1.6$), and one redshift bin with

¹Note that the sky coverage for individual tracers may be substantially lower than $\sim 7,500$ sq. deg. due to masks and cuts; see [67] for details.

²The DESI DR1 sample contains a total of 2,432,022 ELGs in two redshift bins, but the ELGs in the lower redshift bin (1,016,365 objects) did not pass the systematics checks [1], so we do not use them in the cosmological analysis.

the QSOs ($0.8 < z < 2.1$). These objects are assembled into large-scale structure catalogs, and the power spectrum in each redshift bin is subsequently computed as discussed below; see [67] and references therein for all details.

In addition to the discrete tracers described above, DESI also uses the spectra of distant QSOs to measure large-scale structure in the intergalactic medium (i.e., the Ly α forest). Measurements of the 3D correlation function of the DR1 Ly α forest data are presented in [66]. At present we only use the baryon acoustic oscillation information in the large-scale clustering of the Ly α forest to constrain the background geometry [66], and do not provide a measurement of growth. For this reason the Ly α forest measurements only enter via their contribution to constraining the expansion history.

2.1.2 DESI full-shape measurements

The goal of our analysis is to extract cosmological information beyond the BAO feature from the measurements [1, 67] of the full-shape clustering of DESI tracers. To that effect, we measure the first few multipole moments of the Fourier-space tracer power spectra relative to the line-of-sight to the observer — the monopole, quadrupole and hexadecapole — which quantify the information imprinted by redshift-space distortions (although we limit our analysis here to the first two of these multipoles). These measurements are obtained with the estimator from [83]. The power spectrum measurements are obtained from the galaxy catalogs (“data”) and from synthetically-generated catalogs with random distribution of points (“randoms”) to which we assign the same selection as for the data, including assigning weights (to points) that account for systematic corrections, and those that implement the Feldman-Kaiser-Peacock (FKP) optimal weighting scheme [84]. We also use the random catalog to compute the window matrix [85, 86] that relates the measured power spectrum multipoles to the theory power spectrum prediction.

Data and random catalogs are constructed as described in [67]; they are both masked for the presence of bright objects, lack of or bad imaging data and spectroscopic observations, and target priorities. Fiber assignment results in variations of the observed density of tracers; this effect is corrected by applying the completeness weights at the catalog level. Despite this correction, fiber assignment impacts the two-point statistics at small angular separations, which we consequently remove from the power spectrum estimation [86]. Both the small-angle structure in the masks and the small-scale angular cuts result in a window matrix that has contributions extending to very small scales; we then “rotate” [86] our power spectrum measurement, covariance matrix, and window matrix to make the latter more diagonal. Imaging systematics (due to galactic dust, imaging depth, and a host of other reasons) are corrected for by systematic weights at the catalog level. The imaging template-fitting techniques used (based on random forest or neural nets) damp large-scale angular modes: we measure this “angular integral constraint” effect in mock realisations, and remove it from the power-spectrum measurements. Finally, the radial selection function imprinted in the random catalog is directly inferred from the observed data, resulting in a “radial integral constraint” which is similarly estimated from mocks and corrected at the power spectrum level. The power spectrum covariance matrix is estimated from a set of 1000 fast approximate mocks (EZmocks, [81]) and rescaled to make the mock-based covariance matrix of the two-point

correlation function. This rescaling, which is enacted by multiplying each element of the covariance matrix by a factor of between 1.11 and 1.39 depending on the tracer, serves to correct the `EZmocks` covariance for certain aspects of DR1 data (e.g. fiber assignment) that are more accurately described by a semi-empirical covariance prediction (RASCALC, [77]). The details of this battery of validation tests are presented in [67] and references therein. Based on these tests, we only use the monopole and quadrupole in our cosmological analysis, and restrict the full-shape analysis to the wavenumber range $0.02 < k/h \text{ Mpc}^{-1} < 0.20$ [1], with a binning width of $\Delta k = 0.005 h \text{ Mpc}^{-1}$. The lower end of the usable wavenumber range is chosen to prevent potential residual systematics which primarily impact large scales [67], while the upper limit, $k_{\text{max}} = 0.2 h^{-1} \text{ Mpc}$, is determined by detailed validation tests of the full-modeling analysis of the power spectrum and its precise value has little impact on the cosmological constraints [70–72].

2.1.3 DESI full-shape blinding

An essential part of our analysis framework was “blinding” the results during the period where data-selection and analysis choices were being made, to avoid the risk of confirmation bias. The blinding process has two components: blinding of the BAO, and blinding of the redshift-space distortions. The blinding procedure is performed at the catalogue level, and was applied consistently to both the BAO-only analysis [64] as well as the full-shape analysis in this paper. The BAO aspect of the blinding procedure follows the work of [87], where the redshifts of the observed galaxies are modified so that they imprint a shift in the anisotropic position of the BAO peak. The redshift-space distortion aspect of blinding, designed by [87] to render the cosmological information about the growth of structure impervious to confirmation bias, is achieved by applying a shift in the growth rate f . A full description of the blinding technique, and how it has been tailored to DESI needs, can be found in [88].

2.2 External data

We now describe the external datasets we combine with the DESI (FS+BAO) measurements. These choices largely follow the DESI DR1 BAO analysis [64], with the important addition of angular clustering and lensing data from the Dark Energy Survey.

We adopt the cosmic microwave background (CMB) data from the official *Planck* (2018) PR3 release [89]. We use as our baseline the temperature (TT) and polarisation (EE) auto-spectra, plus their cross-spectra (TE), as incorporated in the `simall`, `Commander` (for multipoles $\ell < 30$) and `plik` (for $\ell \geq 30$) likelihoods. As part of our robustness tests for constraints on the neutrino mass, we also alternatively consider two independent analyses of the latest *Planck* PR4 data release: the high- ℓ `CamSpec` likelihood [90, 91], and the `LoLLiPoP` (low- ℓ) and `HiLLiPoP` (high- ℓ) likelihoods [92, 93]. We complement the CMB likelihood with the information from the reconstruction of the lensing power spectrum as measured using the connected 4-point function of the CMB temperature and polarisation. We adopt data from the combination of `NPIPE` PR4 *Planck* CMB lensing reconstruction [94] and the Data Release 6 of the Atacama Cosmology Telescope (ACT) [95–97].³ In what follows, we will denote results obtained using temperature and polarisation information from *Planck*, and CMB

³The likelihood is available from https://github.com/ACTCollaboration/act_dr6_lenslike.

lensing information from the *Planck*+ACT combination, simply as “CMB”. Where necessary, we will explicitly label results that do not use CMB lensing reconstruction as “CMB-nl”.⁴

In the analyses that do not include the CMB information, we also add the prior on the physical baryon density, $\Omega_b h^2$, coming from Big Bang Nucleosynthesis (BBN). The theoretical BBN prediction for the abundances of light elements, especially deuterium (D) and Helium (^4He), is sensitive to the baryon density. Measurements of these abundances therefore lead to a constraint on the baryon density, but one which depends on details of the theoretical framework, particularly the crucial input of nuclear interaction cross-sections. As we did in our BAO-only paper [64], we adopt a recent analysis [98] that makes use of the new `PRyMordial` code [99] to recompute the predictions while marginalising over uncertainties in the reaction rates. We adopt the joint constraint on $\Omega_b h^2$ and the number of relativistic species N_{eff} , and fix the latter parameter to its fiducial value of 3.044 in models where we are not allowing for the presence of additional light relics.⁵

We also add information from type Ia supernovae (SN Ia), which serve as standardisable candles offering an alternative way to measure the expansion history of the universe. Here we utilise the same three SN Ia datasets that we studied in the DESI DR1 BAO paper; these are the largest compilations of supernova data that have been consistently reduced and analysed. The first SN Ia dataset we consider is the PantheonPlus⁶ compilation [100], with 1550 spectroscopically-confirmed SN Ia in the redshift range $0.001 < z < 2.26$, where we use the public likelihood from [101]. The second SN Ia dataset that we adopt is the Union 3 compilation [102], containing 2087 SN Ia in the redshift range $0.01 < z < 2.26$, 1363 of which are in common with PantheonPlus, and which uses a likelihood analysis and treatment of statistical and systematic errors based on Bayesian hierarchical modelling. The third SN Ia dataset is the Year 5 supernova analysis from the Dark Energy Survey (henceforth “DES-SN5YR”) [103]. This analysis starts with a homogeneously-selected sample of 1635 photometrically-classified SN Ia with redshifts $0.1 < z < 1.13$. This is complemented by 194 low-redshift SN Ia (which are in common with the PantheonPlus sample [100]) spanning $0.025 < z < 0.1$. We include all three SN Ia datasets in our analysis; however, in certain cases where there is no meaningful dependence of the result on the choice of SN Ia data, we only adopt one of the three datasets to avoid unnecessary redundancy.

Additionally, we consider external information from the combination of angular galaxy clustering and weak gravitational lensing — the so-called “3 × 2-pt” datavector that consists of three two-point correlation functions (galaxy-galaxy, galaxy-shear, and shear-shear). We use results from the Dark Energy Survey Year-3 (DESY3) analysis [104], which is based

⁴For clarification, the TT, EE and TE power spectra always include the effect of gravitational lensing; here we emphasise that our fiducial CMB dataset additionally includes the CMB lensing reconstruction, while the CMB-nl version does not.

⁵To be precise, the joint constraint on $\Omega_b h^2$ and N_{eff} has the respective mean values (0.02196, 3.034), and the corresponding covariance (<https://tinyurl.com/29vzc592>)

$$\mathbf{C} = \begin{bmatrix} 4.03112260 \times 10^{-7} & 7.30390042 \times 10^{-5} \\ 7.30390042 \times 10^{-5} & 4.52831584 \times 10^{-2} \end{bmatrix}.$$

⁶We denote the originally named Pantheon+ dataset as PantheonPlus in order to avoid ambiguities with the ‘+’ symbol used to denote the combinations of datasets.

on observations of about 100 million source galaxies, and about 10 million lens galaxies in the fiducial MagLim sample, over a footprint of 4143 square degrees. The DESY3 analysis employs photometric redshifts to identify the “source” galaxies and divide them into four tomographic bins, and the “lens” galaxies that are subdivided into six redshift bins, although the two highest-redshift lens samples are not used in the fiducial analysis. The positions of lens galaxies are used to compute the galaxy angular clustering signal (i.e. galaxy-galaxy correlations); the shear of source galaxies is used to measure cosmic shear (i.e. shear-shear correlations); and finally the shear of source galaxies correlated with positions of the lens galaxies gives the shear-galaxy correlations. The dataset also includes the ratio of galaxy-shear correlations at small scales in the so-called shear-ratio data vector. The DES 3×2 -pt analysis mitigates information from scale-dependent bias, baryon feedback effects and nonlinearities, which are challenging to model sufficiently accurately, using a combination of scale cuts and theoretical modeling using halofit [105, 106]. The analysis also marginalises over nuisance parameters that encode imperfect knowledge of certain astrophysical effects (such as galaxy biases, photo- z distribution shifts, intrinsic galaxy alignments and multiplicative shear biases in each source tomographic bin). For modified gravity, when we use DES 3×2 -pt data, we employ a different likelihood tailored to modified-gravity (MG) analysis with similar conservative scale cuts as imposed in the DESY3 MG study [107]. We assume that the DESY3 (3×2 -pt) data are uncorrelated with DESI (FS+BAO).

In addition to the DESY3 3×2 -pt data, we also make use of the so-called “ 6×2 -pt” data from DES Y3, which complement the galaxy clustering, cosmic shear, and galaxy-galaxy lensing with the information from the CMB lensing. Specifically, the 6×2 -pt datavector extends the 3×2 -pt one by further adding the galaxy-CMB lensing, shear-CMB lensing and CMB lensing-CMB lensing two-point correlation functions. We adopt the data vector from the DESY3 6×2 -pt analysis [108] which uses CMB lensing data from *Planck* and around 1,800 square degrees of the South Pole Telescope (SPT) [109] footprint. We use the same modelling and scale cut choices as the DESY3 analyses. When combining the DESY3 6×2 -pt likelihood with the CMB, we use the CMB data without lensing (CMB-nl) in order to avoid double-counting the CMB lensing information.

2.3 Modeling and likelihood

Having described the DESI and external datasets, we now discuss the likelihood pipeline, including the parameter space that we constrain and other details of cosmological inference. We start with a brief overview of how we theoretically model the power spectra in our “Full Modeling” approach.

2.3.1 Full-shape modeling approach

We use a perturbation-theory approach to full-shape clustering analysis (referred to as Full Modeling here and in companion papers). [Ref. [1] also describes another approach called ShapeFit [110], which we use for testing and validation of our pipeline, but not for producing the cosmological results in this paper.]

The idea behind Full Modeling is to directly fit a model to the full-shape power spectrum multipoles [34, 35, 37]. In this approach, we model the linear matter power spectrum using a

set of cosmological parameters (see section 2.3.3), and complement it with a set of nuisance parameters that describe the anisotropic power spectrum in the mildly nonlinear regime as well as various astrophysical or instrumental systematic uncertainties (e.g. galaxy bias).

Our theoretical model for two-point galaxy clustering is built around cosmological perturbation theory (PT; [111, 112]). Within PT, the growth of structure is treated systematically by expanding order-by-order in the amplitude of the initial fluctuations, with nonlinearities at small scales encoded using a series of “counterterms” that are constrained by the symmetries of the equations of motion (often known as “effective-field theory techniques” [112]). Biased tracers of large-scale structure, like galaxies or neutral hydrogen in the intergalactic medium, are treated in a consistent manner by identifying the contributions to their clustering signature allowed by fundamental symmetries at each order in PT [113]. Currently the redshift-space power spectrum of galaxies can be modeled with accuracy well beyond the expected statistical uncertainty in DESI, with the models being extensively tested against simulations [114, 115], compared to each other, and tested on earlier surveys such as BOSS and eBOSS. We have tested and compared several perturbation theory codes, and chosen to use the Eulerian PT implementation in `velocileptors` [116] as our default, though the results should be indistinguishable using other codes (see further discussion below).

Each of the theory codes that we employ computes the 1-loop PT predictions for the power spectrum multipoles, including mode-coupling due to quasi-linear evolution, nonlinear bias and redshift-space distortions. The framework includes the aforementioned counterterms that describe the impact of small-scale physics on the observed clustering, the stochastic terms that describe the shot noise and fingers of god in this formalism, and infrared resummation that describes the broadening of the BAO peak due to large-scale flows. The models have been extensively developed and are described in some detail in [70–73] with references therein to the original literature. They are compared to each other, and to a series of simulations, in [1, 75].

A particular advantage of perturbative models of large-scale structure is that they rely on a minimal set of theoretical assumptions to consistently model a wide range of clustering data. They can thus be relied upon for robust inference. A drawback of this approach is that these models tend to require a large number of free parameters. If a signal in the data can be explained by a complex bias model rather than, or in addition to, changes in the underlying cosmology, the models will explore this possibility in the fits. The majority of the cosmological information then originates from scales that are almost linear and protected by fundamental symmetries. Unfortunately, some of the “nuisance parameters” are partially degenerate with cosmological parameters influencing the shape of the linear theory power spectrum (e.g. Ω_m and h). This degeneracy can cause a parameter “projection effect”, where the peak of the marginalised posterior is offset from the global maximum of the posterior (maximum *a posteriori* value, MAP). Of particular concern are degeneracies with non-linear bias parameters, stochastic terms and counterterms that describe the impact of poorly-understood, small-scale physics on the observed clustering. The origin and impact of these effects is discussed in detail in supporting papers to this work [1, 70] as well as in [117] for example. We do not show any results that are subject to significant projection effects, but we illustrate how such effects can occur in appendix A.

In addition to systematics related to theoretical modeling, we quantify several further potential systematic effects using mock catalogues. These mocks are built from the `AbacusSummit` suite of simulations [118, 119] with a galaxy-halo connection prescription based on halo occupation distribution (HOD) models calibrated on the DESI Early Data Release [51] which are described in [120–122]. We have identified and studied seven sources of systematic effects that could bias our cosmological constraints: i) theoretical modeling mentioned above [70], ii) description of the galaxy-halo connection [74], iii) assumptions related to the fiducial cosmology [80], iv) imaging systematics due to inhomogeneities in the target selection [79], v) fibre assignment incompleteness [86, 123, 124], vi) spectroscopic redshift uncertainties and catastrophic redshift errors [125, 126], and vii) covariance matrix estimation [76–78]. Of these, the two most dominant sources of systematic effects are uncertainties associated with the imaging systematics, and the galaxy-halo connection as described by the HOD formalism. To help alleviate imaging systematics we adopt an additional nuisance parameter; more details can be found in [79]. In order to propagate the systematic errors from the HOD to the constraints on cosmological parameters, we estimate the effects on the power spectrum by fitting the model to mocks generated using a variety of plausible galaxy-halo connection models and a single underlying set of cosmology parameters, and quantifying the shifts in the inferred cosmological parameter values from these fits [74]. These systematic contributions are directly added to the statistical power spectrum covariance matrix introduced in section 2.1. Detailed quantification of the systematic error budget for cosmological parameters in the Λ CDM model is presented by [1].

In addition to the power spectrum measurements, we include distance-scale information from the post-reconstruction correlation function in the region around the BAO peak. The DESI DR1 BAO measurements are described in detail in [65], and have already been used in the DR1 BAO cosmological analysis [64].⁷ The joint covariance between the power spectrum and the post-reconstruction BAO measurements is estimated from the 1000 `EZmocks`. The post-reconstruction-BAO part of the resulting covariance matrix is replaced with that estimated from the BAO fits to the data. To this covariance we further add systematic contributions to both the power spectrum and post-reconstruction BAO, as summarised above for the power spectrum, and as detailed in [65] for the BAO measurements. The full pipeline for DESI DR1 modeling analysis and the cosmological constraints in Λ CDM from full-shape alone, and full-shape combined with BAO, are presented in [1].

2.3.2 Likelihood and priors

The combined DESI full-shape and BAO likelihood is implemented using the theoretical modeling summarised in section 2.3.1 and fully described by [1]. The key data input to the likelihood are the measurements of the monopole and quadrupole of the power spectrum, restricted to scales $0.02 < k/(h^{-1}\text{Mpc}) < 0.2$. These measurements are performed for each of the six data samples, with corresponding six redshift bins that are listed in section 2.1.1. In each redshift bin, the clustering measurements are complemented by post-reconstruction BAO parameters [65], and we make use of the complete covariance matrix that covers the

⁷For the analysis in the present paper, we rerun the BAO measurements with the most up-to-date catalogs as described in appendix B of [67].

power spectrum measurements, the post-reconstruction BAO parameters, and their mutual correlation (see section 2.3.1). The measurements in the six redshift bins are considered independent (see section 2.3.2 of [64] for a justification and quantification of inter-bin correlations), and their log-likelihoods are summed to compute the total likelihood. We combine this likelihood with the Ly α BAO likelihood [66], as we did in our DR1 BAO analysis [64].

The procedure that we just described comprises our DESI (FS+BAO) likelihood. When we combine our results with external data from Dark Energy Survey clustering analyses or from the CMB, we adopt the likelihoods provided by these respective collaborations. For type Ia supernova datasets, we assume that the likelihood in the data (distances to individual supernovae) is Gaussian; this assumption has been validated to some extent with simulations (e.g. [127]).

We next describe the non-cosmological “nuisance” parameters in our analysis. To enable the modeling of redshift-space distortions in our likelihood we adopt the Eulerian PT model in `velocileptors` [70]. The Eulerian `velocileptors` redshift-space distortion model produces posteriors that are nearly indistinguishable from those of the Lagrangian PT given DESI DR1 precision, whilst being significantly faster (a single-model evaluation for six redshift bins takes $\simeq 0.5$ s, computing 1-loop terms once and rescaling them to each redshift).

This model also includes the scale-dependent impact of massive neutrinos on the growth rate.⁸ We describe galaxy bias with three (Lagrangian-bias) parameters per redshift bin: b_1 , b_2 , and b_s . The third-order non-local bias parameter is expected to be small and is degenerate with the other nuisance parameters, so we set it to zero; see [70] for tests validating this choice. In practice, we sample and impose priors on $(1 + b_1)\sigma_8$, $b_2\sigma_8^2$, and $b_s\sigma_8^2$ (with each σ_8 evaluated at the effective redshift of the corresponding bin), as the data are sensitive to these combinations. Next, we include stochastic parameters SN_0 and SN_2 which marginalise over small-scale physics (halo exclusion effects, conformity, and virialisation), and enter additively to the anisotropic power spectrum as SN_0 and $SN_2 k^2 \mu^2$ (where μ is the cosine between the wavenumber \mathbf{k} and the line-of-sight to the galaxy pair); see eq. (3.6) in [70]. The priors on the stochastic parameters are Gaussian with mean zero; the widths of these Gaussians are given in table 1, and are further scaled by the estimated shot noise (for SN_0), or by the product of the shot noise with a typical velocity variance and the satellite fraction (for SN_2) [70]. To describe the modeling uncertainties associated with non-linear structure formation, beyond the cutoff scale adopted in one-loop integrals of the perturbation theory model, we include two “counterterm” parameters, α_0 and α_2 corresponding, respectively, to the monopole and the quadrupole. We give each of these parameters a Gaussian prior centered at zero with a standard deviation of 12.5; this prior width is chosen to correspond to the value at which the counterterm corrections represent 50% of the value of the linear power spectrum contribution at the maximum wavenumber $k_{\max} = 0.2 h^{-1}\text{Mpc}$ used in our fits. As a sanity check, we tested the likelihood using the PT code FOLPS [71, 128] as an alternative to `velocileptors`, finding nearly identical results. The bottom part of table 1 summarises the nuisance parameters and their priors. Overall, we adopt three bias, two

⁸https://github.com/sfschen/velocileptors/blob/master/velocileptors/EPT/ept_fullresum_varyDz_nu_fftw.py.

counterterm, and two stochasticity parameters in each of the six redshift bins, resulting in the grand total of 42 nuisance parameters.

Our fiducial constraints, which we refer to as DESI (FS+BAO), are based on this combined power spectrum and post-reconstruction BAO likelihood. Whenever we do *not* add the CMB data to DESI, we include two non-trivial external priors as a default. First, we combine DESI with the external constraint on the physical baryon density $\Omega_b h^2$ that comes from measurements of the primordial deuterium abundance and helium fraction interpreted in the standard model of Big Bang Nucleosynthesis (BBN) [98]; the model predictions were generated using the `PRyMordial` code [99]; see section 2.2 for more details. Second, we add a weak Gaussian prior on the spectral index n_s , centered at the *Planck* mean value $n_s = 0.9649$ and with a width, $\sigma(n_s) = 0.042$, chosen to be 10 times wider than the posterior width obtained from *Planck* temperature, polarisation and lensing spectra [129].⁹ The loose n_s prior to which we refer to as n_{s10} , is therefore implemented as

$$n_s \sim \mathcal{N}(0.9649, 0.042^2) \quad (n_{s10} \text{ prior}), \quad (2.1)$$

where $\mathcal{N}(x, \sigma^2)$ refers to the Gaussian normal distribution with mean x and standard deviation σ . When we combine DESI data with the CMB, we do not apply the BBN likelihood and n_{s10} prior, as the CMB already tightly constrains these two parameters.

2.3.3 Cosmological inference

Our inference procedure largely follows that described in the DESI DR1 BAO paper [64]; the main difference is that we now marginalise over many more nuisance parameters which are specific to the full-shape analysis. We use the cosmological inference code `cobaya` [130, 131], to which we incorporate the PantheonPlus, Union3 and DES-SN5YR SN Ia likelihoods, as well as our DESI likelihood. We use CMB likelihoods based on public packages that are either included in the public `cobaya` version or available directly from the respective teams. Within `cobaya` we use the Boltzmann code `CAMB` [132, 133] to produce model power spectra as a function of parameters. For modified-gravity analyses we employ the code `ISiTGR` [134, 135] which is based on `CAMB` and also called within `cobaya`. When including the CMB data (the combined *Planck*+ACT lensing likelihood) we use higher precision settings as recommended by ACT. We perform Bayesian inference using the Metropolis-Hastings MCMC sampler [136, 137] in `cobaya`, running four chains in parallel. We use `getdist`¹⁰ [138] to derive the constraints presented in this paper. We occasionally wish to calculate $\Delta\chi_{\text{MAP}}^2 \equiv -2\Delta \log \mathcal{L}$, defined to represent the difference (times -2) of log-posteriors $\log \mathcal{L}$ at the maximum posteriori (MAP) parameter-space points. Such MAP points are estimated with the `iminuit` [139, 140] algorithm, as implemented in `cobaya`, starting from the points with maximum log-posterior found in the MCMC chains. More details about the code settings and extraction of the cosmological-parameter values is provided in section 2.5 of the DESI DR1 BAO paper [64].

Table 1 summarises the cosmological parameters that are sampled over in different runs and the priors that are placed on them. We assume a flat geometry (zero curvature)

⁹The loose n_s prior was originally imposed to stabilise the results of the ShapeFit analyses used in our tests [1]. This prior has a small effect on the cosmological results from the Full-Modeling analyses presented in this paper. In cases when we add this prior, we still impose hard prior cutoffs $n_s \in [0.8, 1.2]$.

¹⁰<https://github.com/cmbant/getdist>.

data or model	parameter	default	prior	comment
DESI (ΛCDM)	$\Omega_{\text{cdm}}h^2$	—	$\mathcal{U}[0.001, 0.99]$	—
	$\Omega_{\text{b}}h^2$	—	$\mathcal{U}[0.005, 0.1]$	—
	H_0	—	$\mathcal{U}[20, 100] \text{ km s}^{-1} \text{ Mpc}^{-1}$	—
	$\ln(10^{10} A_{\text{s}})$	—	$\mathcal{U}[1.61, 3.91]$	—
	n_{s}	—	$\mathcal{N}(0.9649, 0.042^2)$	<i>Planck</i> 10σ
with CMB (ΛCDM)	$\Omega_{\text{cdm}}h^2$	—	$\mathcal{U}[0.001, 0.99]$	—
	$\Omega_{\text{b}}h^2$	—	$\mathcal{U}[0.005, 0.1]$	—
	$100\theta_{\text{MC}}$	—	$\mathcal{U}[0.5, 10]$	replaces H_0
	$\ln(10^{10} A_{\text{s}})$	—	$\mathcal{U}[1.61, 3.91]$	—
	n_{s}	—	$\mathcal{U}[0.8, 1.2]$	—
	τ	0.0544	$\mathcal{U}[0.01, 0.8]$	—
Beyond ΛCDM (dynamical DE)	w_0	−1	$\mathcal{U}[-3, 1]$	—
	w_a	0	$\mathcal{U}[-3, 2]$	—
(massive ν)	$\sum m_{\nu}$ (eV)	0.06	$\mathcal{U}[0, 5]$	—
	μ_0	0	$\mathcal{U}[-3, 3]$	—
(modified gravity)	Σ_0	0	$\mathcal{U}[-3, 3]$	—
nuisance (DESI)	$(1 + b_1)\sigma_8$		$\mathcal{U}[0, 3]$	each z -bin
	$b_2\sigma_8^2$		$\mathcal{N}[0, 5^2]$	each z -bin
	$b_s\sigma_8^2$		$\mathcal{N}[0, 5^2]$	each z -bin
	α_0		$\mathcal{N}[0, 12.5^2]$	each z -bin*
	α_2		$\mathcal{N}[0, 12.5^2]$	each z -bin*
	SN_0		$\propto \mathcal{N}[0, 2^2]$	each z -bin*
	SN_2		$\propto \mathcal{N}[0, 5^2]$	each z -bin*

Table 1. Parameters and priors used in our analysis. The first block shows the parameters and their priors in the DESI analyses without the CMB, while the second block shows the slightly extended parameter set used whenever CMB is used either alone or in combination with DESI. The third and fourth blocks respectively show parameters used in various beyond- Λ CDM analyses and nuisance parameters in the DESI analysis. Here, \mathcal{U} refers to a uniform prior in the range given, whilst $\mathcal{N}(x, \sigma^2)$ refers to the Gaussian normal distribution with mean x and standard deviation σ . In addition to the flat priors on w_0 and w_a listed in the table, we also impose the requirement $w_0 + w_a < 0$ in order to enforce a period of high-redshift matter domination. Similarly, an extra prior $\mu_0 < 2\Sigma_0 + 1$ is included for modified-gravity parameters μ_0 and Σ_0 (see section 5.2). Nuisance-parameter combinations listed in the second column are independently varied for each of the six tracer/redshift bins. The asterisk next to the counterterms α_0 and α_2 and stochastic parameters SN_0 and SN_2 indicates that these parameters are marginalised over analytically. The constant of proportionality in front of SN_0 and SN_2 priors indicates that these priors as written are further scaled with corresponding physically motivated terms; see text for details. Note that the n_{s} (*Planck* 10σ) prior is added by default in the DESI (FS+BAO) analysis, but dropped in favor of a flat, non-informative prior once DESI data are combined with the CMB.

throughout. For the basic DESI (FS+BAO) analysis and assuming the Λ CDM model, we vary five cosmological parameters: the Hubble constant H_0 , the physical densities of baryons and cold dark matter $\Omega_b h^2$ and $\Omega_{\text{cdm}} h^2$, and the amplitude and spectral index of the primordial density perturbations, A_s and n_s . When we add the CMB likelihood, instead of H_0 we vary the parameter θ_{MC} which is an approximation to the acoustic angular scale θ_* , and we add the optical depth to reionization parameter, τ . In models beyond Λ CDM we extend this basic cosmological parameter set with additional variables: in the dynamical dark-energy ($w_0 w_a$ CDM) model we have two additional dark-energy parameters w_0 and w_a , in the massive-neutrinos model the sum of the neutrino masses $\sum m_\nu$, and in the class of modified-gravity parametrisation we consider, we introduce additional freedom in the linearly-perturbed Einstein's equations given by parameters μ_0 and Σ_0 (see section 5 for their definitions). Finally, we have a set of nuisance parameters which are required to describe the full-shape clustering signal. In all, our analysis in the base Λ CDM model includes a total of five cosmological parameters and 42 nuisance parameters.

In Λ CDM, we combine the DESI (FS+BAO) analysis with the DESY3 3×2 -pt and 6×2 -pt analyses at the likelihood level, because there is negligible covariance between the multipoles and the projected statistics [141]. We rerun the DESY3 3×2 -pt and 6×2 -pt analyses with the same priors as the DESI (FS+BAO) analysis using the publicly available CosmoSIS [142] pipelines. For each cosmological model, we use the same modeling and scale cuts. For these combinations we then use `CombineHarvesterFlow`¹¹ [143] to fit normalising flows to the DES chains and re-weight the DESI (FS+BAO) chains to compute the joint posteriors. To ensure the results are not affected by undersampling the joint high-density region [143], we randomly split the DES and DESI chains in half and repeat this procedure on both pairs of chains and check the results remain unchanged.

For our modified-gravity inference, we use a DESY3 3×2 -pt likelihood that has been tailored to this specific analysis and validated against the DESY3 modified-gravity results by [107]. This likelihood has been included in our main pipeline using `cobaya`.

3 Dark energy constraints

We focus on two dark energy models: a spatially flat model with a cosmological constant (Λ CDM), and a flat model where the dark energy equation of state is allowed to vary with time and is modeled by two parameters ($w_0 w_a$ CDM). We choose not to study the models with non-zero spatial curvature (Ω_K) or with a constant dark energy equation of state (w CDM) in this paper as we did in the analysis of DESI BAO-only data [64]. Instead, we focus on the two aforementioned models that are of most interest: Λ CDM because it is the standard model of cosmology, and $w_0 w_a$ CDM because of its ability to phenomenologically describe a wide variety of physical models, and because of our earlier findings that show some preference for this model over Λ CDM [64].

3.1 Λ CDM model

We start by constraining the cosmological parameters in a flat Λ CDM model. Here, only a single parameter describes dark energy: $\Omega_{\text{de}} \equiv 1 - \Omega_{\text{m}}$, where Ω_{m} and Ω_{de} are respectively

¹¹<https://github.com/playlor16/CombineHarvesterFlow>.

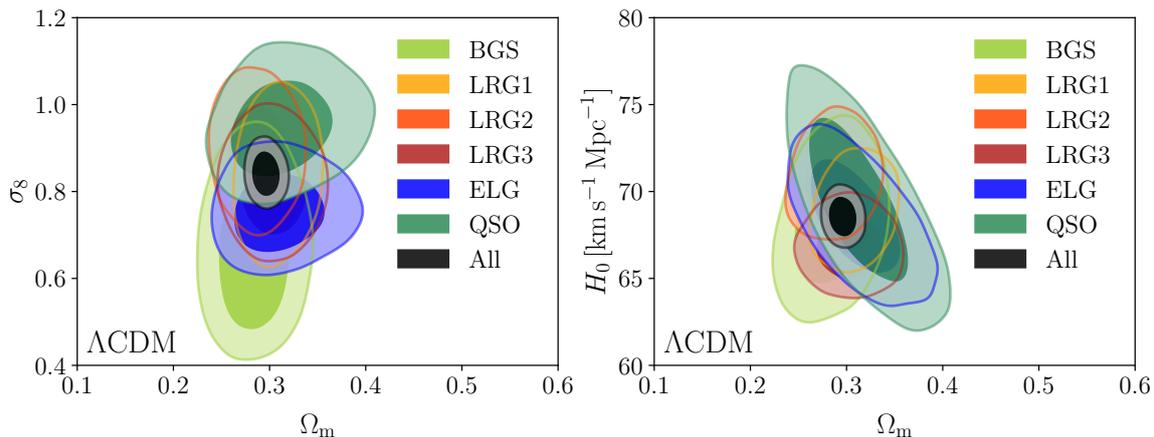


Figure 1. 68% and 95% credible intervals in the Ω_m - σ_8 plane (left panel) and Ω_m - H_0 plane (right panel) from the combined DESI full-shape and BAO analysis, assuming the Λ CDM background. We show the constraints from individual DESI tracers (with the BBN and loose n_s priors), and the combined measurement that includes all the tracers shown and the Ly α BAO data.

the total matter and dark-energy densities relative to critical. The other parameters that we vary, and their respective priors, are listed in table 1. Other than Ω_m , the parameters of most interest to us are the amplitude of mass fluctuations σ_8 (and the combination $S_8 \equiv \sigma_8(\Omega_m/0.3)^{0.5}$), and the Hubble constant H_0 .

Figure 1 shows the DESI-only constraints (combined with the BBN and loose n_s priors) in the Ω_m - σ_8 plane (left panel) and Ω_m - H_0 plane (right panel) from the combined DESI full-shape and BAO analysis, marginalised over all other parameters. The individual contours, with their 68% and 95% credible intervals, show measurements from individual DESI tracers: bright galaxy survey (BGS); luminous red galaxies in three redshift bins (LRG1, LRG2 and LRG3), emission line galaxies (ELG) and quasars (QSO). The combined constraints from these tracers, including their covariance and also including the geometrical (BAO) information from the Ly α forest, is shown by the small black contour in each panel. This figure illustrates the excellent mutual agreement between individual tracers, as well as their complementarity. The combined measurements on these three parameters from all DESI tracers are,

$$\left. \begin{aligned} \Omega_m &= 0.2962 \pm 0.0095, \\ \sigma_8 &= 0.842 \pm 0.034, \\ H_0 &= (68.56 \pm 0.75) \text{ km s}^{-1} \text{ Mpc}^{-1}, \end{aligned} \right\} \text{ DESI (FS+BAO)+BBN+}n_{s10}, \quad (3.1)$$

where, recall, we add the BBN and loose n_s priors to DESI-only data by default. The addition of full-shape information therefore leads to a significant improvement in the constraints on matter density, from $\sigma(\Omega_m) = 0.015$ in the DESI (BAO) + BBN case [64], to $\sigma(\Omega_m) = 0.0095$ when the full-shape information (along with our loose n_s prior) is added. This measurement is almost entirely independent of the constraint from the CMB alone ($\sigma(\Omega_m) = 0.0066$, for our fiducial CMB combination).

Note that the results of eq. (3.1) include a 4% constraint on the amplitude of mass fluctuations, σ_8 , in flat Λ CDM. The BAO alone are a purely geometric probe and thus

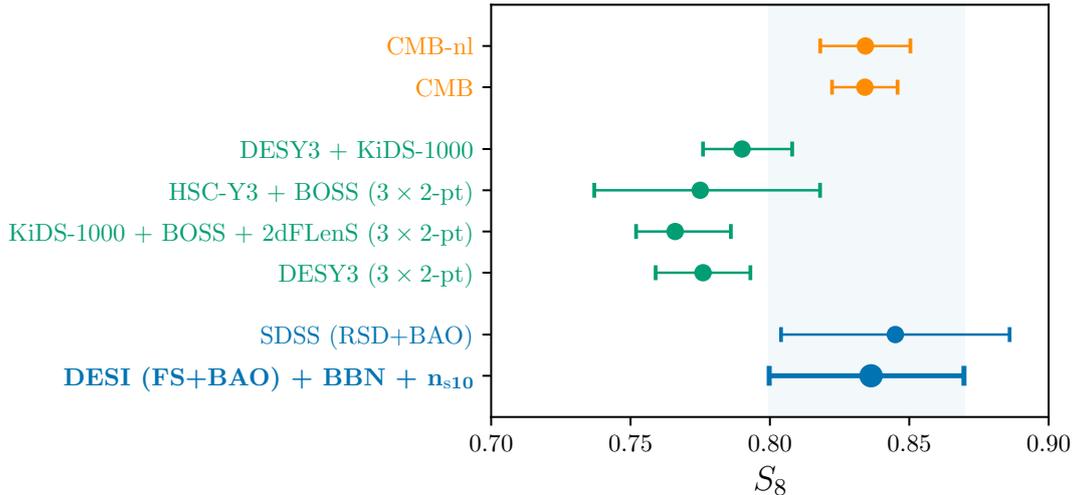


Figure 2. Constraints on the parameter $S_8 \equiv \sigma_8(\Omega_m/0.3)^{0.5}$, assuming the Λ CDM background. The whisker on the bottom (and the corresponding blue-shaded region) shows our fiducial 68% constraint from DESI DR1 (FS+BAO), combined with the BBN and loose n_s priors. The first two whiskers from the top, in orange, show the constraints from the CMB, without and with CMB lensing information. The following four whiskers, in green, show the results from weak lensing probes, while the second to bottom whisker, in blue, shows the constraints from the SDSS combination of redshift-space distortions and BAO. See text for details.

insensitive to σ_8 , so the present constraint comes from using the full-shape clustering measurements of DESI’s tracers. DESI’s constraint on σ_8 is broadly consistent with that from other cosmological measurements, although with an error much larger than that from the CMB data alone (which give $\sigma_8 = 0.8133 \pm 0.0050$).

We next study constraints on the derived parameter $S_8 \equiv \sigma_8(\Omega_m/0.3)^{0.5}$. This parameter combination (unlike σ_8 by itself) is accurately measured by weak gravitational lensing probes of the large-scale structure, and is thus a good meeting point to compare results from galaxy clustering (and CMB and SNIa) surveys with those from weak gravitational lensing. Figure 2 shows the marginalised 68% constraints on S_8 . The bottom whisker in the figure shows our fiducial 68% constraint from DESI DR1 FS+BAO (and the shaded region shows the same)

$$S_8 = 0.836 \pm 0.035 \quad (\text{DESI (FS+BAO)+BBN+}n_{s10}). \quad (3.2)$$

Next, the top two whiskers in figure 2 show constraints from the CMB, without and with CMB lensing information. The following four rows display results from the combined shear analysis from the Dark Energy Survey and Kilo-Degree Survey (DESY3+KiDS-1000; [144]); constraints from shear measured in the Hyper Suprime-Cam Year-3 data combined with galaxy clustering from the Baryon Oscillation Spectroscopic Survey (HSC Y3 + BOSS; [145]); constraints from shear measured in the photometric KiDS-1000 survey combined with spectroscopic galaxy clustering in BOSS and the 2-degree Field Lensing Survey (2dFLenS) (KiDS-1000+BOSS+2dFLenS; [146]); and the combined shear and clustering 3×2 -pt analysis from the DESY3 data [104]. The second-to-bottom whisker shows the constraints from

the SDSS RSD+BAO analysis [27] that used a full-shape analysis that was substantially different in detail to our Full Modeling. We observe an excellent agreement between DESI DR1 (FS+BAO) and the CMB, both of which are slightly higher than the values inferred from the weak lensing surveys. There remains some modest difference between our constraint in eq. (3.2) and that from the lensing probes; for example, the DESY3 measurement is $S_8 = 0.776 \pm 0.017$ [104]. We do not discuss these discrepancies further in this paper. We also note the excellent agreement between our results and the SDSS combined redshift-space-distortion and BAO analysis [27], which found $S_8 = 0.845 \pm 0.041$. Finally, it is interesting to compare our constraints to those resulting from *cross-correlating* a spectroscopically-calibrated LRG sample selected from the DESI Legacy imaging survey [147] with lensing probed by galaxies or the CMB. Cross-correlating with CMB lensing gives $S_8 = 0.775 \pm 0.02$ [148, 149], lower than but consistent with our result (see also [150] for a related analysis that gives similarly consistent constraints). Cross correlating the DESI galaxy positions with DES galaxy shapes gives $S_8 = 0.850^{+0.042}_{-0.050}$ [151], again consistent with our result.

We next discuss the results of combining DESI (FS+BAO) with external probes. When we combine DESI with the CMB information we obtain,

$$\left. \begin{aligned} \Omega_m &= 0.3056 \pm 0.0049, \\ \sigma_8 &= 0.8121 \pm 0.0053, \\ H_0 &= (68.07 \pm 0.38) \text{ km s}^{-1} \text{ Mpc}^{-1}, \end{aligned} \right\} \text{ DESI (FS+BAO) + CMB.} \quad (3.3)$$

These results are fully consistent with those from DESI (FS+BAO) alone, but have much smaller errors: the uncertainties in the matter density and Hubble constant are roughly halved, whilst the uncertainty in σ_8 decreases by more than a factor of five.

We also find that the combined information from DESI full-shape clustering and the BAO, when combined with the CMB, generally improves the constraints from CMB alone: the Ω_m , H_0 , and S_8 errors decrease by 30%, but the σ_8 error increases slightly (by $\sim 10\%$). Our tests on synthetic chains from data with no stochasticity confirm that the observed improvement in Ω_m , H_0 , and S_8 precision, and lack thereof in σ_8 precision, is expected when DESI (FS+BAO) is added to CMB alone. Hence, a mild worsening of the constraints on σ_8 when the DESI data are added is not unexpected.

Combining DESI data with that of the CMB *without* CMB lensing reconstruction (that is, without the ACT+*Planck* lensing likelihood) gives similar results as the full DESI+CMB fits. The only exception is the amplitude of mass fluctuations σ_8 , which now has a $\sim 30\%$ larger error ($\sigma_8 = 0.8086 \pm 0.0071$) in the DESI+CMB-nl analysis, relative to the precision in DESI+CMB. This is to be expected, given that CMB lensing measures the depth of the lensing potential at a range of redshifts ($z \sim 0.5\text{--}3$, very roughly), and is hence sensitive to this parameter. See table 2 for more details.¹²

The final external dataset we consider is the combined analysis of weak gravitational lensing and galaxy angular clustering data from the Dark Energy Survey, DESY3 (3×2 -pt). These lensing and galaxy clustering measurements are sensitive to the power spectrum and growth of structure in a way that is complementary to DESI (FS+BAO), with a different

¹²Note also that we choose not to combine the SN Ia data with that of DESI+CMB in the Λ CDM model, as the two respective measurements disagree in the value of Ω_m (see the discussion in section 4.1 of [64]).

model/dataset	Ω_m	σ_8	S_8	H_0 [km/s/Mpc]	w_0	w_a
Flat ΛCDM						
DESI (FS+BAO) + BBN+ n_{s10}	0.2962 ± 0.0095	0.842 ± 0.034	0.836 ± 0.035	68.56 ± 0.75	—	—
DESI+CMB-nl	0.3045 ± 0.0053	0.8086 ± 0.0071	0.815 ± 0.012	68.14 ± 0.40	—	—
DESI+CMB	0.3056 ± 0.0049	0.8121 ± 0.0053	0.8196 ± 0.0090	68.07 ± 0.38	—	—
DESI+DESY3 (3×2 -pt) + BBN+ n_{s10}	0.2980 ± 0.0070	$0.807^{+0.016}_{-0.020}$	$0.804^{+0.011}_{-0.015}$	$68.67^{+0.69}_{-0.77}$	—	—
DESI+DESY3 (6×2 -pt) + BBN+ n_{s10}	0.2986 ± 0.081	0.799 ± 0.016	0.797 ± 0.011	$68.66^{+0.63}_{-0.73}$	—	—
DESI+DESY3 (6×2 -pt) + CMB-nl	0.3009 ± 0.0034	$0.8028^{+0.0050}_{-0.0045}$	0.8039 ± 0.0056	68.40 ± 0.27	—	—
Flat $w_0 w_a$CDM						
DESI+PantheonPlus + BBN+ n_{s10}	0.3084 ± 0.0089 (0.3117)	0.820 ± 0.035 (0.829)	0.831 ± 0.036 (0.845)	68.4 ± 1.1 (67.8)	-0.875 ± 0.072 (-0.874)	$-0.61^{+0.42}_{-0.36}$ (-0.48)
DESI+Union3 + BBN+ n_{s10}	0.320 ± 0.012 (0.328)	$0.805^{+0.033}_{-0.037}$ (0.809)	0.831 ± 0.036 (0.846)	67.6 ± 1.2 (66.6)	-0.74 ± 0.12 (-0.68)	$-1.12^{+0.58}_{-0.48}$ (-1.15)
DESI+DES-SN5YR + BBN+ n_{s10}	0.3183 ± 0.0090 (0.3214)	$0.808^{+0.033}_{-0.037}$ (0.815)	$0.832^{+0.034}_{-0.038}$ (0.843)	67.7 ± 1.0 (67.2)	-0.761 ± 0.080 (-0.759)	$-1.03^{+0.47}_{-0.40}$ (-0.92)
DESI+CMB +PantheonPlus	0.3061 ± 0.0064 (0.3091)	0.8227 ± 0.0087 (0.8210)	0.8309 ± 0.0091 (0.8334)	68.34 ± 0.67 (67.92)	-0.858 ± 0.061 (-0.847)	$-0.68^{+0.27}_{-0.23}$ (-0.64)
DESI+CMB +Union3	0.3156 ± 0.0090 (0.3246)	0.8152 ± 0.0099 (0.8073)	0.8360 ± 0.0097 (0.8397)	67.35 ± 0.92 (66.44)	-0.742 ± 0.096 (-0.667)	$-1.02^{+0.36}_{-0.32}$ (-1.20)
DESI+CMB +DES-SN5YR	0.3142 ± 0.0063 (0.3171)	0.8163 ± 0.0083 (0.8157)	0.8353 ± 0.0092 (0.8386)	67.48 ± 0.62 (67.11)	-0.761 ± 0.065 (-0.749)	$-0.96^{+0.30}_{-0.26}$ (-0.92)

Table 2. Cosmological parameter results from DESI DR1 (FS+BAO) data alone (labeled “DESI” in the table), and in combination with external datasets. We show results in the baseline flat Λ CDM model and in the (w_0, w_a) parameterisation of the dark energy equation of state. Constraints are quoted as the marginalised means and 68% credible intervals in each case. For flat $w_0 w_a$ CDM model, where mild projection effects are observed in some cases, we also show the best-fit (MAP) value of the parameter in parentheses just below the credible interval. In this and other tables, the shorthand notation “CMB” is used to denote the addition of temperature and polarisation data from *Planck* and CMB lensing data from the combination of *Planck* and ACT, while n_{s10} refers to the loose prior on the spectral index defined in eq. (2.1).

set of systematic errors. We also consider the 6×2 -pt analysis from DESY3, which adds further information from lensing of the CMB (see section 2.2 for details). The addition of the DESY3 (6×2 -pt) data to DESI (FS+BAO) produces,

$$\left. \begin{aligned} \Omega_m &= 0.2986 \pm 0.0081, \\ \sigma_8 &= 0.799 \pm 0.016, \\ H_0 &= (68.66^{+0.63}_{-0.73}) \text{ km s}^{-1} \text{ Mpc}^{-1}. \end{aligned} \right\} \begin{array}{l} \text{DESI (FS+BAO)+BBN+}n_{s10} + \\ \text{DESY3 (6} \times \text{2-pt).} \end{array} \quad (3.4)$$

The addition of DESY3 (6×2 -pt) data therefore improves DESI’s constraint on σ_8 by about a factor of two (compare to eq. (3.1)). It also pulls the central value of σ_8 down by about one standard deviation; the same trend is observed in S_8 , for which the measurement error is improved by nearly a factor of three by the addition of DES data (see table 2). The downward pull in σ_8 and S_8 is unsurprising as the lensing information present in DESY3 data favors lower values of these parameters [104]. We also note that the addition of the DESY3 (6×2 -pt) data does not appreciably change the precision of DESI’s constraints on Ω_m and H_0 . Moreover, the constraints and trends remain similar if we replace the DESY3 (6×2 -pt) by the DESY3 (3×2 -pt) data in these tests (see table 2), but the $\sim 20\%$ improvement in the σ_8 and S_8 errors when going from 3×2 -pt to 6×2 -pt data indicates that the CMB lensing information present in the latter DESY3 analysis is significant. Note finally that the uncertainty in Ω_m *increases* by 20% as one goes from 3×2 -pt to 6×2 -pt data; this is because CMB lensing (present in 6×2 -pt) brings the matter density parameter down, and thus eases the mild tension between the higher value preferred by the 3×2 -pt data and the lower value preferred by DESI that led to a correspondingly tight error bar on Ω_m . Figure 3 further illustrates these results, and compares the DESI+DESY3 (6×2 -pt) constraints to those obtained from the CMB temperature and polarisation (without CMB lensing¹³).

It is also interesting to compare the combination of DESI and DESY3 (6×2 -pt) measurements in eq. (3.4) to those from DESY3 (6×2 -pt) alone [108]. The addition of DESI full-shape and BAO data improves the constraints from DESY3 (6×2 -pt) alone by about a factor of two in Ω_m and σ_8 , though only by about 10% in S_8 , the parameter to which gravitational-lensing surveys are sensitive.

We also investigate adding CMB (with no lensing reconstruction) data to the combination of DESI full-shape and BAO along with DESY3 (6×2 -pt). We find

$$\left. \begin{aligned} \Omega_m &= 0.3009 \pm 0.0034, \\ \sigma_8 &= 0.8028^{+0.0050}_{-0.0045}, \\ H_0 &= (68.40 \pm 0.27) \text{ km s}^{-1} \text{ Mpc}^{-1}. \end{aligned} \right\} \begin{array}{l} \text{DESI (FS+BAO) + CMB-nl +} \\ \text{DESY3 (6} \times \text{2-pt).} \end{array} \quad (3.5)$$

We see that the addition of CMB-nl to the combination of DESI and DESY3 (6×2 -pt) leads to about a factor of three improvement in the measurement error of cosmological parameters. Specifically, the matter density is determined to 1% accuracy, σ_8 and S_8 are

¹³Recall that when we combine DESY3 (6×2 -pt) data with the CMB, we use CMB without its lensing reconstruction (hence “CMB-nl”) in order not to double-count information with the CMB lensing present in the 6×2 -pt analysis; see section 2.2.

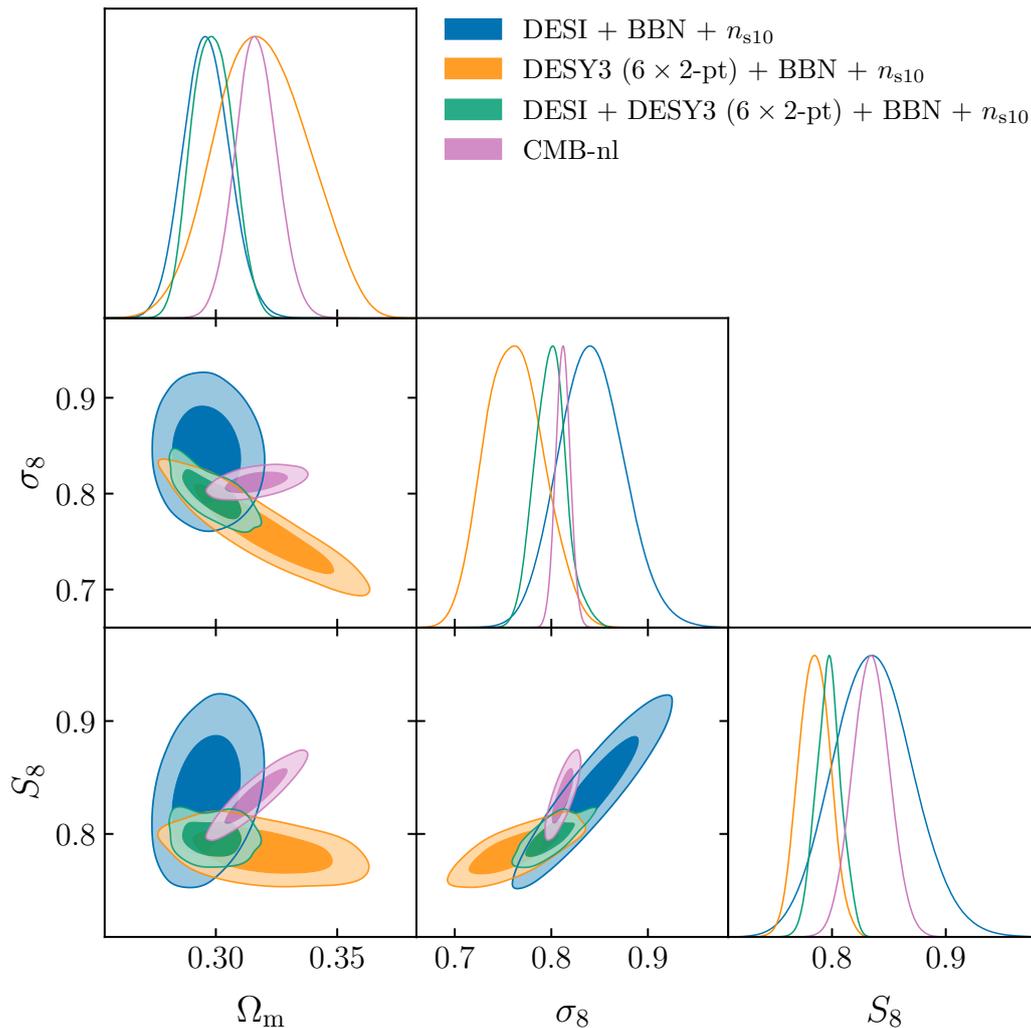


Figure 3. Projected constraints on Ω_m , σ_8 , and S_8 in the Λ CDM model, with 68% and 95% credible intervals shown in each case. The blue contours display the DESI (FS+BAO) constraints (with the BBN and loose n_s priors). The orange contours show constraints from the DESY3 (6×2 -pt) analysis which combines galaxy clustering, cosmic shear and CMB lensing. The green contours show the combination of DESI and DESY3 (6×2 -pt) data. For comparison we also show the CMB temperature and polarisation constraints without the lensing reconstruction as purple contours.

determined to 0.6%, and the Hubble constant is pinned down to 0.4%. These are the strongest constraints in the Λ CDM model presented in this paper, and they show the remarkable power of modern survey data to measure key parameters of the cosmological model. We also find that the addition of DESI data improves the measurement of the Hubble constant by $\sim 20\%$ compared to CMB-nl+DESY3 (6×2 -pt) data alone. Our constraints are further illustrated in the Ω_m - σ_8 plane in figure 4.

Our baseline measurement of the Hubble constant in eq. (3.1) indicates that DESI alone (helped with the BBN and n_{s10} priors) prefers lower values of H_0 , in agreement with independent measurements by the CMB [129]. When DESI and the CMB (no lensing) are

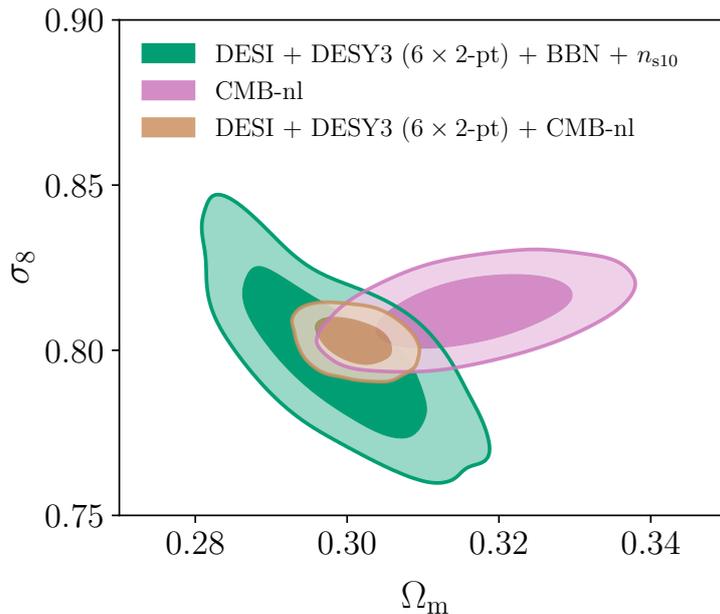


Figure 4. Constraints projected to the Ω_m - σ_8 plane in the Λ CDM model. The green contour shows constraints from the DESI full-shape and BAO analysis, combined with the BBN and n_{s10} priors, and further complemented with the DESY3 (6×2 -pt) data. The pink contour shows the CMB without lensing reconstruction. The brown contour shows the combination of the two, that is, DESI combined with DESY3 (6×2 -pt) and CMB-nl.

combined together and further helped with the DESY3 (6×2 -pt) data, the central value of H_0 does not appreciably change but the errors decrease by about a factor of three relative to DESI alone; the resulting measurement of the Hubble constant (eq. (3.5)) is in a 4.5σ tension with the much higher value preferred by the distance-ladder measurements that use Cepheid variables and nearby SN Ia [152]. We will study cosmological-parameter tensions in more detail in a dedicated supporting paper [153].

Note finally that the Hubble constant constraints determined in the FS+BAO analysis, and its combinations with external probes, depend on sound horizon physics. It will be of interest to compare our results to analyses that use sound-horizon-independent methods, and either marginalise over the sound horizon [154, 155], use only energy densities [156], or use power spectrum or correlation function features that depend on the epoch of matter-radiation equality [157–159].

3.2 w_0w_a CDM model

The combination of DESI DR1 baryon acoustic oscillations with cosmic microwave background and type Ia supernova datasets demonstrated a preference for a time-varying dark energy equation of state [64]. Here we report how these w_0 - w_a results are updated once we add full-shape information to the DESI BAO data. Moreover, we investigate whether generalising the cosmological model from Λ CDM to w_0w_a CDM loosens the constraints on the other parameters of interest, such as σ_8 and (as we will study in section 4) the sum of the neutrino masses.

We study the time-varying dark energy equation of state in the parameterisation [160, 161]

$$w(a) = w_0 + w_a(1 - a), \quad (3.6)$$

where w_0 and w_a are the two beyond- Λ CDM parameters describing the temporal evolution of the dark-energy equation of state, and a is the scale factor. We make use of the parametrised post-Friedmann approach [162] to compute the dark energy perturbations when calculating the CMB angular power spectrum. We do not show the measurements derived from DESI (FS+BAO) data alone, as they are significantly affected by parameter projection effects (see [1] and appendix A). The reason that the projection effects are much more pronounced in the FS+BAO constraints in w_0w_a CDM than in the equivalent BAO-alone analysis [64] is the presence of many additional nuisance parameters in the full-shape analysis which allow additional freedom and open new degeneracy directions. Therefore, we only consider DESI full-shape clustering and BAO in combination with other data when testing the w_0w_a CDM model. We find that the combination of DESI and CMB is also subject to strong projection effects, but that further addition of data from type Ia supernovae removes them (see again appendix A).

Therefore, when allowing the extra freedom in the expansion history allowed by the w_0w_a CDM model, we only present results for the combination of our DESI data with the CMB as well as the various type Ia supernova datasets. With PantheonPlus, we find

$$\left. \begin{aligned} w_0 &= -0.858 \pm 0.061, \\ w_a &= -0.68_{-0.23}^{+0.27}, \end{aligned} \right\} \begin{array}{l} \text{DESI (FS+BAO)+CMB} \\ \text{+PantheonPlus.} \end{array} \quad (3.7)$$

The combination with Union3 supernova data results in

$$\left. \begin{aligned} w_0 &= -0.742 \pm 0.096, \\ w_a &= -1.02_{-0.32}^{+0.36}, \end{aligned} \right\} \begin{array}{l} \text{DESI (FS+BAO)+CMB} \\ \text{+Union3.} \end{array} \quad (3.8)$$

Finally, the combination with DES-SN5YR gives

$$\left. \begin{aligned} w_0 &= -0.761 \pm 0.065, \\ w_a &= -0.96_{-0.26}^{+0.30}, \end{aligned} \right\} \begin{array}{l} \text{DESI (FS+BAO)+CMB} \\ \text{+DES-SN5YR.} \end{array} \quad (3.9)$$

These results are summarised in figure 5, which shows that the outcomes of the combined analyses that include DESI with BAO-only information are consistent with those that include both the full-shape and the BAO information. Moreover, when the DESI (BAO)+CMB+SN Ia combination is supplemented with the full-shape information from DESI, the constraints on the dark energy equation of state parameters improve: the dark energy figure of Merit [163]¹⁴ for combinations involving PantheonPlus, Union3, and DES-SN5YR increases by a factor of 1.16, 1.23 and 1.15 respectively. Thus, the respective credible-region areas in the w_0 - w_a plane are reduced by about 20% when the full-shape data is added to the BAO.

¹⁴The Dark Energy Task Force Figure of Merit (FoM) is defined as the inverse area of the 95% posterior contour in the w_0 - w_a plane. For a Gaussian posterior, $\text{FoM} \propto |\det \mathbf{C}|^{-1/2}$, where \mathbf{C} is the projected 2×2 covariance in the (w_0, w_a) subspace, and this is the definition we use to calculate the ratios of FoMs.

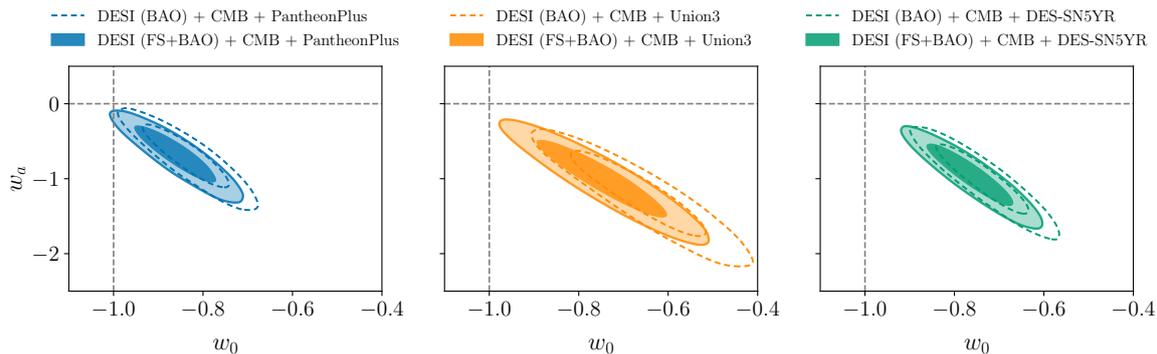


Figure 5. Constraints on w_0 and w_a , assuming a w_0w_a CDM model with a time-varying dark energy equation of state parameterisation (eq. (3.6)). The contours represent the 68% and 95% credible intervals. The solid blue, orange, and green contours represent the combination of DESI (FS+BAO) and CMB with three respective SN Ia data sets: PantheonPlus, Union3 and DES-SN5YR. The dashed blue, orange, and green contours show the same respective combinations, but with the DESI full-shape clustering and BAO replaced by DESI (BAO). The figure shows how the addition of the full-shape information to the BAO-only data improves the precision of the constraints. The measurements of these two parameters remain mutually consistent, and prefer $w_0 > -1$ and $w_a < 0$.

Figure 5 also illustrates that, when combining our full-shape and BAO results with that from the cosmic microwave background and type Ia supernovae, the results remain fully consistent with the same probe combination that contains DESI BAO-only [64], and continue to indicate a preference for a departure from the Λ CDM values of ($w_0 = -1, w_a = 0$). Calculating the difference between the maximum *a posteriori* value of the w_0w_a CDM models in our chains, and the MAP value of the models that enforce Λ CDM ($w_0 = -1, w_a = 0$), we find values of $\Delta\chi_{\text{MAP}}^2 = -8.8, -14.5$ and -17.5 for the combinations of DESI and CMB with PantheonPlus, Union3 and DES-SN5YR respectively. These values of $\Delta\chi_{\text{MAP}}^2$ correspond to preferences for w_0w_a CDM over Λ CDM at the significance levels of 2.5σ (PantheonPlus), 3.4σ (Union3), and 3.8σ (DES-SN5YR). These preferences are similar to the ones we found for the combination of DESI BAO-only data with CMB and SN Ia [64]. Note also that the change in the preference for departures from Λ CDM in the combination of DESI and CMB with Union3 data, when going from BAO to full-shape plus BAO, is smaller than what figure 5 visually implies because of the noticeable projection effects in this combination (compare the mean and MAP values of w_0 and w_a for this combination of probes in table 2).

Figure 6 illustrates the dependence of the measurements of Ω_m and σ_8 on the cosmological background assumed. We show the DESI full-shape plus BAO constraints in the Ω_m - σ_8 plane assuming Λ CDM background, and also constraints in the w_0w_a CDM model for DESI in combination with each of the three SN Ia samples. Overall, the measurements of Ω_m and σ_8 are consistent for all four cases shown. In more detail, the contours for the w_0w_a CDM model shift slightly toward higher Ω_m and lower σ_8 values compared to the DESI-only Λ CDM results presented in eq. (3.1). The numerical results for these measurements are presented in table 2. Note in particular that the values of S_8 in the w_0w_a CDM background from the combination of DESI and SN Ia datasets remain fully consistent with its value in Λ CDM from DESI alone, as the changes in σ_8 and Ω_m (that enter in the definition of S_8) effectively cancel out.

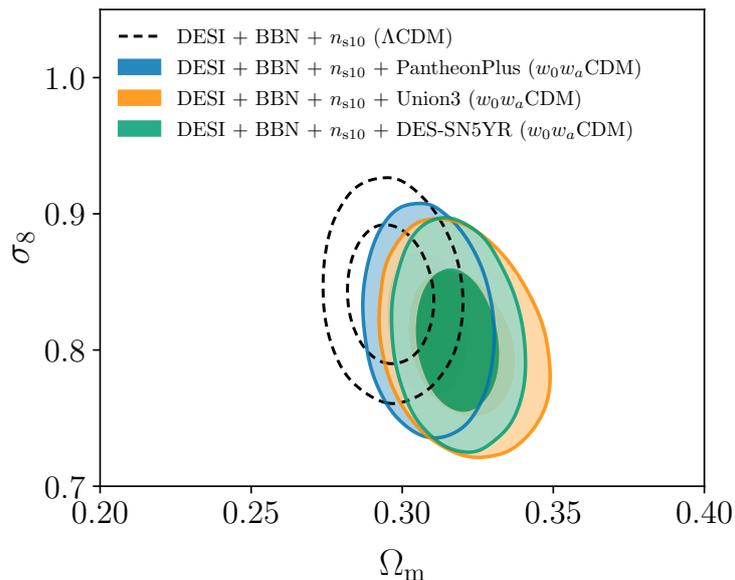


Figure 6. Constraints in the Ω_m - σ_8 plane illustrating the dependence on the assumed cosmological background. The dashed black contour shows the result of analysing DESI full-shape clustering and BAO in Λ CDM model. The blue, orange and green contours represents the respective combinations of DESI with PantheonPlus, Union3 and DES Year-5 supernova datasets, all assuming the w_0w_a CDM model. In all cases, we include the usual BBN and n_{s10} priors. The colored contours highlight a modest shift to smaller σ_8 values as the cosmological background changes from Λ CDM to w_0w_a CDM. Overall, however, the constraints on both parameters remain consistent between all cases.

Figure 7 shows a more detailed scan of the parameter space in our w_0w_a CDM analysis, showing the projection onto the five-dimensional subspace spanned by $\Omega_m, H_0, \sigma_8, w_0$ and w_a . We show results for the combination of DESI (FS+BAO) and CMB with each of the three SN Ia datasets. The numerical constraints are presented in table 2. When the cosmological background changes from Λ CDM to w_0w_a CDM, the central values in $\Omega_m, \sigma_8/S_8$ and H_0 remain unchanged within the errors. The error bars in these parameters increase but remain small, with percent-level precision in each for all three DESI+CMB+SN Ia combinations. This shows the robustness of our constraints on these key parameters to variations in the underlying cosmological model.

In summary, DESI full-shape clustering and BAO data, in combination with the CMB and SN Ia, continue to show hints of a departure from the Λ CDM model. The degree of preference for this departure from the standard cosmological model depends on the choice of the SN Ia dataset. This preference has already been investigated in the recent full-shape re-analysis of BOSS data that combines it with similar external data as our study [164], and it will be interesting to investigate in detail how the choice of dataset and full-shape methodology affects the results. Studying models with more freedom in the dark-energy sector than allowed in w_0w_a CDM is also promising [165, 166], especially as the data get better. Looking ahead, DESI’s forthcoming Year-3 data analysis will shed significant new light on dynamical dark energy.

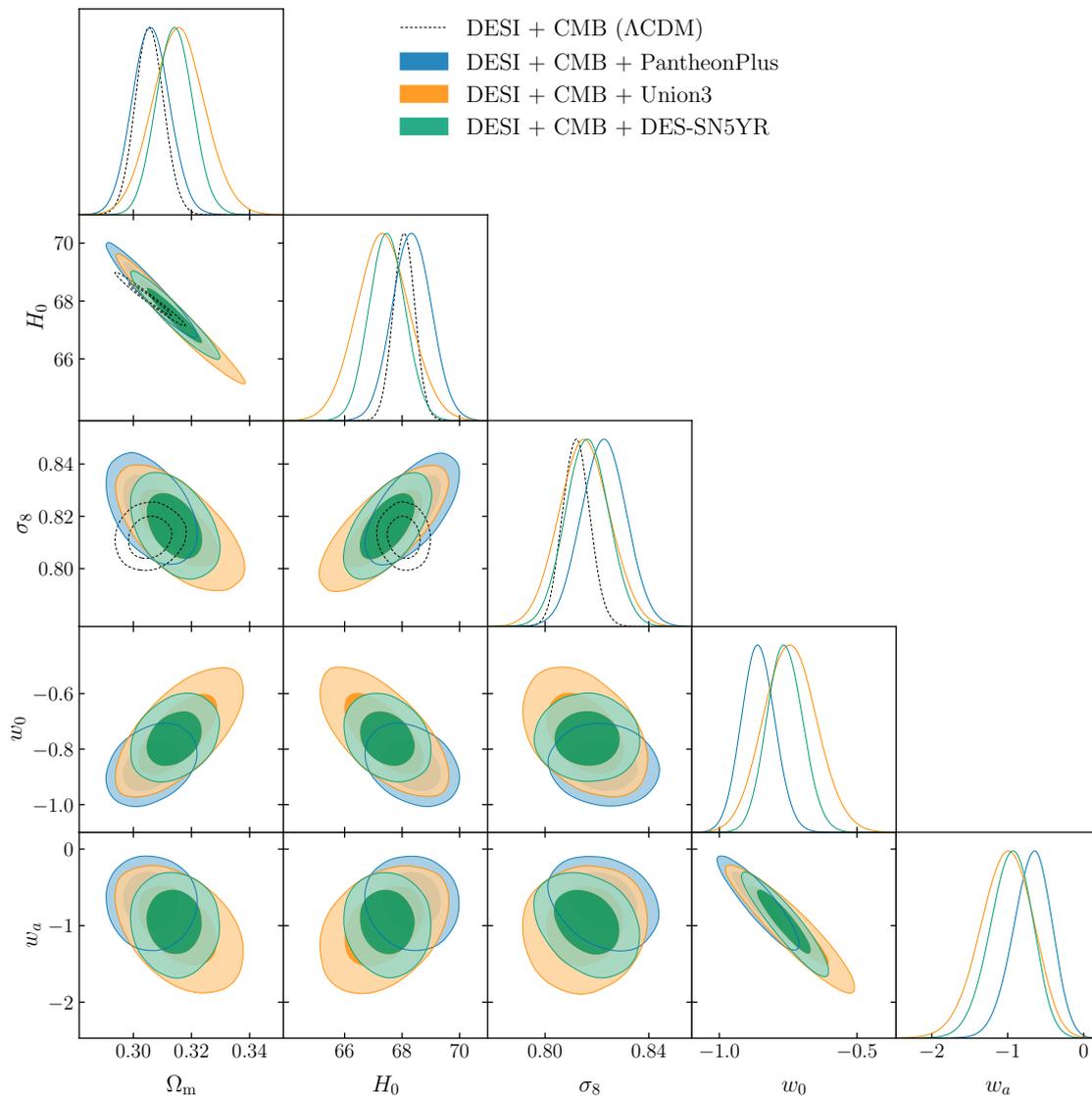


Figure 7. Constraints in the w_0w_a CDM model, projected onto the five-dimensional subspace spanned by Ω_m , H_0 , σ_8 , w_0 and w_a . We show the combination of DESI full-shape and BAO and the CMB with each of the three SN Ia datasets: PantheonPlus (blue), Union3 (orange) and DES-SN5YR (green). For comparison, the black dotted contours illustrate the constraints obtained from DESI and the CMB within w_0w_a CDM.

4 Neutrino constraints

In this section we exploit the DESI full-shape data combined with other datasets in order to set limits on the sum of neutrino masses and on the number of relativistic species in the early universe.

4.1 Sum of neutrino masses

The existence of massive neutrinos, implied by the discovery of neutrino oscillations [167–169], is direct evidence for physics beyond the Standard Model. Major efforts are underway to

constrain neutrino properties in laboratory experiments, but neither the ordering of the neutrino masses nor their absolute scale is known. Oscillation experiments are sensitive to the mass squared differences [170–173] and set a lower bound on the sum of the three neutrino masses. If the smallest mass splitting is between the lowest mass eigenstates, neutrino masses are said to have the normal ordering (NO) and satisfy $\sum m_\nu \geq 0.059$ eV. The other possibility, known as the inverted ordering (IO), implies $\sum m_\nu \geq 0.10$ eV. The strongest model-independent upper bound on the absolute mass scale comes from the KATRIN experiment [174], which constrains the effective electron anti-neutrino mass to $m_\beta < 0.45$ eV (90% CL). Assuming three quasi-degenerate neutrinos, this is equivalent to $\sum m_\nu < 1.35$ eV (90% CL), which is an order of magnitude higher than typical upper limits from cosmology.

Cosmological probes are sensitive to a number of distinctive signatures of cosmic background neutrinos, which enable independent and complementary constraints on $\sum m_\nu$ [175–177] and can be broadly separated into effects on the background expansion and on the growth of fluctuations. At the background level, massive neutrinos affect the expansion history in a unique way, contributing as radiation in the early universe and as non-relativistic matter at recent epochs. Neutrino masses can be tightly constrained from this signature alone. Massive neutrinos also have a strong effect on the growth of cosmic structure. After becoming non-relativistic, neutrinos retain large thermal velocities and cannot be contained in regions smaller than a typical free-streaming length, which is inversely proportional to their mass. Since, at this stage, neutrinos contribute fully to the expansion as non-relativistic matter, but only partially to the clustering, the growth of density perturbations is reduced on small scales. This is manifested as a scale-dependent suppression of the matter power spectrum, which scales as $\Delta P(k)/P(k) \propto -\Omega_\nu/\Omega_m$ [175, 178, 179] and affects equally the broadband shape of the power spectrum and the amplitude of the BAO, with constraints from the latter effect being potentially less prone to parameter projection effects [41].

The DESI full-shape power spectrum analysis allows the sum of neutrino masses to be constrained independently of the CMB. Assuming a Λ CDM background and three degenerate neutrino species, we find an upper bound

$$\sum m_\nu < 0.409 \text{ eV} \quad (95\%, \text{ DESI (FS+BAO)+BBN}+n_{s10}). \quad (4.1)$$

We reiterate that we adopt two external priors when DESI is not combined with CMB data: an external BBN prior on $\Omega_b h^2$ and a weak prior on the spectral index, n_s , that corresponds to ten times the uncertainty (10σ) from *Planck*. On the scales measured by DESI, the broadband suppression from neutrinos is degenerate with n_s . Moreover, n_s is also degenerate with H_0 and with Ω_m . Adding a stronger (1σ) Gaussian prior on n_s , representing a limited use of CMB information, improves the upper bound to

$$\sum m_\nu < 0.300 \text{ eV} \quad (95\%, \text{ DESI (FS+BAO)+BBN+tight } n_s \text{ prior}). \quad (4.2)$$

This is similar to the constraint from the CMB without CMB lensing, $\sum m_\nu < 0.265$ eV (95%), as can be seen from the top left panel of figure 8, but relies on the growth of fluctuations instead of primarily on effects on the background expansion (see [180] for an analysis of the

contributions of geometric and growth information from the CMB). The physics behind these constraints will be addressed in greater detail in [181]. The combination of CMB temperature and polarisation data and CMB lensing, as before simply denoted as “CMB”, yields an upper bound of $\sum m_\nu < 0.218 \text{ eV}$ (95%). The tightest limits are obtained from the full combination of DESI and CMB measurements, demonstrating their strong complementarity. First, using CMB and DESI BAO only yields

$$\sum m_\nu < 0.082 \text{ eV} \quad (95\%, \text{ DESI (BAO only)+CMB}). \quad (4.3)$$

Compared to the equivalent result for the same data combination reported in [64], this figure is slightly higher due to a change in the external CMB lensing likelihood as we switched from version v1.1 of the ACT lensing likelihood to version v1.2. Nevertheless, it represents the strongest cosmological bound from CMB and BAO information only. Adding the full-shape information improves the upper bound further to

$$\sum m_\nu < 0.071 \text{ eV} \quad (95\%, \text{ DESI (FS+BAO)+CMB}). \quad (4.4)$$

The tight limit in eq. (4.3) from DESI BAO + CMB arises from the preference of DESI data for high values of H_0 and low values of Ω_m , which suppresses the bounds on $\sum m_\nu$ due to the geometric degeneracy between these parameters. This trend is further reinforced by the DESI full-shape analysis, which improves the precision in H_0 and especially Ω_m once combined with DESI BAO. Hence, the improvement in the constraint seen in eq. (4.4) is not directly associated with the suppression of the power spectrum, but with a greater pull towards low Ω_m and high H_0 . This is illustrated in the top right panel of figure 8, which shows how the degeneracy in the H_0 - $\sum m_\nu$ plane is broken by the DESI + CMB combination.

The marginalised posterior distribution of $\sum m_\nu$ is shown in the bottom left panel of figure 8. As was the case for the DESI BAO analysis [64], the posterior peaks at a value near $\sum m_\nu = 0 \text{ eV}$, which is excluded by neutrino oscillations. This is true for all data combinations reported above, including the CMB-independent result in eq. (4.1). Although similar behavior had already been seen in *Planck* and SDSS data [e.g. 22, 27, 41, 182], the results have always been compatible with the oscillation constraints.

Since the release of the DESI BAO results, a number of recent studies have identified stronger tensions between the constraints from cosmology and neutrino oscillations, either by combining the DESI BAO measurements with additional data sets [183, 184] or by considering the possibility of neutrino masses beyond the experimental limits, including apparent negative values, as an indicator of systematics, new neutrino properties or a non-standard cosmological expansion [185–187]. Compared to the DESI (BAO) analysis, the full-shape information leads to a 15% stronger constraint and a slight increase in the tension. Nevertheless, our baseline DESI + CMB result remains compatible with the lower bound for the normal ordering at the $\sim 2\sigma$ level. Results from the full-shape analysis for extended neutrino models will be presented in [181].

The bounds given above were all obtained under the assumption of three degenerate neutrino species with a prior that $\sum m_\nu > 0$. The 95% upper limits would increase considerably

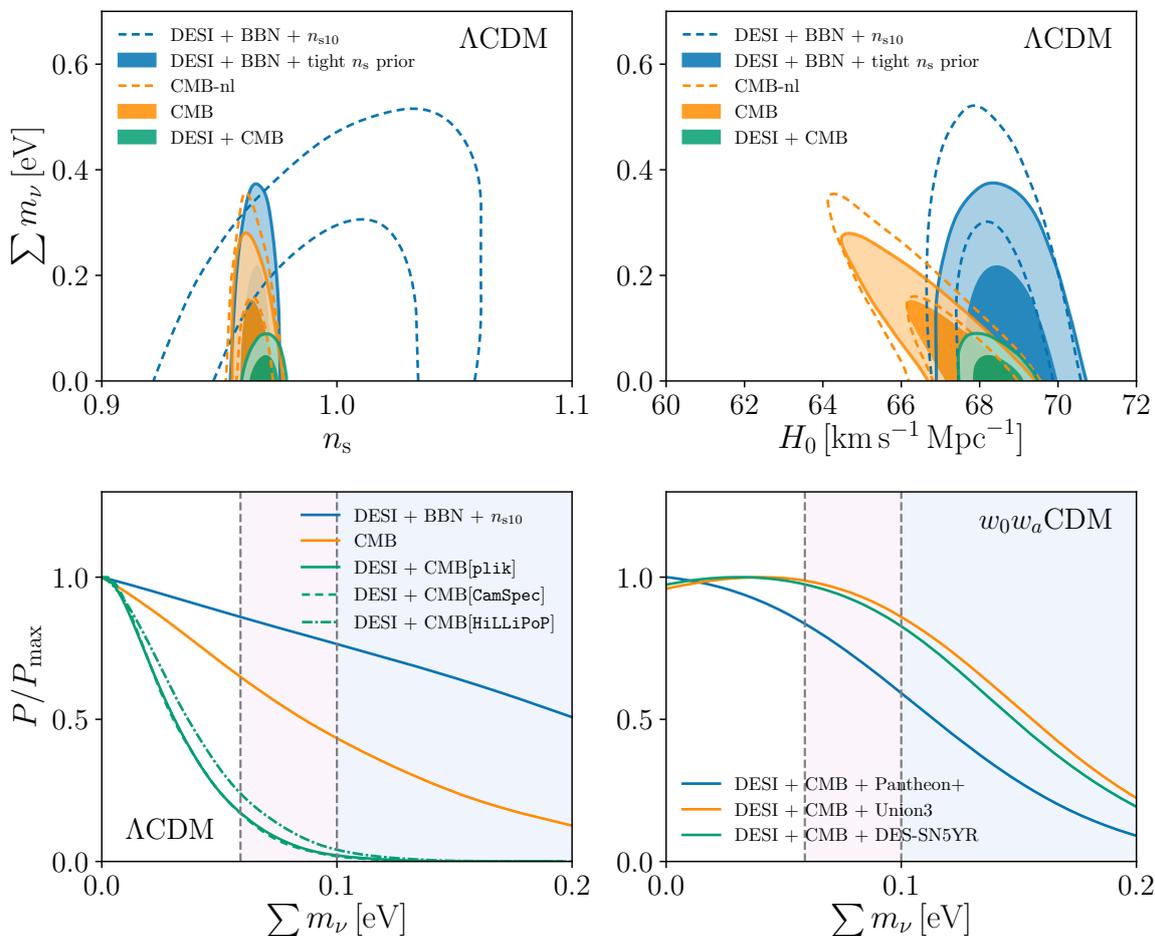


Figure 8. *Top left panel:* constraints in the n_s - $\sum m_\nu$ plane. The blue dashed contours show the 68% and 95% credible intervals for the fiducial DESI (FS+BAO) dataset, accompanied, as usual, by the BBN prior on $\Omega_b h^2$ and a loose prior on n_s . The filled blue contours illustrate the improvement when the n_s prior is tightened to be that from *Planck* (rather than 10 times weaker). The dashed orange contours show the results from CMB without the lensing reconstruction, while the filled orange contours show the constraints from CMB with lensing. Finally, the green contours show the DESI+CMB combination. *Top right panel:* constraints in the H_0 - $\sum m_\nu$ plane for the same data combinations as in the top left panel, illustrating that the DESI+CMB combination breaks the geometric degeneracy between H_0 and $\sum m_\nu$. *Bottom left panel:* one-dimensional posteriors on the sum of the neutrino masses. We show constraints from DESI (FS+BAO) alone, CMB alone, and DESI+CMB for three alternative choices of the CMB likelihood. The minimal masses for the normal or inverted mass ordering scenarios, corresponding respectively to $\sum m_\nu \geq 0.059$ eV and $\sum m_\nu \geq 0.10$ eV, are shown by the vertical dashed lines and the shaded regions. *Bottom right panel:* same as for the bottom left panel, but for the $w_0 w_a$ CDM background and showing constraints from the combination of DESI, CMB, and SN Ia as labelled.

if more restrictive priors, motivated by neutrino oscillations ($\sum m_\nu \geq 0.059$ eV or $\sum m_\nu \geq 0.10$ eV depending on the ordering), were imposed [64]. This is a consequence of the fact that much of the posterior volume is in the unphysical range ($\sum m_\nu < 0.059$ eV). The preference for the normal ordering over the inverted ordering from DESI + CMB stands at a modest $\Delta\chi_{\text{MAP}}^2 \simeq -2$ level. The implications for the neutrino mass ordering of a combined analysis of cosmological and laboratory data will be discussed in detail in a forthcoming publication [181].

A potential systematic affecting neutrino mass bounds is the so-called CMB lensing anomaly [185, 186, 188–190], which refers to a small oscillatory feature in the *Planck* data that could be explained by additional gravitational lensing of the CMB [129, 191–193]. Two independent analyses of the latest *Planck* PR4 data release, the high- ℓ **CamSpec** likelihood [90, 91] and the low- ℓ **LoLLiPoP** and high- ℓ **HiLLiPoP** likelihoods [92, 93], are less affected by this anomaly. To investigate the robustness of our constraints, we repeat the analysis but replace the high- ℓ PR3 **plik** (TTTEEE) likelihood with PR4 **CamSpec**. This combination yields a bound that is nearly identical to our baseline result in eq. (4.4),

$$\sum m_\nu < 0.069 \text{ eV} \quad (95\%, \text{ DESI (FS+BAO) + CMB[CamSpec]}). \quad (4.5)$$

Replacing both the low- ℓ **simall** (EE) likelihood with **LoLLiPoP** and the high- ℓ **plik** likelihood with **HiLLiPoP** has a somewhat larger effect, relaxing the bound to

$$\sum m_\nu < 0.081 \text{ eV} \quad (95\%, \text{ DESI (FS+BAO) + CMB[HiLLiPoP]}). \quad (4.6)$$

The associated posteriors are shown in the bottom left panel of figure 8. Overall, the posteriors are fairly consistent, but the likelihoods that are least affected by the lensing anomaly (**LoLLiPoP** and **HiLLiPoP**) yield a slightly greater upper bound.

It is important to note that constraints on $\sum m_\nu$ depend strongly on the assumed dark energy model. Due to parameter degeneracies, allowing the dark energy equation of state, w , to vary can considerably relax the upper bound on $\sum m_\nu$ [194], although this depends on the conditions imposed on w [183, 195, 196]. In this section, as in section 3, we restrict attention to the w_0w_a CDM model in which the dark energy equation of state is a function of time described by two parameters. For the combination with DES-SN5YR in w_0w_a CDM, the upper bound on the neutrino mass relaxes to

$$\left. \begin{aligned} \sum m_\nu < 0.196 \text{ eV} \quad (95\%), \\ w_0 = -0.753 \pm 0.070, \\ w_a = -1.02^{+0.37}_{-0.29}, \end{aligned} \right\} \begin{array}{l} \text{DESI (FS+BAO) + CMB +} \\ \text{DES-SN5YR,} \end{array} \quad (4.7)$$

while the constraints on w_0 and w_a do not differ significantly from those obtained for fixed $\sum m_\nu$. The bottom right panel of figure 8 shows the marginalised posterior distribution of $\sum m_\nu$ in the w_0w_a CDM model for DESI + CMB combined with the three supernova datasets: PantheonPlus, Union3, and DES-SN5YR. In the case of Union3 and DES-SN5YR, the maximum of the posterior is recovered near the physical mass range. However, in all cases, the data can accommodate larger neutrino masses, alleviating the tension with neutrino

model / dataset	Ω_m	H_0 [km s ⁻¹ Mpc ⁻¹]	$\sum m_\nu$ [eV]	N_{eff}
ΛCDM + $\sum m_\nu$				
DESI (FS+BAO)+BBN+ n_{s10}	$0.2991^{+0.0098}_{-0.011}$	68.40 ± 0.78	< 0.409	—
DESI+CMB	0.3026 ± 0.0052	68.35 ± 0.41	< 0.071	—
ΛCDM + N_{eff}				
DESI+CMB	0.3028 ± 0.0059	68.9 ± 1.1	—	3.18 ± 0.16
$w_0 w_a$CDM + $\sum m_\nu$				
DESI+CMB+PantheonPlus	0.3064 ± 0.0067	68.33 ± 0.68	< 0.175	—
DESI+CMB+Union3	0.3167 ± 0.0095	67.30 ± 0.93	< 0.201	—
DESI+CMB+DES-Y5SN	0.3151 ± 0.0067	67.45 ± 0.63	< 0.196	—
$w_0 w_a$CDM + N_{eff}				
DESI+CMB+PantheonPlus	0.3068 ± 0.0066	68.0 ± 1.1	—	2.97 ± 0.17
DESI+CMB+Union3	0.3167 ± 0.0093	66.8 ± 1.2	—	2.94 ± 0.17
DESI+CMB+DES-Y5SN	0.3152 ± 0.0065	67.0 ± 1.0	—	2.94 ± 0.17

Table 3. Cosmological parameter estimates and constraints from DESI DR1 full-shape clustering and BAO data, in combination with external datasets, when considering extensions in the neutrino sector of the Λ CDM and $w_0 w_a$ CDM models (“DESI” in the table stands for DESI DR1 (FS + BAO)). Results with two-sided error bars refer to the marginalised means and 68% credible intervals; upper bounds on $\sum m_\nu$ refer to 95% limits. All constraints on $\sum m_\nu$ assume a model with three degenerate mass eigenstates and a minimal prior, $\sum m_\nu > 0$ eV. The empty $\sum m_\nu$ and N_{eff} fields indicate that fixed values of $\sum m_\nu = 0.06$ eV and $N_{\text{eff}} = 3.044$ respectively were adopted.

oscillations [186, 189]. These results are in line with those of the DESI (BAO) analysis [64]. The constraints for the combinations with Union3 and PantheonPlus are given in table 3.

4.2 Number of relativistic species

We also report constraints on the effective number of relativistic species, N_{eff} . This parameter is defined in terms of the energy density, ρ_ν , due to neutrinos before their non-relativistic transition, which is given by

$$\rho_\nu = N_{\text{eff}} \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} \rho_\gamma, \quad (4.8)$$

where ρ_γ is the energy density of photons. The standard model prediction for the non-instantaneous decoupling of three neutrino species is $N_{\text{eff}} = 3.044$ [197–201]. However, the possibility of new degrees of freedom contributing additional dark radiation adding to N_{eff} motivates extending the Λ CDM model to include N_{eff} as a free parameter. Although N_{eff} can be constrained through its imprint on the shape and phases of the BAO oscillations [202–204], DESI primarily contributes constraining power by breaking the degeneracies with H_0 and Ω_m . For the Λ CDM+ N_{eff} model, we obtain a constraint from DESI BAO and CMB data of

$$N_{\text{eff}} = 3.07 \pm 0.17 \quad (\text{DESI (BAO)+CMB}). \quad (4.9)$$

As was the case for the bounds on $\sum m_\nu$, this value differs slightly from the constraint for the same data combination reported in [64] due to the switch from version v1.1 of the ACT CMB lensing likelihood to version v1.2. The addition of the full-shape information leads to

$$N_{\text{eff}} = 3.18 \pm 0.16 \quad (\text{DESI (FS+BAO)+CMB}). \quad (4.10)$$

This amounts to a slight reduction in uncertainty and an upward shift in the central value relative to the constraint from DESI (BAO) + CMB (eq. (4.9)), which can be attributed to the preference of DESI data for high H_0 and low Ω_m .

Finally, we present the constraints on N_{eff} in the $w_0 w_a$ CDM model in table 3. Compared to Λ CDM, the uncertainty on N_{eff} increases only slightly and the constraints remain within 1σ of $N_{\text{eff}} = 3.044$.

5 Modified-gravity constraints

DESI full-shape clustering data are sensitive to the growth of large-scale structure, and can hence constrain deviations from general relativity, which we analyse here. We briefly describe the formalism that we use for the modified-gravity parametrisation and then provide results from DESI alone and in combination with other available data sets.

5.1 Modified-gravity formalism and parameterisation

A common and promising approach to testing deviations from general relativity (GR) is to add physically motivated phenomenological parameters to the perturbed Einstein's gravitational field equations and test the deviations of such parameters from their GR predicted values. Whilst many such modified-gravity (MG) parameterisations have been proposed in the literature (see e.g. the reviews [15, 17, 205] and references therein), we focus here on one that is based on the coupling of gravitational potentials to the source content of spacetime.

In the conformal Newtonian gauge, the flat Friedmann-Lemaître-Robertson-Walker metric with scalar perturbations can be written as

$$ds^2 = a(\tau)^2[-(1 + 2\Psi)d\tau^2 + (1 - 2\Phi)\delta_{ij}dx^i dx^j], \quad (5.1)$$

where Ψ and Φ are two gravitational potentials, τ is the conformal time, x^i and x^j are spatial coordinates, and the sum over i and j is implied.

The Einstein field equations applied to the line element, eq. (5.1), yield two equations describing the coupling and evolution of the gravitational field potentials. The first equation relates the potential Ψ to the space-time sources and reduces at late times (i.e. in the absence of anisotropic stresses) to

$$k^2\Psi = -4\pi G a^2 \mu(a, k) \sum_i \rho_i \Delta_i, \quad (5.2)$$

where ρ_i is the density of matter species i , and Δ_i is its comoving overdensity whose evolution describes the growth of inhomogeneities. The phenomenological MG function $\mu(a, k)$ is added to modify the strength of the gravitational interaction and thus the growth rate of structure.

The second perturbed Einstein equation relates the two gravitational potentials, Ψ and Φ and their coupling to source energy densities and stress shear. In the late-time universe and assuming general relativity, it is expected that the anisotropic stress becomes negligible and that the two potentials are nearly equal at the present time. However, this may not be the case in modified-gravity models, where the two potentials can be different and the equation takes the form $(\Phi - \eta(a, k)\Psi) \approx 0$ where the gravitational slip function $\eta(a, k)$ parametrizes a possible deviation from GR. Combining this expression with eq. (5.2), one gets the equation that is particularly useful for the motion of massless particles in a gravitational field (and hence for, e.g., gravitational lensing) and reads

$$k^2(\Phi + \Psi) = -8\pi G a^2 \Sigma(a, k) \sum_i \rho_i \Delta_i, \quad (5.3)$$

where we introduced the MG function $\Sigma(a, k) \equiv \mu(a, k)(\eta(a, k) + 1)/2$. On the left-hand side, $(\Phi + \Psi)$ is equal to twice the so-called Weyl potential that governs the motion of massless particles while, on the right-hand side, the function $\Sigma(a, k)$ modifies the equation from its general-relativistic form. In general relativity, $\Sigma(a, k) = \mu(a, k) = 1$.

Whilst these MG parameterisations are a general function of a and k , it is challenging to constrain free functions in the MG sector with current data, particularly their scale dependence. Therefore, we limit our analysis to a model with a scale-independent Σ and μ , but allow for their time (scale factor) dependence. We adopt the commonly used time dependence for the MG functions, see e.g. [206, 207], where

$$\mu(a) = 1 + \mu_0 \frac{\Omega_{\text{DE}}(a)}{\Omega_{\Lambda}}, \quad \Sigma(a) = 1 + \Sigma_0 \frac{\Omega_{\text{DE}}(a)}{\Omega_{\Lambda}}, \quad (5.4)$$

where the MG parameters μ_0 and Σ_0 take the value of zero in general relativity. This functional form is motivated by the desire to establish a connection between the observed cosmic acceleration and modification to gravity at late times. Consequently, the time dependence of μ and Σ is set to be proportional to the dark energy density. Whilst these forms of time dependencies have been widely used in the literature and constitute a good basis to compare constraints across many surveys and works, they are not free from limitations and may be less effective at capturing models that depart significantly from the evolution in eq. (5.4). Other parameterisations and discussions, including functional and binning forms, can be found in, e.g., [15, 17, 205]. Note that in the class of modified-gravity models that we consider, the evolution of linear matter growth is scale-independent at late times on scales well below the horizon. As a result, the large-scale scale-independence is expected to be preserved for galaxies in redshift space to a very good approximation; see, e.g., [208].

Finally, we note that our μ - Σ model (ansatz in eq. (5.2) and eq. (5.3)) is defined specifically in linear theory. Efforts to extend these models to non-linear scales are ongoing but are faced with multiple challenges. This is not expected to affect our results from DESI full-shape clustering and BAO, as the scale cuts in our full-shape analysis, $0.02 < k/h \text{ Mpc}^{-1} < 0.20$, ensure that nonlinearities (that is, the one-loop terms in the effective field theory expansion) are small. The `velocileptors` prescription used for our full-shape analysis has been tested against the MG non-linear code `fkpt` [209], showing a good agreement in loop corrections for small deviations from GR. The external data that we use also rely on linear scales: SN Ia

and the primary CMB are manifestly linear; CMB lensing is almost entirely in the linear regime; and finally, for DESY3 (3×2 -pt) analysis we use the same conservative scale cuts as used in the MG analysis by DES [107] that limits the information to linear scales. Therefore, our constraints on modified gravity rely on linear scales where this model is well-defined.

5.2 Constraints on modified gravity

Our constraints on μ_0 and Σ_0 from DESI (with BBN and n_s priors), and DESI in combination with external data, are presented in figure 9 and table 4. We derive our measurements assuming the Λ CDM model for the background evolution. We find that DESI constrains μ_0 to be consistent with the zero value predicted by general relativity for the motion of massive particles and their clustering, yielding:¹⁵

$$\mu_0 = 0.11_{-0.54}^{+0.45} \quad (\text{DESI (FS+BAO)+BBN}+n_{s10}). \quad (5.5)$$

The marginalised mean value is centered close to the GR zero value but with 68% credible interval that still allows for substantial possible deviations around it. This consistency of μ_0 with zero also holds for combinations of DESI with external datasets. We see in figure 9 (left panel) that DESI provides no constraints on the “lensing” MG parameter Σ_0 , leading to a horizontal green band in the μ_0 - Σ_0 plane.

Adding DESI data to CMB (no lensing), and the combination of weak lensing and galaxy clustering from the Dark Energy Survey (DESY3 3×2 -pt), can break degeneracies between cosmological parameters and reduce the uncertainties in Σ_0 . The results from these combinations are consistent with the zero value predicted by GR for Σ_0 and μ_0 . The combination of DESI+CMB-nl+DESY3 (3×2 -pt) gives the following tight constraints on the two MG parameters:

$$\left. \begin{aligned} \mu_0 &= 0.04 \pm 0.22, \\ \Sigma_0 &= 0.044 \pm 0.047, \end{aligned} \right\} \begin{array}{l} \text{DESI (FS+BAO)+CMB-nl+} \\ \text{DESY3 (3} \times \text{2-pt).} \end{array} \quad (5.6)$$

This result, further illustrated on the right panel of figure 9, showcases the gains from adding DESI data: complementing CMB-nl and DESY3 (3×2 -pt) data with DESI full-shape and BAO improves the constraints on μ_0 by a factor of 2.5, and those on Σ_0 by a factor of 2.

Inspecting the CMB-only MG constraints (see table 4), we find the same pattern as in previous studies [107, 129, 210, 211], where constraints on the Σ_0 parameter from *Planck* PR3 are in some tension with the zero value predicted by GR. In [129, 210], this was attributed to the anomalous amount of lensing in the CMB captured with the A_{lens} parameter [191–193]

¹⁵The shaded gray region in the top left of figure 9 results from a hard prior $\mu_0 < 2\Sigma_0 + 1$ that is imposed when running our MCMC chains. As noted in previous works (e.g. [107, 207]), this is necessary to avoid an MG parameter space where MG software based on **CAMB** encounters numerical errors when integrating the evolution of perturbations. However, this prior is of no consequence for the above results and their interpretation: while the horizontal green DESI band and the almost vertical orange CMB contour approach and just hit against this prior as shown in the figure, the other data set combinations have smaller contours and are not affected by the prior.

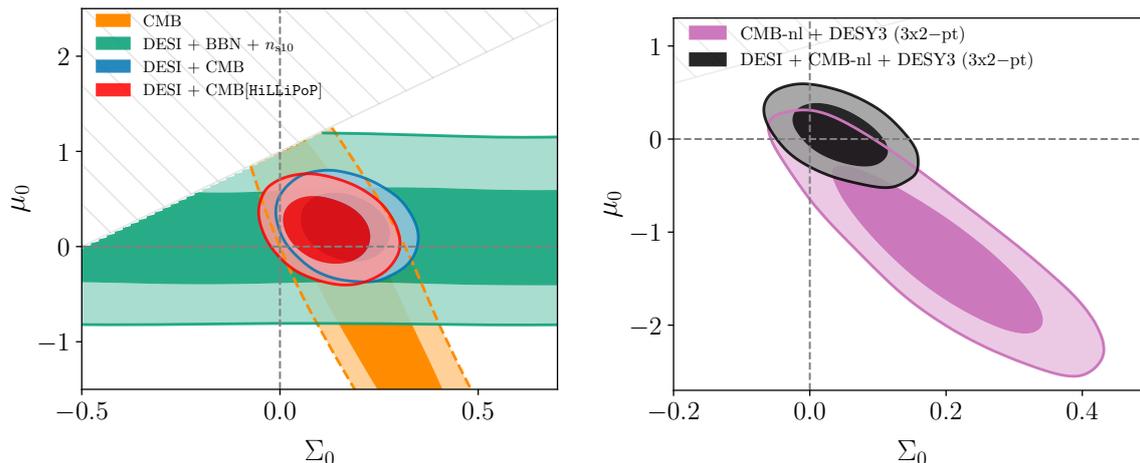


Figure 9. 68% and 95% credible-interval constraints on modified-gravity parameters μ_0 and Σ_0 , assuming a Λ CDM background. *Left panel:* we show the constraints from the CMB alone in orange, and those from DESI full-shape clustering and BAO in green. The blue contour shows the constraints from the combination of DESI and the CMB with our fiducial `plik` likelihood, while the red contour shows the same with the alternate `LoLLiPoP-HiLLiPoP` CMB likelihood. Note that DESI data constraints are consistent with, and centered around, the general relativistic value of $\mu_0 = 0$, and this consistency is maintained once external data is added. Whilst DESI alone does not directly constrain Σ_0 , when it is added to other datasets, like the CMB, it breaks other parameter degeneracies and helps tighten the constraints on Σ_0 (see full discussion of results in section 5.2). *Right panel:* the purple contour shows the combination of CMB-nl with the galaxy clustering and weak lensing (3×2 -pt) likelihood from DESY3, whilst the black contour shows the same combination with the addition of DESI. In both panels, the shaded area on the top left shows the hard prior $\mu_0 < 2\Sigma_0 + 1$ that is imposed due to computational limitations of publicly-available modified-gravity codes based on `CAMB`; this prior does not affect our main results from the combination of probes (see section 5.2).

model/dataset	Ω_m	σ_8	H_0 [km/s/Mpc]	μ_0	Σ_0
Flat $\mu_0\Sigma_0\Lambda$CDM					
DESI (FS+BAO)+ BBN+ n_{s10}	0.2957 ± 0.0097	0.839 ± 0.034	68.53 ± 0.75	$0.11^{+0.45}_{-0.54}$	no constraint
CMB-nl	0.3041 ± 0.0093	$0.742^{+0.13}_{-0.092}$	68.21 ± 0.71	$-0.66^{+1.5}_{-0.83}$	$0.47^{+0.16}_{-0.22}$
CMB-nl [HiLLiPoP]	0.3060 ± 0.0076	$0.737^{+0.13}_{-0.084}$	67.93 ± 0.57	$-0.73^{+1.4}_{-0.79}$	$0.23^{+0.13}_{-0.20}$
DESI+CMB-nl	0.2985 ± 0.0055	0.822 ± 0.024	68.63 ± 0.43	0.23 ± 0.24	$0.388^{+0.11}_{-0.086}$
DESI+CMB	0.3023 ± 0.0053	0.824 ± 0.024	68.32 ± 0.41	0.21 ± 0.24	0.166 ± 0.074
DESI+CMB-nl [HiLLiPoP]	0.3006 ± 0.0051	0.824 ± 0.024	68.33 ± 0.40	0.22 ± 0.24	$0.148^{+0.097}_{-0.12}$
DESI+CMB [HiLLiPoP]	0.3028 ± 0.0050	0.825 ± 0.024	68.18 ± 0.38	0.18 ± 0.24	$0.119^{+0.068}_{-0.076}$
DESI+CMB-nl+DESY3 (3×2 -pt)	0.3027 ± 0.0051	0.808 ± 0.023	68.28 ± 0.40	0.04 ± 0.22	0.044 ± 0.047

Table 4. Constraints on modified-gravity parameters μ_0 and Σ_0 from DESI (FS+BAO) data alone (with the usual BBN and n_{s10} priors), CMB alone, and DESI in combination with external datasets. We assume the flat Λ CDM model for the background. We quote marginalised means and 68% credible intervals in each case.

which we already discussed in our section 4 above. Adding the CMB lensing reconstruction alleviates this tension, as also found in [129, 210]. Although this trend is not driven by DESI, the addition of DESI data to *Planck* breaks parameter degeneracies and makes the preference for nonzero Σ_0 stronger (at the 3σ level). However, as shown by the red contours in figure 9 and our numbers in table 4, we find that this tension goes away when using the more recent CMB likelihoods for low- ℓ LoLLiPoP and high- ℓ HiLLiPoP which also recently alleviated the problem of the anomalous A_{lens} [92, 93]. Therefore, we find that this tension is indeed linked to the lensing-anomaly issue when the *Planck* PR3 likelihood is used. We illustrate and further discuss this point in our paper dedicated to detailed modified gravity analyses [69].¹⁶

It is worth noting that combining DESI with the CMB and weak lensing and galaxy clustering from the DES gives constraints on MG parameters that are comparable to those derived in [107] from a similar (but pre-DESI) combination of probes.¹⁷ This is due to the fact that the precision on μ_0 with current DESI DR1 is approximately as powerful as that of the entire SDSS-IV dataset from two decades of observations [27]. It is also interesting that, while previous constraints on μ_0 using SDSS-IV exhibited a tension with general relativity at a level slightly over 1σ (see figure 9 in [107]), we do not observe such a preference using DESI data.

In sum, we find that DESI full-shape data constrain the modified-gravity parameter μ_0 , reflecting the sensitivity of the clustering signal to the growth of structure and the motion of massive particles where this parameter is involved. The DESI constraint on μ_0 is consistent with the predictions of general relativity. Whilst DESI does not constrain the parameter Σ_0 , the addition of external probes that are sensitive to gravitational lensing breaks degeneracies and produces tight constraints in the μ_0 - Σ_0 plane. Our constraints on modified-gravity parameters are summarised in figure 9 and table 4, and a further extended analysis is presented in an accompanying paper dedicated to modified gravity [69].

6 Conclusions

This is the second paper on cosmological results based on the 1st release (DR1) of data from the Dark Energy Spectroscopic Instrument (DESI). In the first paper [64], we presented results from the analysis of baryon acoustic oscillations in DR1. In the present paper we add information from the broadband clustering of DESI tracers, which we refer to as the “full-shape” analysis. The major new consequence of adding the full-shape information to the baryon acoustic oscillations (BAO) is that DESI DR1 data now become directly sensitive to the temporal growth of structure, and hence to the amplitude of mass fluctuations σ_8 and other parameters that characterise cosmic growth.

Our data includes clustering from luminous red galaxies, emission line galaxies, quasars, and the Ly α forest observed in an area of 7,500 square degrees (less for some tracers) analysed

¹⁶While this paper was in DESI internal review, the paper [212] appeared on the arXiv, showing similar findings regarding the Σ_0 tension being alleviated when using LoLLiPoP and high- ℓ HiLLiPoP likelihoods, but using a different modified-gravity software than the one we use in our analysis.

¹⁷The combination in the DES extensions paper [107] also included SN Ia which we do not add to our MG external-data combination, as the DESI and SN Ia values of Ω_m are in tension in Λ CDM background (and remain so in the $\mu_0\Sigma_0\Lambda$ CDM model).

in six redshift bins in the range $0 < z < 4$. The full-shape methodology has been thoroughly validated in a series of supporting papers, leading to decisions on scale cuts and the treatment and parameterisation of systematic errors. These decisions are incorporated in the full-shape likelihood, which is subsequently combined with the BAO likelihood from [64]. We refer to the results from the total likelihood as DESI (FS+BAO). When not combining with the CMB data, we complement this combination of full-shape and BAO data from DESI with the baryon density ($\Omega_b h^2$) prior from big bang nucleosynthesis, as well as a loose prior on the scalar spectral index n_s .

Assuming the Λ CDM cosmological model, we find that the combination of DESI full-shape and BAO data pins down matter density to $\Omega_m = 0.2962 \pm 0.0095$, a $\sim 3\%$ measurement that is in general agreement with measurements from other cosmological probes. The amplitude of mass fluctuations is $\sigma_8 = 0.842 \pm 0.034$, and we also constrain the derived parameter $S_8 \equiv (\Omega_m/0.3)^{0.5} = 0.836 \pm 0.035$; these measurements are in excellent agreement with previous galaxy-clustering analyses as well as those from the CMB, and slightly higher than, albeit generally consistent with, constraints from weak gravitational lensing. The Hubble constant we get from DESI combined with the BBN and weak n_s prior is $H_0 = (68.56 \pm 0.75) \text{ km s}^{-1} \text{ Mpc}^{-1}$, in concordance with CMB and previous BAO measurements, and in continuing disagreement with the much higher values obtained by inferences from the local universe [152].

When DESI is combined with external data while still assuming Λ CDM, the constraints tighten while generally remaining in concordance with DESI-only results. The addition of the combination of cosmic shear, galaxy clustering, and CMB lensing from the Dark Energy Survey Year-3 data — the DESY3 (6×2 -pt) analysis — shifts the σ_8 (and S_8) values downward by about one standard deviation, while tightening the error bars by about a factor of 2 (and 3). When, instead, the CMB data are added to DESI, the errors in key cosmological parameters tighten even more. Finally, the most comprehensive combination that we consider in Λ CDM, when DESI full-shape and BAO data are combined with the CMB and the DESY3 (6×2 -pt) analysis, leads to a significant tightening of the errors and to parameter determinations with 1% precision in Ω_m , 0.6% in σ_8 and S_8 , and 0.4% precision in H_0 .

We next study $w_0 w_a$ CDM, the model that allows a time-varying equation of state of dark energy. The combination of full-shape and BAO data with the CMB and type Ia supernovae leads to a tightening of $\sim 20\%$ in the area in the w_0 – w_a plane relative to the same combinations when the full-shape information is not included. The preference for a departure from the Λ CDM prediction ($w_0 = -1, w_a = 0$) remains, and we find that the best-fit w_0 – w_a model is favored by $\Delta\chi_{\text{MAP}}^2 = -8.8$ when DESI (FS+BAO)+CMB is combined with PantheonPlus supernovae, by $\Delta\chi_{\text{MAP}}^2 = -14.5$ for the combination with Union3, and by $\Delta\chi_{\text{MAP}}^2 = -17.5$ for the combination with DES-SN5YR (here MAP refers to the maximum *a posteriori* parameter values at which the respective fits are evaluated). These correspond to preferences for $w_0 w_a$ CDM over Λ CDM at the significance levels of 2.5σ (PantheonPlus), 3.4σ (Union3), and 3.8σ (DES-SN5YR). These combined constraints therefore continue to show preference for a departure from the Λ CDM model at very similar statistical levels as the same combinations with DESI BAO alone in [64]. We also check that the aforementioned combined-probe measurements of Ω_m , σ_8/S_8 , and H_0 in the $w_0 w_a$ CDM model remain consistent with those in the Λ CDM model, with only a modest degradation in precision when the more general $w_0 w_a$ CDM background is allowed.

Full-shape information allows us to improve the constraints on the sum of the neutrino masses, as neutrinos affect not only the geometry but also the growth of cosmic structure. We find an upper limit of $\sum m_\nu < 0.409 \text{ eV}$ at 95% confidence from DESI full-shape and BAO data combined with the BBN and loose n_s priors. When DESI is combined with the CMB, we obtain $\sum m_\nu < 0.071 \text{ eV}$ (again at 95%), a constraint that is $\sim 15\%$ stronger than that with DESI BAO data alone combined with the CMB. This strong limit arises from the preference of both DESI BAO and DESI full-shape data for high values of H_0 and low values of Ω_m , which suppress the value of $\sum m_\nu$ due to parameter degeneracy. The upper limit on $\sum m_\nu$ is negatively correlated with the amount of lensing observed in the CMB data, and weakens if we adopt CMB likelihoods that show less evidence for excess lensing in the CMB.

Finally, DESI full-shape data and its sensitivity to the growth of structure allow us to test the theory of gravity. We study a model where departures from general relativity are modeled by two modified-gravity parameters, μ_0 and Σ_0 . DESI full-shape and BAO data constrain the parameter that governs the clustering of massive particles, $\mu_0 = 0.11^{+0.45}_{-0.54}$, which is consistent with the zero value predicted by general relativity. DESI alone is insensitive to the other modified-gravity parameter, Σ_0 , that governs the motion of massless particles, but helps constrain it when combined with external data. DESI, in combination with CMB data, the combined lensing and clustering (3×2 -pt) analysis from the DES Year-3 observations, and DES-SN5YR supernova data, finds $\mu_0 = 0.04 \pm 0.22$ and $\Sigma_0 = 0.044 \pm 0.047$. These results are consistent with the zero-value predictions of general relativity. Interestingly, we find that the combined constraint on the parameter Σ_0 is a factor 4.7 better than that on the parameter μ_0 to which the full-shape analysis is sensitive, which indicates that forthcoming DESI data should be very effective in reducing the uncertainty on the latter parameter.

Whilst this paper wraps up the key cosmological results from DESI first data release (DR1), many ongoing or recently completed DESI projects complete the picture by studying some of the aforementioned results in more detail, or presenting complementary cosmological and astrophysical analyses. Looking ahead, BAO and full-shape analyses to follow from the three years of DESI observations are expected to contribute major new information, provide improved constraints on the cosmological parameters and models discussed in this paper, and shed new insights into dark energy, modified gravity, and neutrino mass.

7 Data availability

Data from the plots in this paper will be available on Zenodo as part of DESI's Data Management Plan.

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The DESI Legacy Imaging Surveys consist of three individual and complementary projects: the Dark Energy Camera Legacy Survey (DECaLS), the Beijing-Arizona Sky Survey (BASS), and the Mayall z-band Legacy Survey (MzLS). DECaLS, BASS and MzLS together include data obtained, respectively, at the Blanco telescope, Cerro Tololo Inter-American Observatory, NSF NOIRLab; the Bok telescope, Steward Observatory, University of Arizona; and the Mayall telescope, Kitt Peak National Observatory, NOIRLab. NOIRLab is operated by the Association of Universities for Research in Astronomy (AURA) under a cooperative agreement with the National Science Foundation. Pipeline processing and analyses of the data were supported by NOIRLab and the Lawrence Berkeley National Laboratory. Legacy Surveys also uses data products from the Near-Earth Object Wide-field Infrared Survey Explorer (NEOWISE), a project of the Jet Propulsion Laboratory/California Institute of Technology, funded by the National Aeronautics and Space Administration. Legacy Surveys was supported by: the Director, Office of Science, Office of High Energy Physics of the U.S. Department of Energy; the National Energy Research Scientific Computing Center, a DOE Office of Science User Facility; the U.S. National Science Foundation, Division of Astronomical Sciences; the National Astronomical Observatories of China, the Chinese Academy of Sciences and the Chinese National Natural Science Foundation. LBNL is managed by the Regents of the University of California under contract to the U.S. Department of Energy. The complete acknowledgments can be found at <https://www.legacysurvey.org/>.

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The authors are honored to be permitted to conduct scientific research on Iolkam Du’ag (Kitt Peak), a mountain with particular significance to the Tohono O’odham Nation.

A Parameter projection effects

In this appendix we give some insight into the parameter projection effects, also known as “prior volume effects”, that influence our study of the w_0w_a CDM model. Projection effects typically occur in the presence of long degeneracy directions in parameter space, when the posterior exhibits strong non-Gaussianity. In those cases, the mean of the marginalised posterior can be significantly offset from the maximum of the posterior (the maximum *a posteriori* (MAP) value). Strong projection effects do not indicate any problem in the analysis, but simply the fact that the statistical constraints on relevant cosmological parameters are weak due to the presence of degeneracies, which can be created by the correlation between the cosmological and nuisance parameters. We have encountered projection effects in our analysis, notably when studying the w_0w_a CDM model. We have not shown in this paper any cosmological results that are subject to strong projection effects, but we now discuss one such result in order to illustrate the effect in this context.

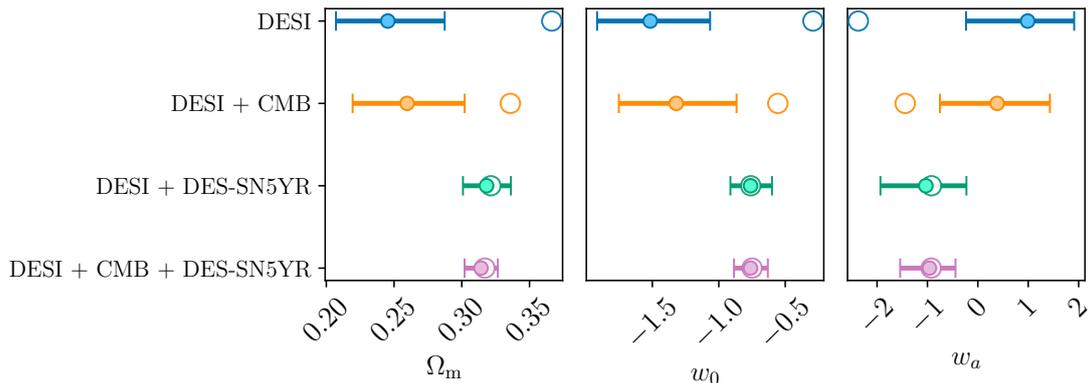


Figure 10. Parameter projection effects in the w_0w_a CDM cosmological model. The solid horizontal lines show the 95% marginalised posteriors, with the solid circle being the mean of the corresponding marginalised posterior, while the open circles show the *maxima* of the corresponding posteriors (MAP values). Note that both DESI alone, and DESI in combination with CMB, have MAP values that are far from the mean of the marginalised posterior, and hence exemplify strong projection effects. However, if DESI data are combined with type Ia supernovae (represented by the DES-SN5YR dataset here), the projection effect disappears and the marginalised means and MAP values agree well. The same result is found for the DESI+CMB+DES-SN5YR combination.

In figure 10 we show the 1D marginalised posteriors (with 95% credible intervals shown) on Ω_m , w_0 , and w_a in the w_0w_a CDM model. We show the results for four combinations of datasets: DESI full-shape clustering and BAO data alone (with the usual BBN and n_{s10} priors), DESI combined with the CMB, DESI combined with the DES Year-5 type Ia supernova dataset, and finally the DESI+CMB+DES-SN5YR combination. In each case and for each parameter, we also show the MAP value as an open circle. We see that the DESI and DESI+CMB combinations are strongly affected by projection effects, as the MAP values for all three parameters lie outside the corresponding Bayesian 95% interval. The conclusion of this analysis is that DESI (FS+BAO) — unlike DESI BAO only — and CMB, separately or in combination, do not give strong constraints on the dark-energy sector in w_0w_a CDM model due to the presence of multiple nuisance parameters in the FS analysis.

However, once we add SN Ia data (in this case, DES-SN5YR) to DESI, the projection effects disappear. Figure 10 shows that the MAP values are now consistent with the Bayesian posterior, due to the more effective breaking of degeneracy in the dark energy sector when SN Ia are added to DESI, relative to the case when CMB data is added instead. The improved constraints — and hence lower projection effects — with the inclusion of SN Ia data, rather than the CMB, can be attributed to the fact the SN Ia are more incisive on the acceleration and complementary to galaxy clustering [213] and they produce tighter constraints in the Ω_m - w_0 - w_a space. Additionally, current SN Ia observations prefer higher values of Ω_m , which further squeezes the contours in the w_0 - w_a plane. We therefore only consider the dark-energy constraints from DESI+SN Ia (and DESI+CMB+SN Ia etc) combinations as useful to report in section 3.2 of this paper. Further discussion of projection effects is given in the companion methodology papers [1, 70].

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The DESI collaboration

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